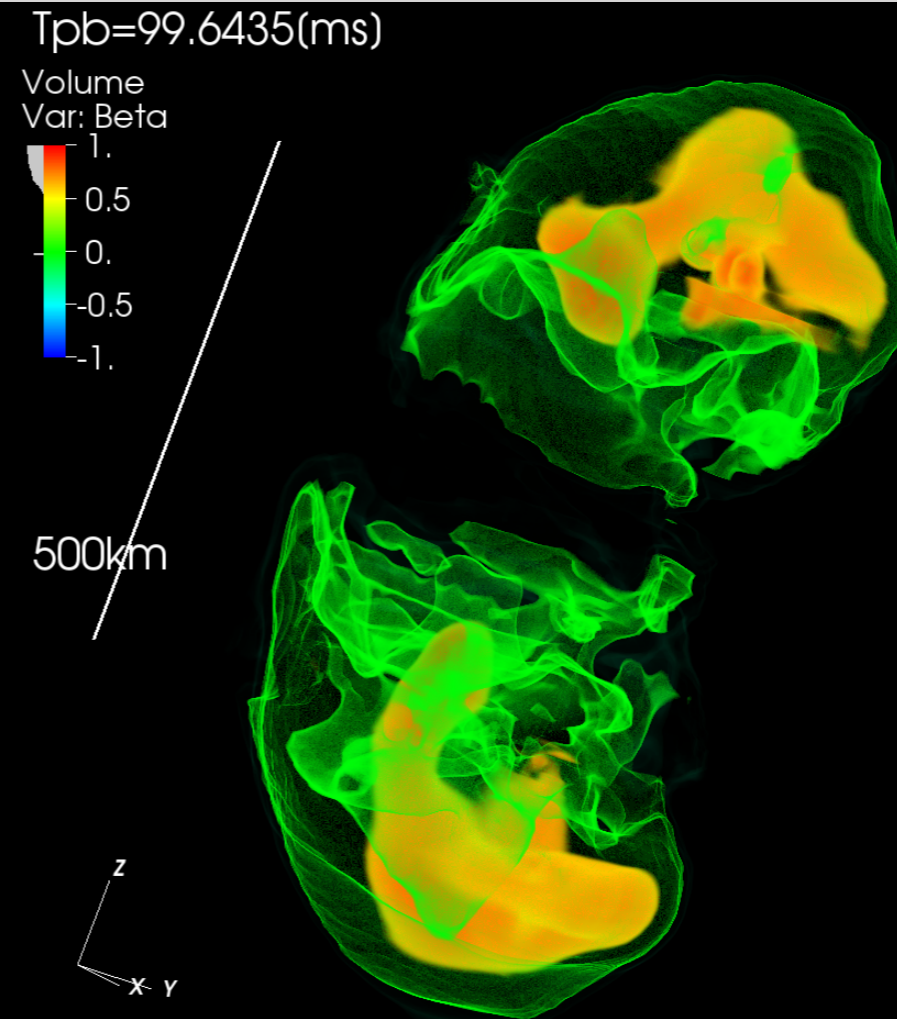
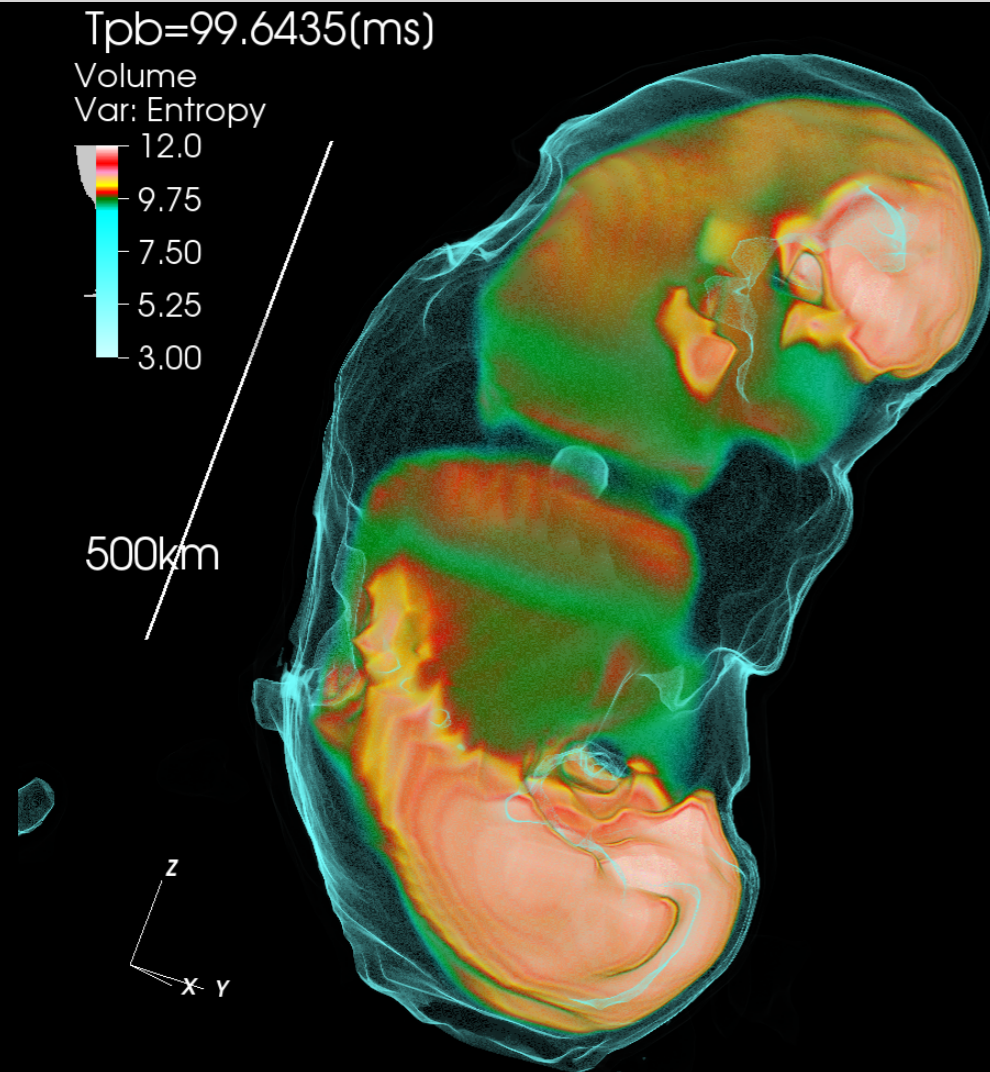


Magneto rotational explosion of a massive star supported by neutrino heating in full 3DGR simulation and its nucleosynthesis



- Takami Kuroda (TU Darmstadt)
- Moritz Reichert (TU Darmstadt)
- Marius Eichler (TU Darmstadt)
- Almudena Arcones (TU Darmstadt)
- Tomoya Takiwaki (NAOJ)
- Kei Kotake (Fukuoka Univ.)



TECHNISCHE
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DARMSTADT



European Research Council
Established by the European Commission

理論懇, 国立天文台, 25-27/12/2019

(1) The standard explosion mechanism

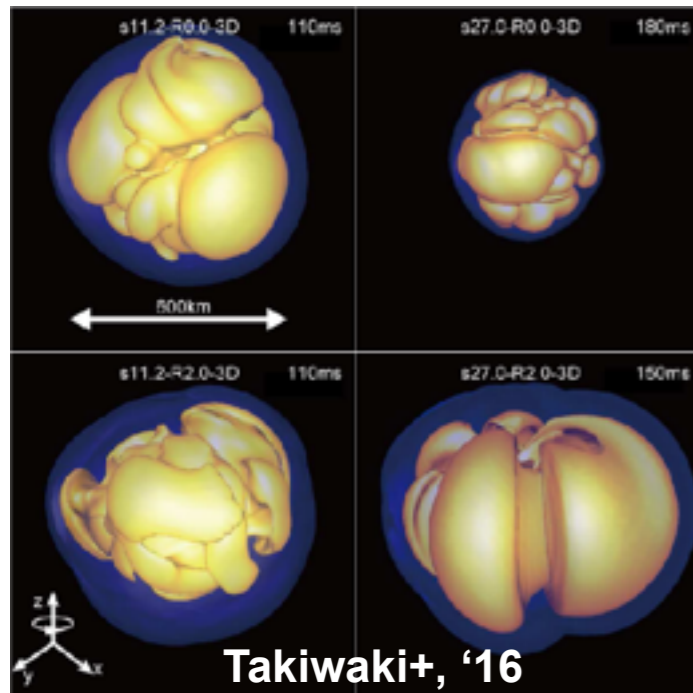
Neutrino driven explosion,

Colgate&White'66, Bethe&Wilson'85

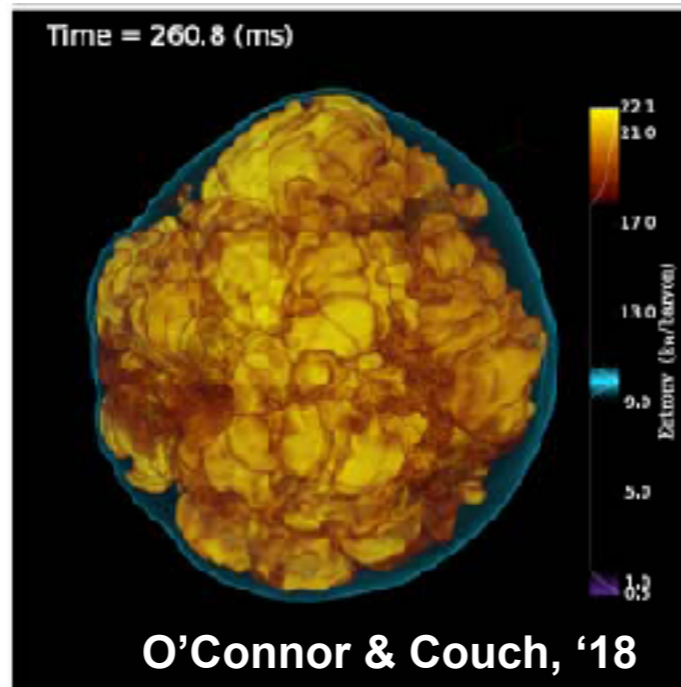
For reviews, Janka'12, Kotake+, '12, Burrows+, '13

**~99% of internal energy
is radiated away
via neutrinos ($\sim 10^{53}$ ergs)**

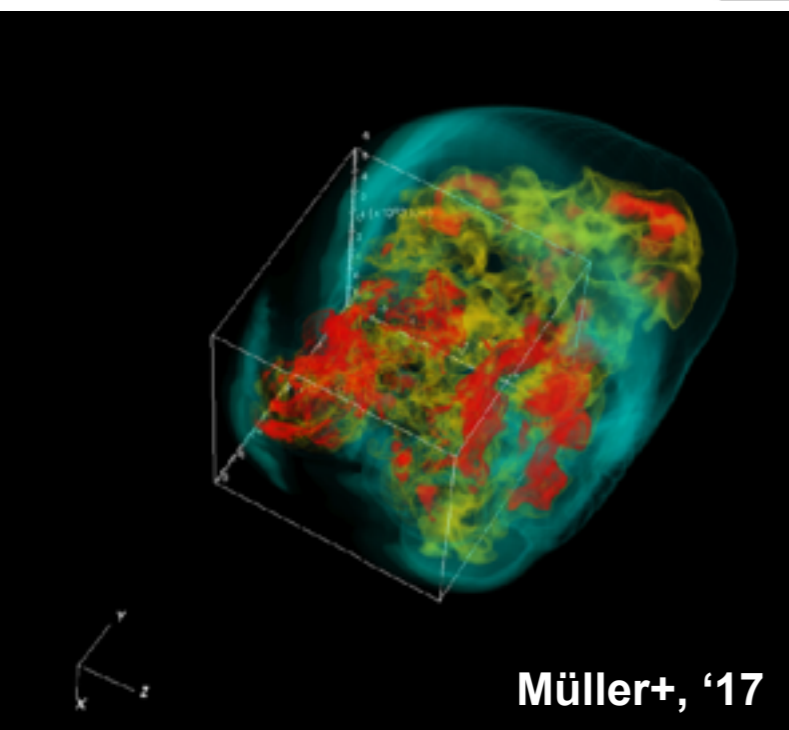
**→ ~10% energy deposition
is enough to explain
 $E_{\text{exp}} \sim 10^{51}$ ergs**



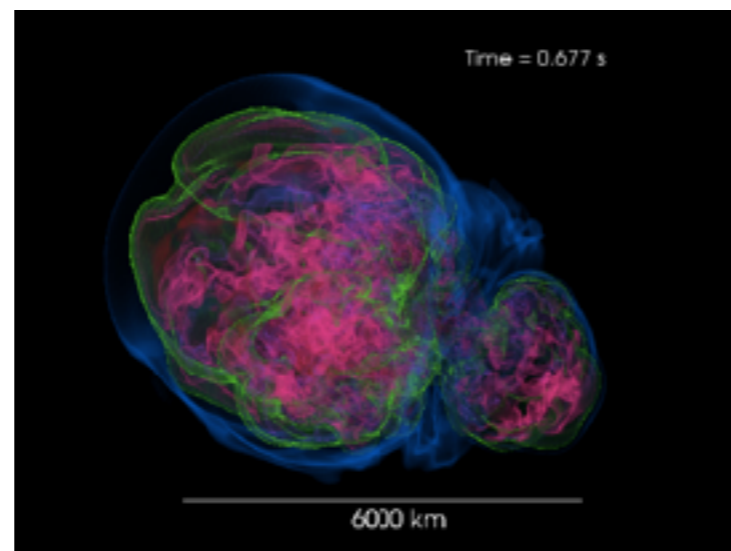
Takiwaki+, '16



O'Connor & Couch, '18



Müller+, '17



Vartanyan+, '19

**Numerical simulations
can't fully explain
canonical explosion energies**

- neutrino-matter interactions?
- resolution problem?
- too short simulation time?

(1) The standard explosion mechanism Neutrino driven explosion,

Colgate&White'66, Bethe&Wilson'85

For reviews, Janka'12, Kotake+, '12, Burrows+, '13

(2) If the magnetic field is strong enough

MHD explosion,

- Angular momentum transfer
- Mass ejection by B pressure
- efficient neutrino heating

2D: Ardeljan+, '00, Kotake+, '04, Obergaulinger+, '06, '17, Burrows+, '07, Takiwaki+, '09,

3D: Mikami+, '08;

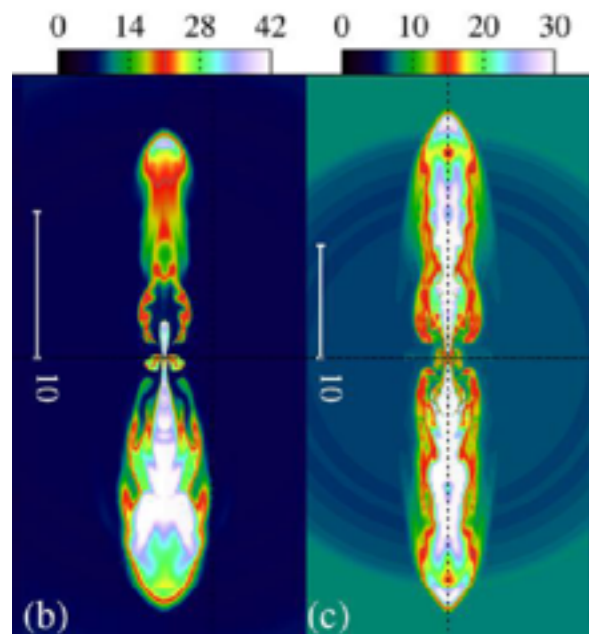
Mösta+, '14;

Obergaulinger+, '19;

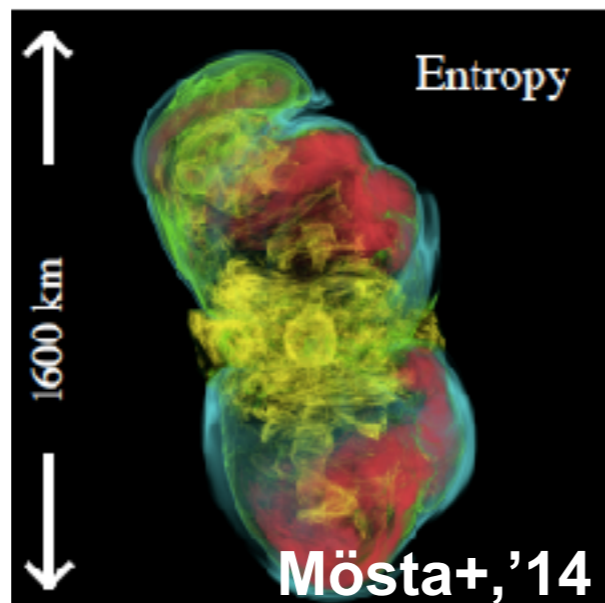
Newtonian, no neutrino, Polytropic EOS

full GR but very simplified neutrino transport

SR with M1 neutrino transport (preliminary result)



Obergaulinger+, '17

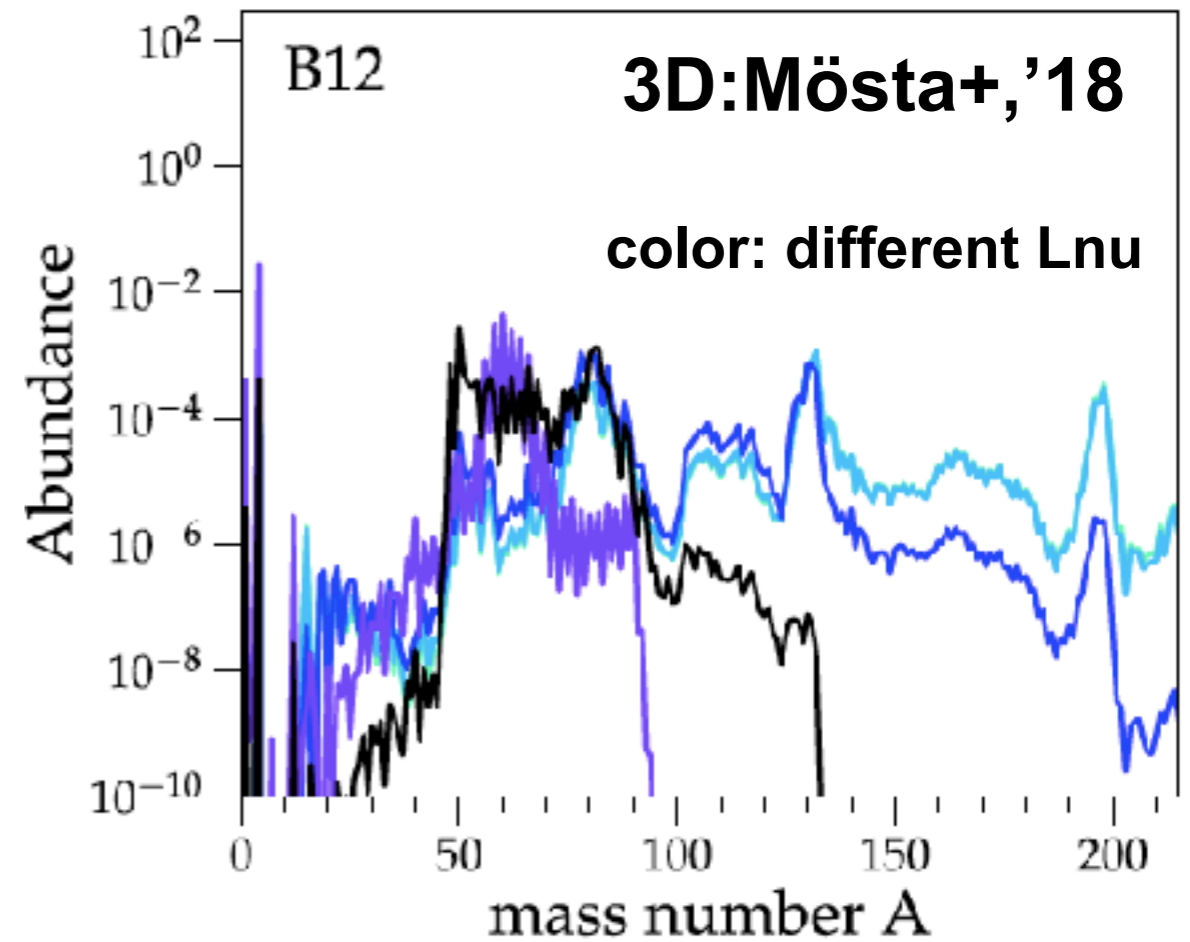
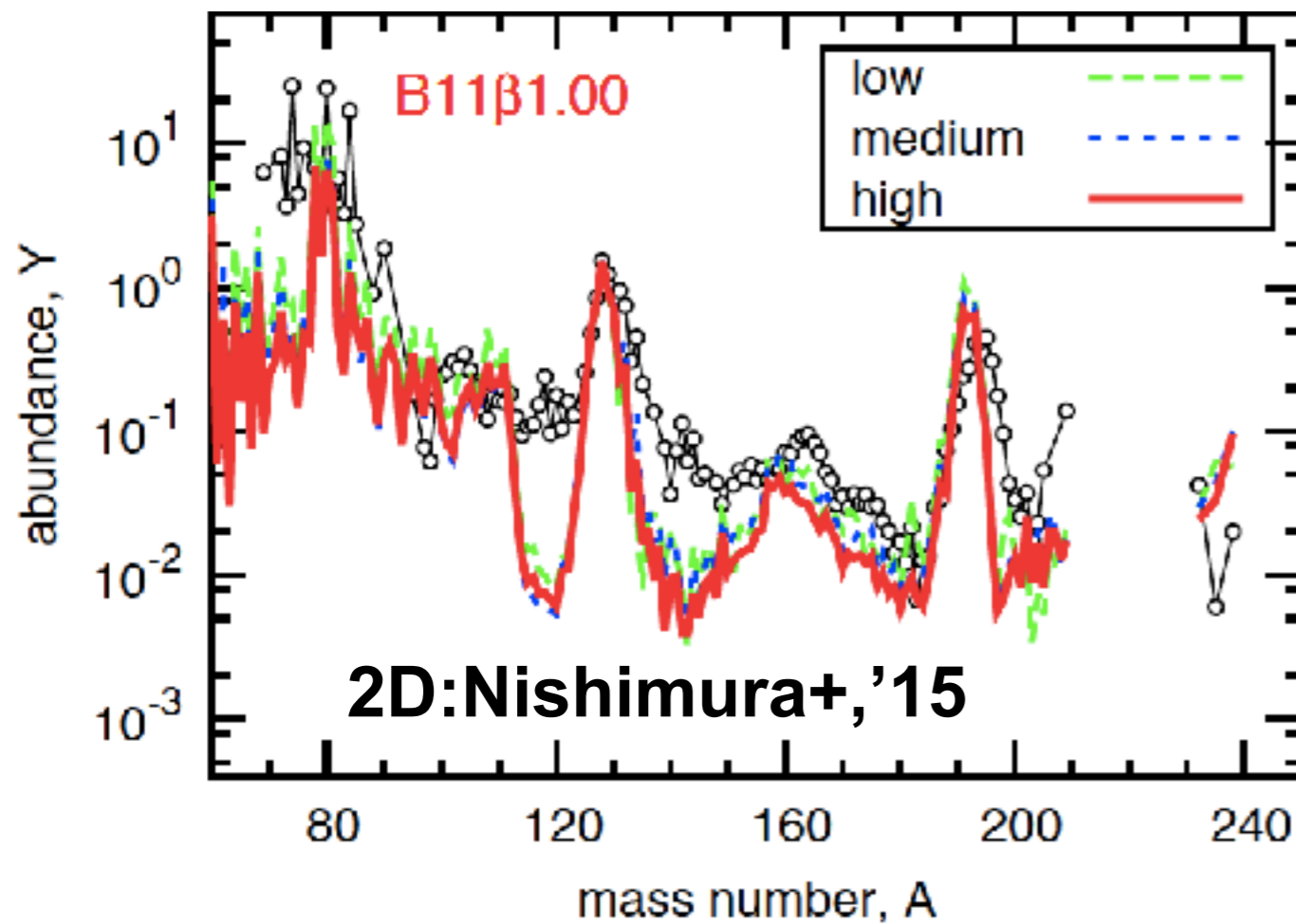


Mösta+, '14

Things still to be explored

- resolution problem -> MRI?
- 2D artefacts -> 3D non-axisymmetries
- microphysics -> neutrino effects

Can the MHD explosion be the r-process site?



- ✱ **Strong MHD jet can potentially produce the 3rd peak (due to low- Y_e ejecta)**
- ✱ **Neutrino radiation significantly influences on the ejecta Y_e**

BSSN equations (17 variables): 4th order accuracy in space and time

$$\begin{aligned}
 (\partial_t - \mathcal{L}_\beta)\tilde{\gamma}_{ij} &= -2\alpha\tilde{A}_{ij} \\
 (\partial_t - \mathcal{L}_\beta)\phi &= -\frac{1}{6}\alpha K \\
 (\partial_t - \mathcal{L}_\beta)\tilde{A}_{ij} &= e^{-4\phi} \left[\alpha(R_{ij} - 8\pi\gamma_{i\mu}\gamma_{j\nu}T_{(\text{total})}^{\mu\nu} - D_i D_j \alpha) \right]^{\text{trf}} + \alpha(K\tilde{A}_{ij} - \tilde{A}_{ij}K) \\
 (\partial_t - \mathcal{L}_\beta)\tilde{\Gamma}^i_k &= -\Delta\alpha + \alpha(\tilde{A}_{ij}\tilde{A}^{ij} + K^2/3) + 4\pi\alpha(n_\mu n_\nu T_{(\text{total})}^{\mu\nu} + \gamma^{ij}\gamma_{i\mu}\gamma_{j\nu}T_{(\text{total})}^{\mu\nu} \\
 &\quad - 16\pi\tilde{\gamma}^{ij}\gamma_{i\mu}n_\nu T_{(\text{total})}^{\mu\nu}) \\
 &\quad + \alpha\left(\frac{2}{5}\tilde{\gamma}^{ij}K_{,j} - 6\tilde{A}^{ij}\phi_{,j} - \tilde{\Gamma}^i_{jk}\tilde{A}^{jk}\right) \\
 &\quad - \frac{1}{3}\tilde{\gamma}^{ij}\beta^k_{,kj} - \tilde{\Gamma}^j\beta^i_{,j} + \frac{2}{3}\tilde{\Gamma}^i\beta^j_{,j} - 2\tilde{A}^{ij}\phi_{,j}
 \end{aligned}$$

determines motion

creates curvature

creates curvature

determines motion

GRMHD (3 variables): 3rd order accuracy in space and time, respectively

GR-Rad (4 variables): 4th order accuracy in space and time, respectively

$$\begin{aligned}
 \partial_t\sqrt{\gamma}S_i + \partial_j\sqrt{\gamma}(S_iv^j) &= -\sqrt{\gamma}[S_0\partial_i\alpha - S_k\partial_i v^k - 2\alpha S_k^k\partial_i\phi \\
 &\quad + \alpha e^{-4\phi}(S_{jk} - P\gamma_{jk})\partial_i\tilde{\gamma}^{jk}/2 + \dots] \quad (9)
 \end{aligned}$$

$$\begin{aligned}
 \partial_t\sqrt{\gamma}E_{(\epsilon)} + \partial_i\sqrt{\gamma}(\alpha F_{(\epsilon)}^i) &+ \sqrt{\gamma}\alpha\partial_\epsilon(\epsilon\tilde{M}_{(\epsilon)}^\mu n_\mu) \\
 \sqrt{\gamma}(\alpha P_{(\epsilon)}^{ij}K_{ij} - F_{(\epsilon)}^i\partial_i\alpha - \alpha S_{(\epsilon)}^\mu n_\mu), &
 \end{aligned}$$

$$\begin{aligned}
 \partial_t\sqrt{\gamma}\tau + \partial_i\sqrt{\gamma}(\tau v^i + P(v^i - v^i)) &= \sqrt{\gamma}[\alpha K S_k^k/3 + \alpha e^{-4\phi}(S_{ij} - P\gamma_{ij})\tilde{A}^{ij} \\
 &\quad - S_i D^i\alpha + \alpha \int d\epsilon S_{(\epsilon)}^\mu n_\mu], \quad (11)
 \end{aligned}$$

$$\begin{aligned}
 \partial_t\sqrt{\gamma}E_{(\epsilon)} + \partial_i\sqrt{\gamma}(\alpha P_{(\epsilon)}^i{}^j - \beta^j F_{(\epsilon)i}) &- \sqrt{\gamma}\alpha\partial_\epsilon(\epsilon\tilde{M}_{(\epsilon)}^\mu \gamma_{i\mu}) \\
 - F_{(\epsilon)j}\partial_i\beta^j + (\alpha/2)P_{(\epsilon)}^{jk}\partial_i\gamma_{jk} + \alpha S_{(\epsilon)}^\mu \gamma_{i\mu}, &
 \end{aligned}$$

neutrino cooling/heating

In GRMHD code, one solves these 3 systems with (26+12*N_{ene}) variables satisfying the Hamiltonian, momentum, & no-monopole constraints

$$\begin{aligned}
 \partial_t(\rho_* Y_e) + \partial_i(\rho_* Y_e v^i) &= \sqrt{\gamma}\alpha m_u \int \frac{d\epsilon}{\epsilon} (S_{(\nu_e, \epsilon)}^\mu - S_{(\bar{\nu}_e, \epsilon)}^\mu) u_\mu, \\
 \partial_t B^i &= \partial_k (B^k v^i - B^i v^k) \quad (12)
 \end{aligned}$$

The basic equations for neutrino transport

$$T_{\mu\nu}^{neutrino} = E n_{\mu} n_{\nu} + F_{\mu} n_{\nu} + F_{\nu} n_{\mu} + P_{\mu\nu} \quad \text{Shibata+'11, TK+'16}$$

(E, F, P: 0th, 1st, 2nd momenta (in Euler))

$$\begin{aligned} & \partial_t \sqrt{\gamma} E_{(\varepsilon)} + \overbrace{\partial_i \sqrt{\gamma} (\alpha F_{(\varepsilon)}^i - \beta^i E_{(\varepsilon)})}^{\text{advection}} + \overbrace{\sqrt{\gamma} \alpha \partial_{\varepsilon} (\varepsilon \tilde{M}_{(\varepsilon)}^{\mu} n_{\mu})}^{\text{gravitational redshift/Doppler}} \\ & = \sqrt{\gamma} (\underbrace{\alpha P_{(\varepsilon)}^{ij} K_{ij} - F_{(\varepsilon)}^i \partial_i \alpha}_{\text{gravitational source}} - \underbrace{\alpha S_{(\varepsilon)}^{\mu} n_{\mu}}_{\text{neutrino-matter interaction}}), \end{aligned} \quad (4)$$

and

$$\begin{aligned} & \partial_t \sqrt{\gamma} F_{(\varepsilon)i} + \partial_j \sqrt{\gamma} (\alpha P_{(\varepsilon)i}^j - \beta^j F_{(\varepsilon)i}) - \sqrt{\gamma} \alpha \partial_{\varepsilon} (\varepsilon \tilde{M}_{(\varepsilon)}^{\mu} \gamma_{i\mu}) \\ & = \sqrt{\gamma} [-E_{(\varepsilon)} \partial_i \alpha + F_{(\varepsilon)j} \partial_i \beta^j + (\alpha/2) P_{(\varepsilon)}^{jk} \partial_i \gamma_{jk} + \alpha S_{(\varepsilon)}^{\mu} \gamma_{i\mu}], \end{aligned} \quad (5)$$

TK+',16

The Opacity Set Included in this Study and their References

Process	Reference	Summarized In
$n\nu_e \leftrightarrow e^-p$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.1
$p\bar{\nu}_e \leftrightarrow e^+n$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.1
$\nu_e A \leftrightarrow e^- A'$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.1
$\nu p \leftrightarrow \nu p$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.2
$\nu n \leftrightarrow \nu n$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.2
$\nu A \leftrightarrow \nu A$	Bruenn (1985), Rampp & Janka (2002)	Appendix A.2
$\nu e^{\pm} \leftrightarrow \nu e^{\pm}$	Bruenn (1985)	Appendix A.3
$e^- e^+ \leftrightarrow \nu\bar{\nu}$	Bruenn (1985)	Appendix A.4
$NN \leftrightarrow \nu\bar{\nu}NN$	Hannestad & Raffelt (1998)	Appendix A.5

Numerical setups

- 20Msun model (WHW07)
- $dx \sim 458\text{m@center}$, $\sim 3.2\text{km@R}=100\text{km}$
- Basic neutrino opacities based on Bruenn'85 (same as TK+, '18)
- $N_{\text{ene}}=12$ bins ($1 < \varepsilon < 300\text{MeV}$)

-
- 3 models (ROB00, R1B00, R1B12)

- SFHo (Steiner+'13)

- Cylindrical rotational law

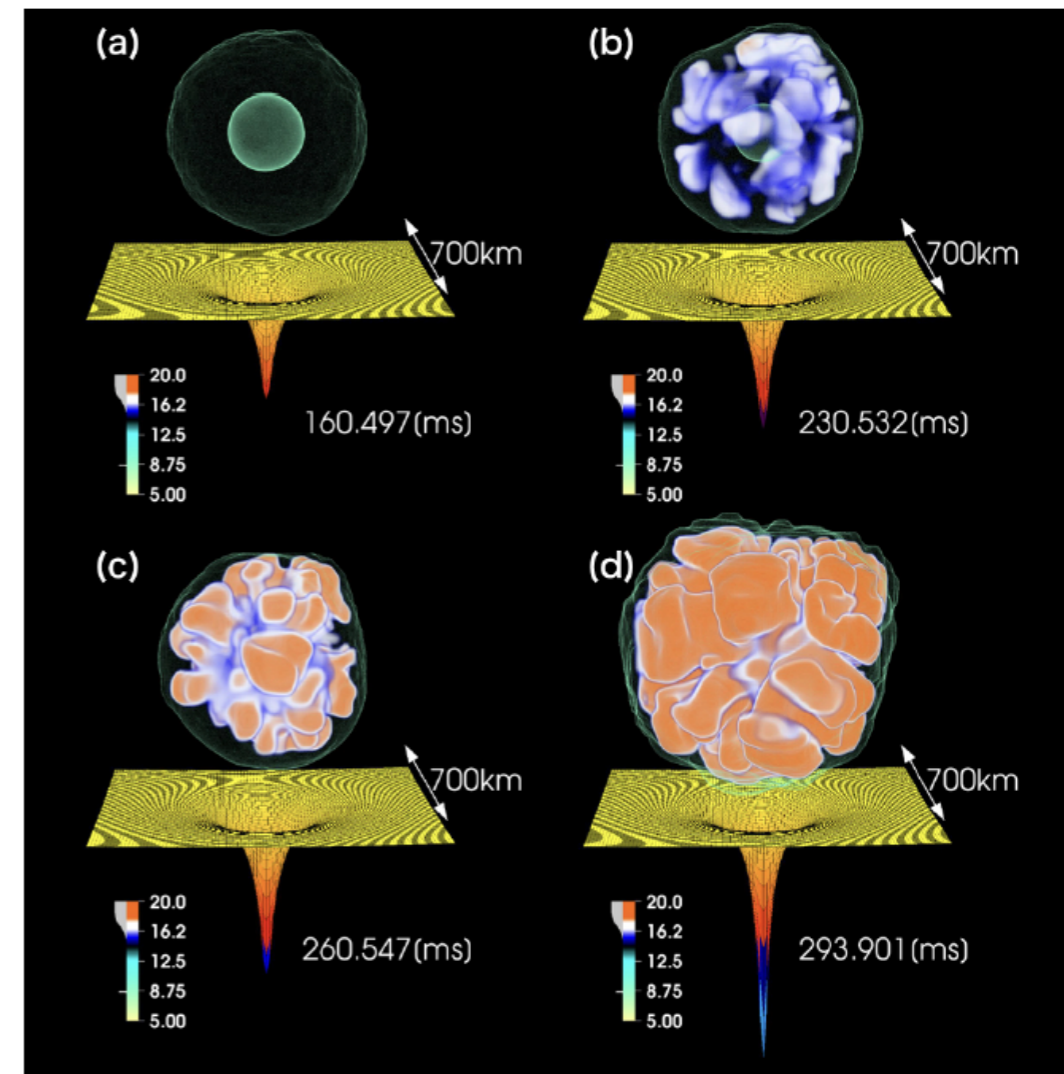
$$\Omega = \Omega_0 \frac{R_0^2}{\varpi^2 + R_0^2} \quad \Omega_0 = 1(\text{rad/s}) \quad (\beta_b \sim 1\%)$$

- Dipole-like B

$$A_\phi = \frac{B_0}{2} \frac{R_0^3}{R^3 + R_0^3} R \sin\theta \quad B_0 = 10^{12}\text{G} \quad (\beta_{\text{mag,b}} \sim 1\%)$$

- CT method for $\text{div}B=0$

-
- XC40 @ NAOJ



TK, Kei Kotake, T. Takiwaki, & F.-K. Thielemann
2018, MNRAS Letter

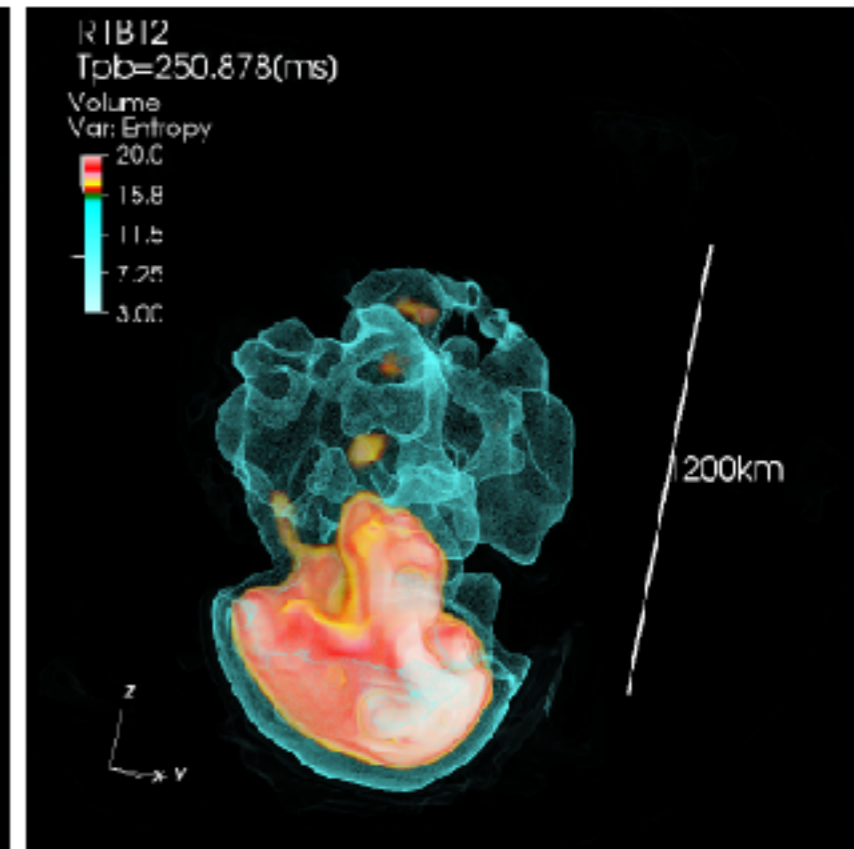
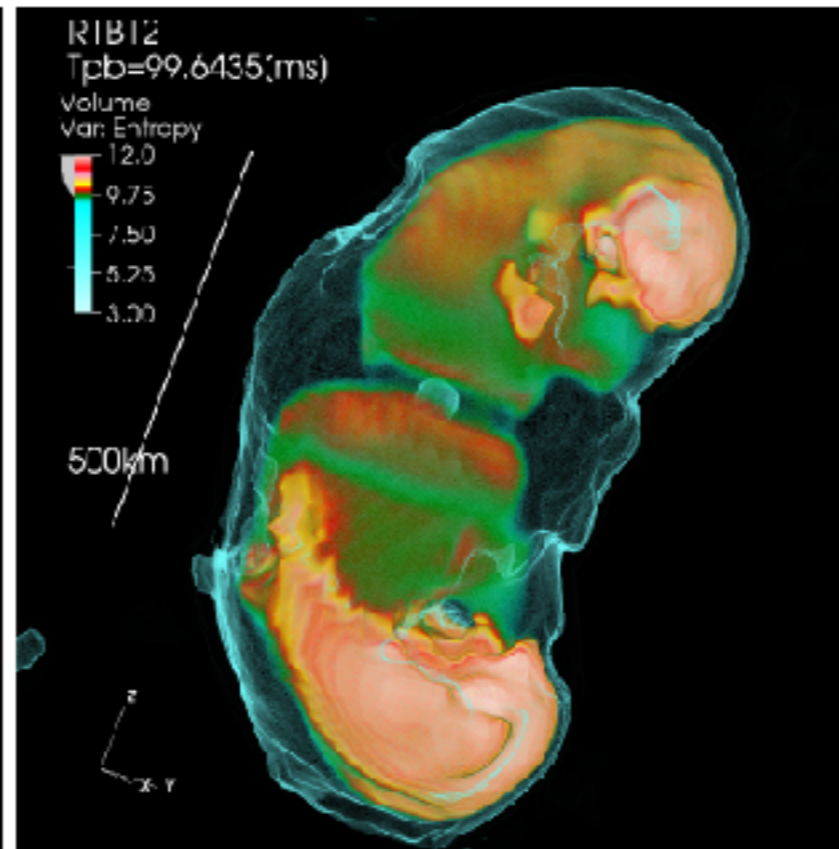
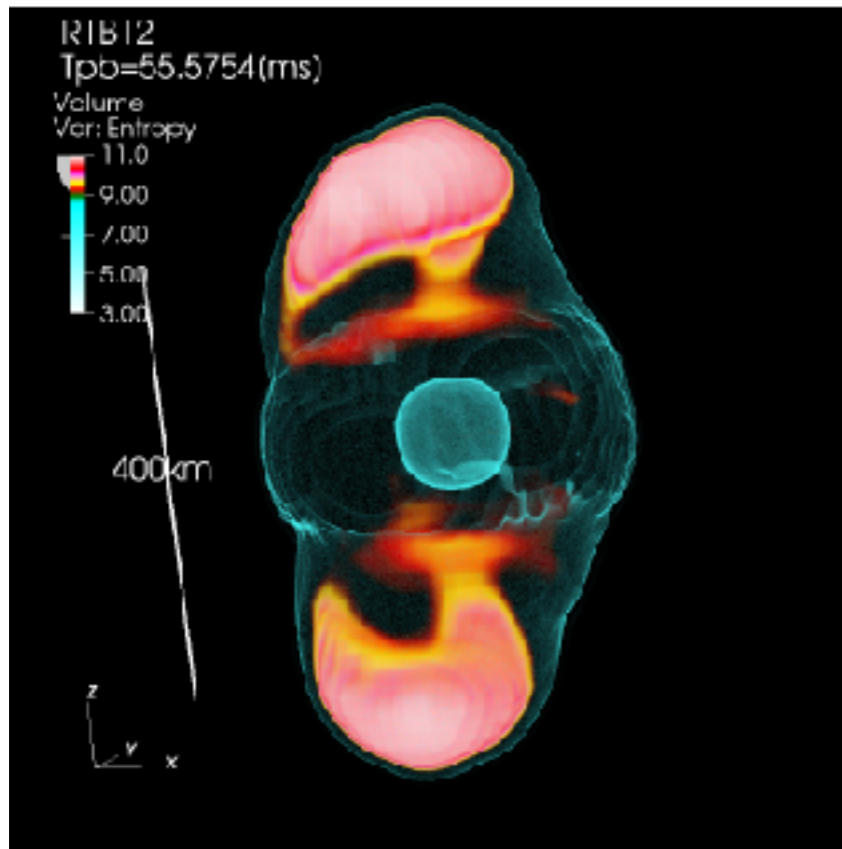
Rotating magnetized model (R1B12)

Tpb=55ms

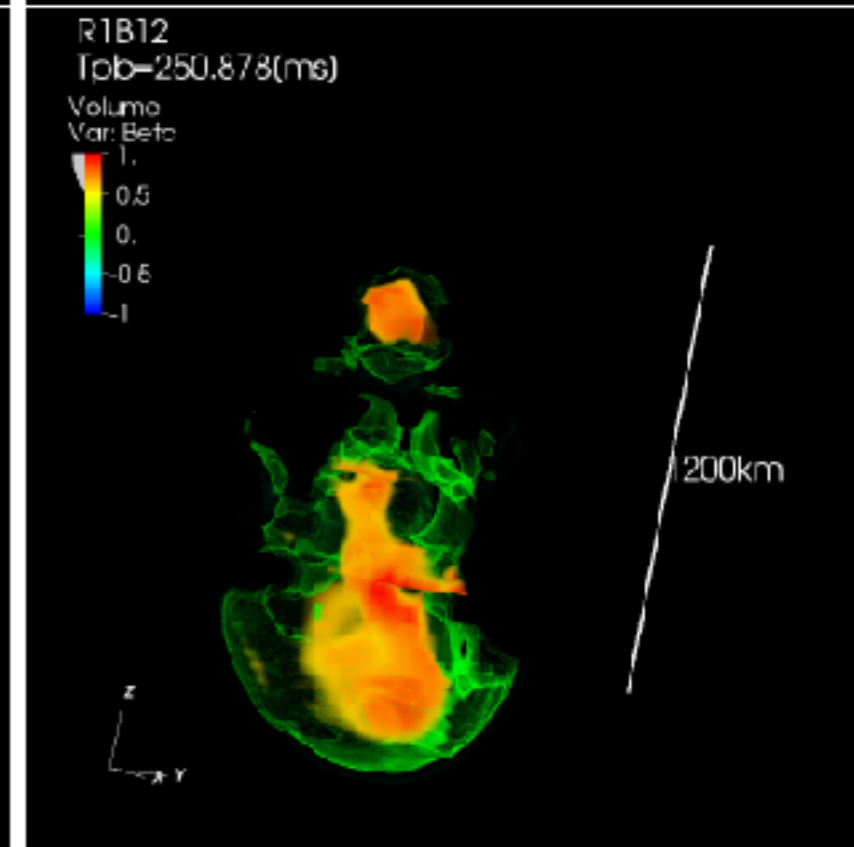
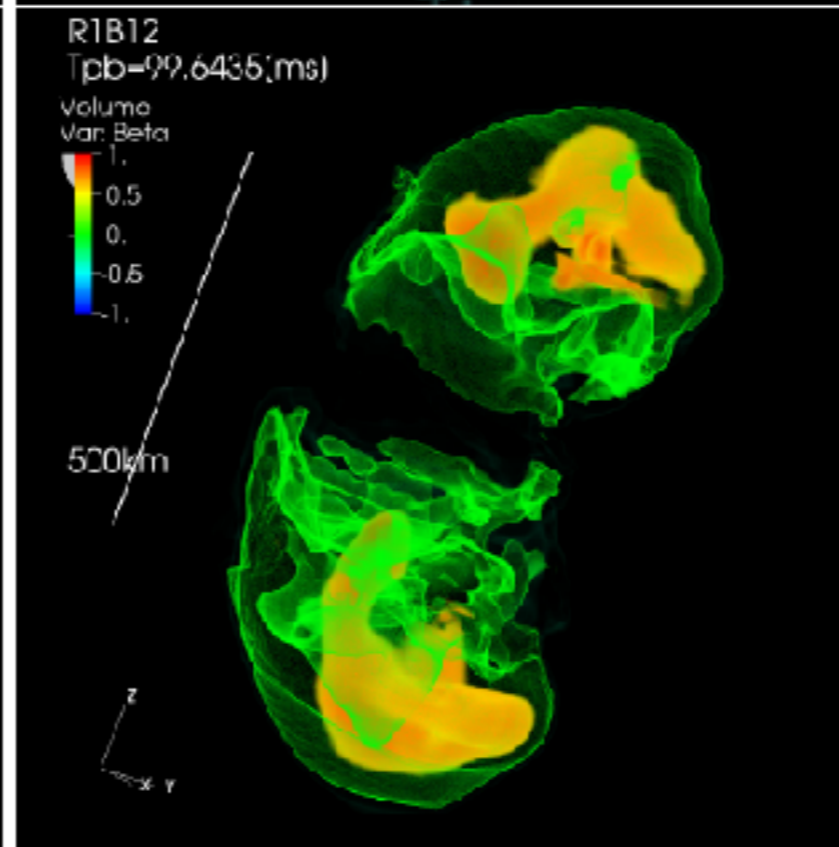
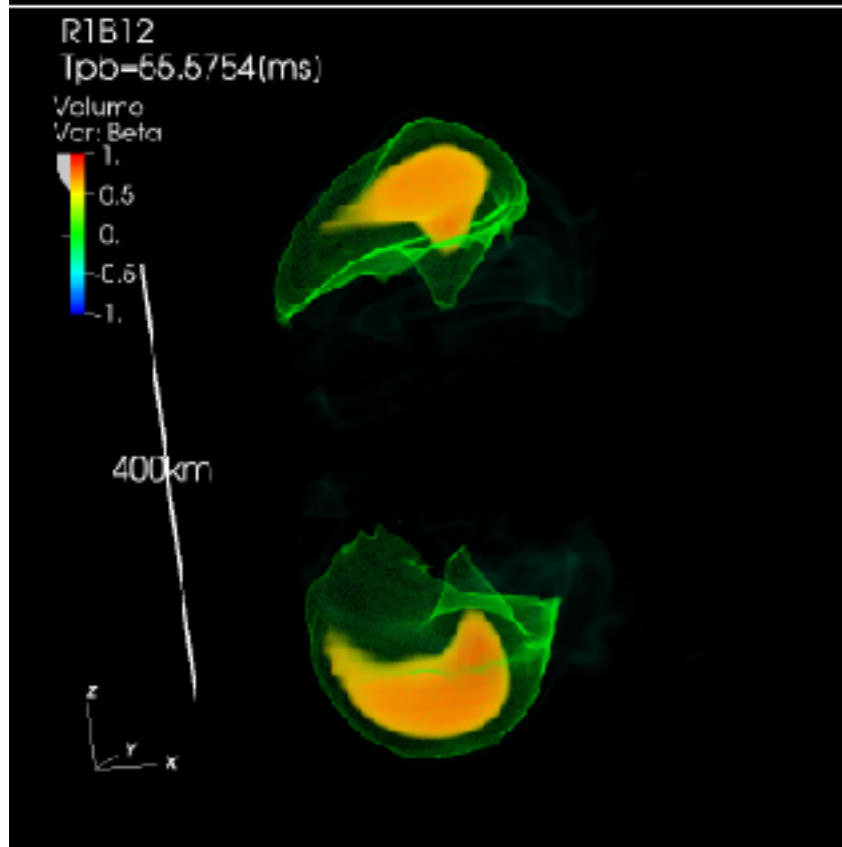
100ms

250ms

Entropy

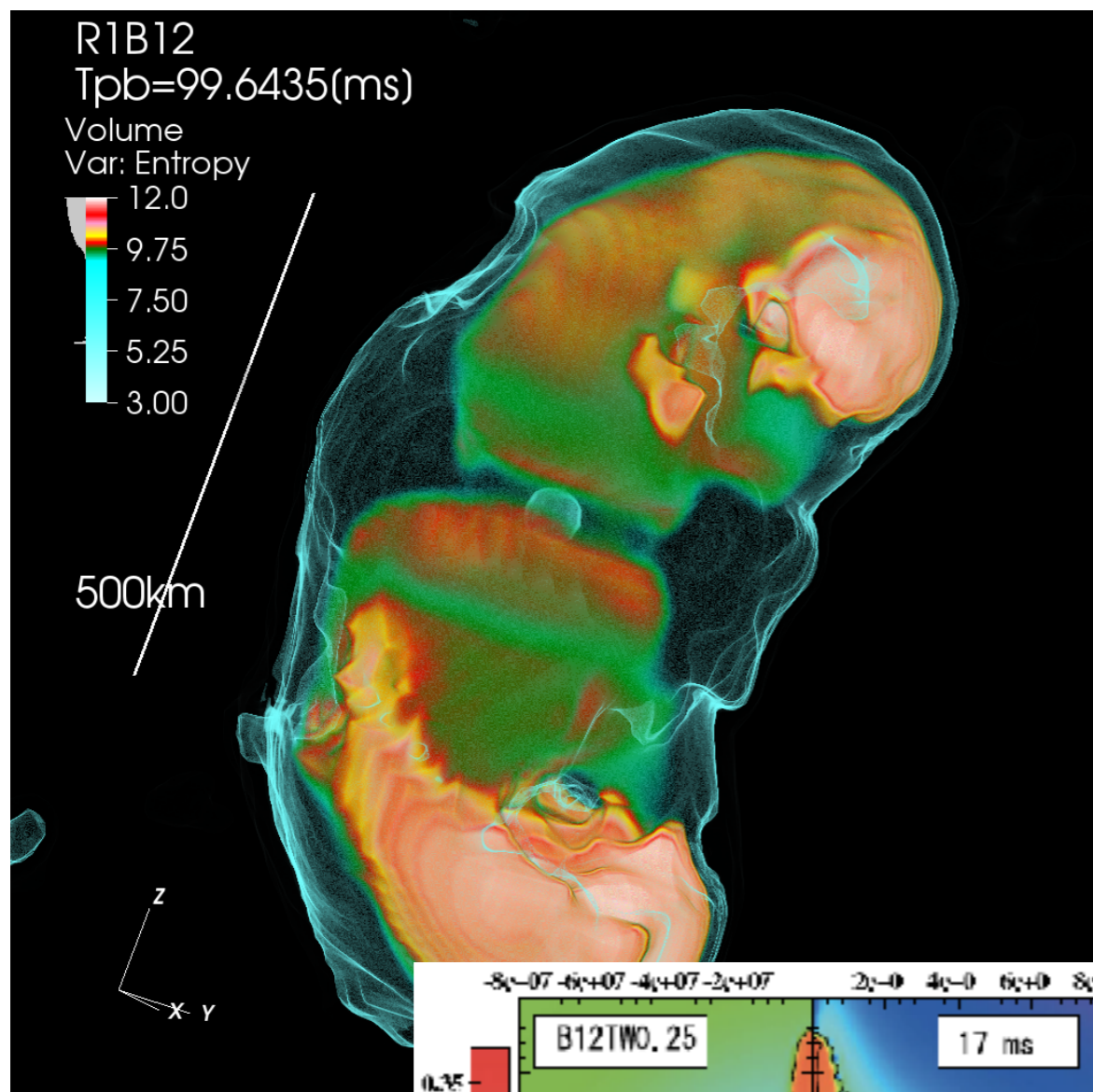


$\log(P_{\text{mag}}/P_{\text{gas}})$

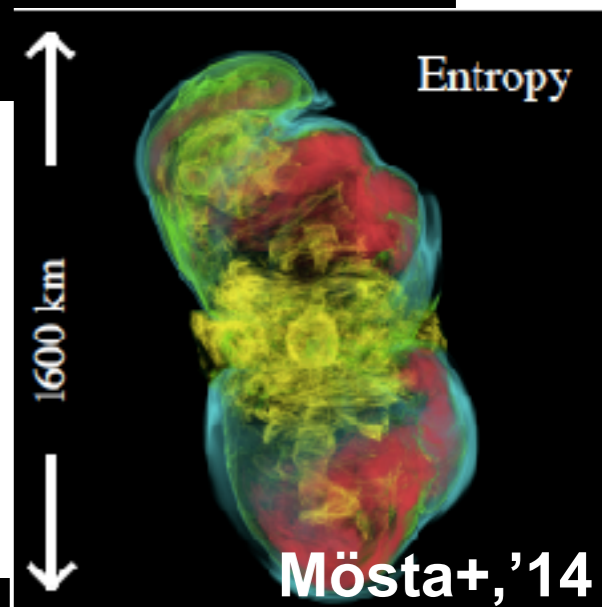
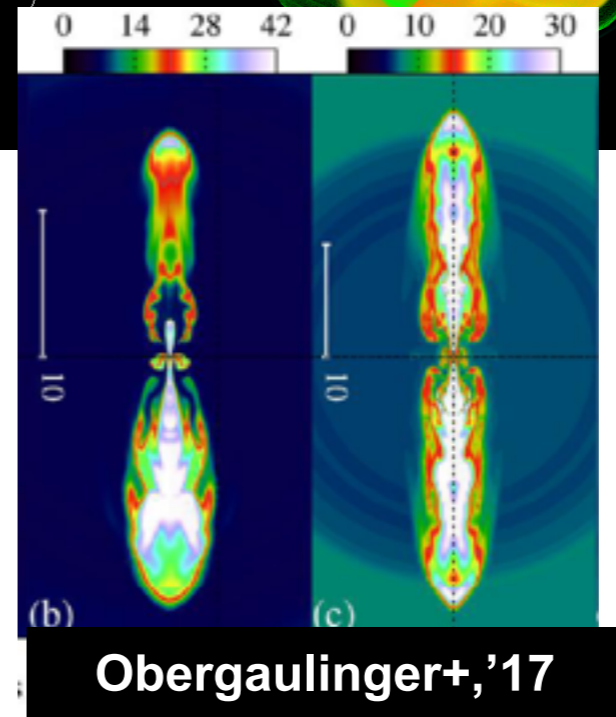
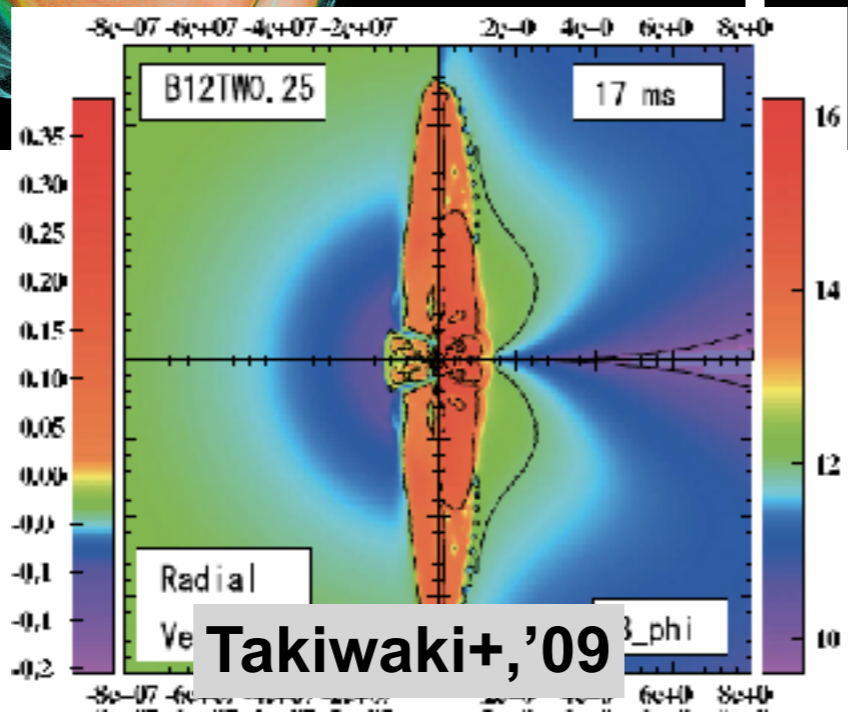
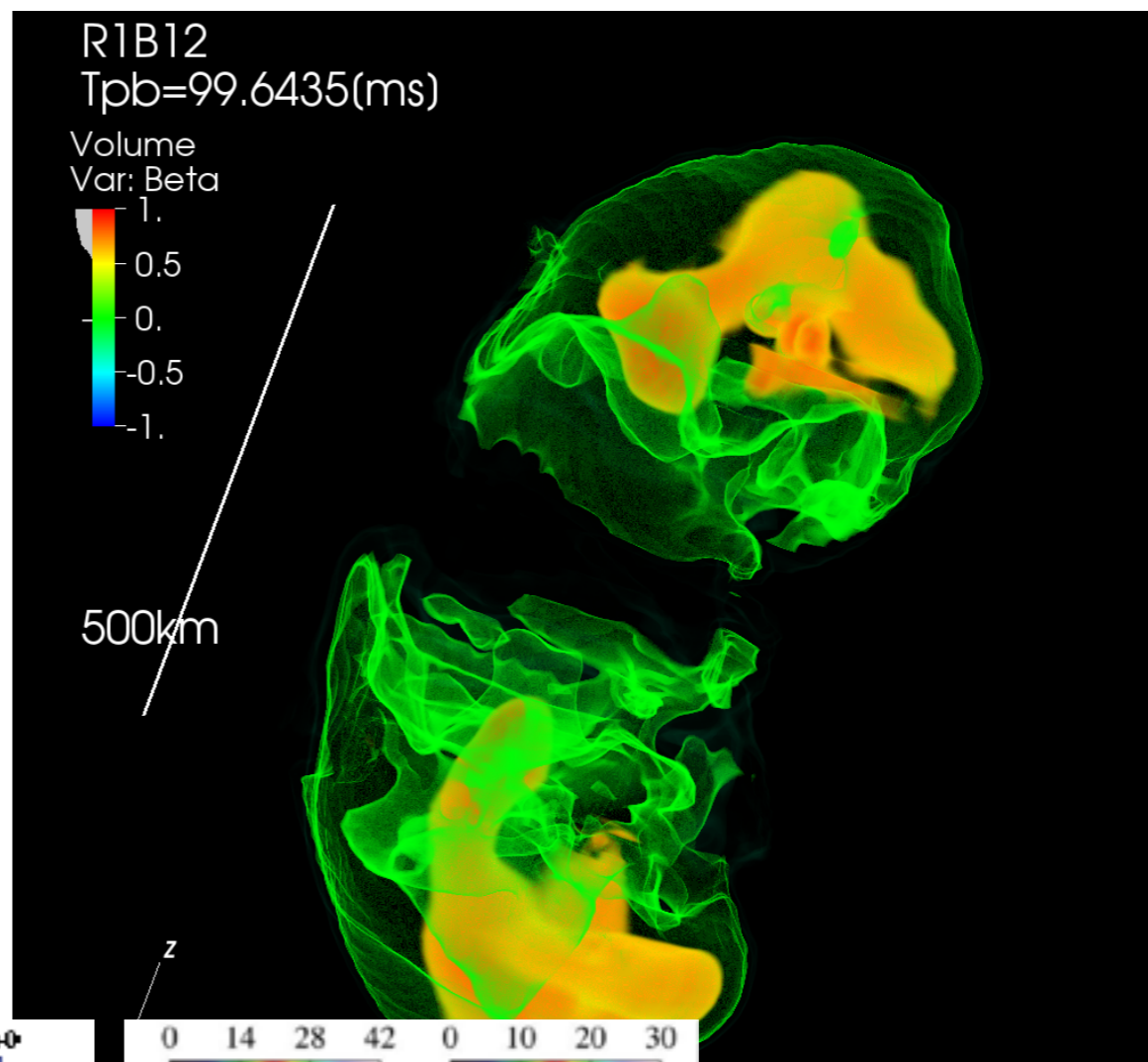


2D-3D comparison

Entropy

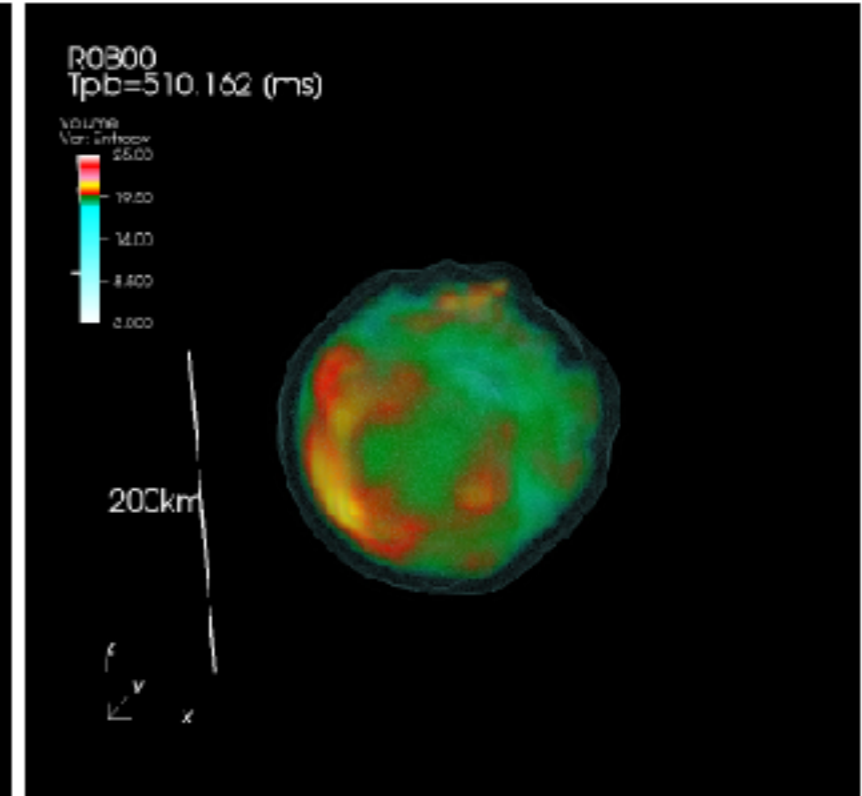
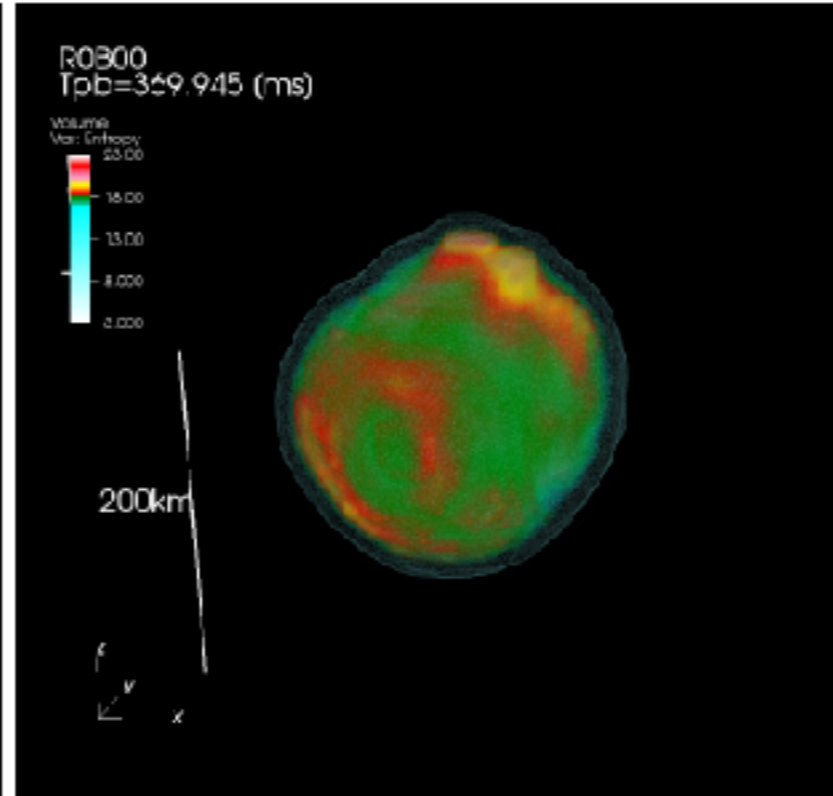
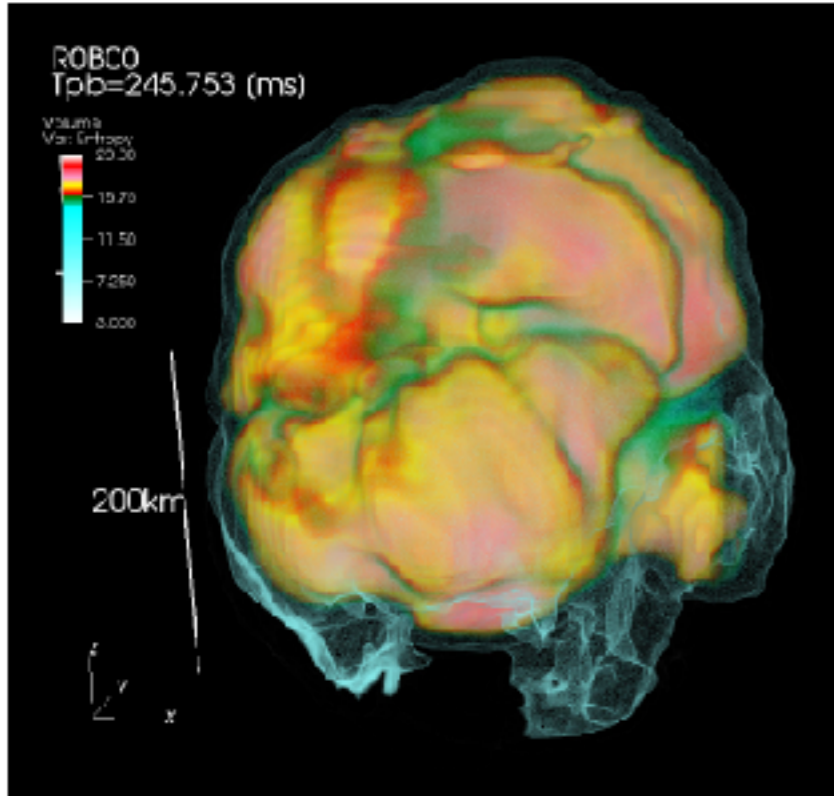


log(P_{mag}/P_{gas})

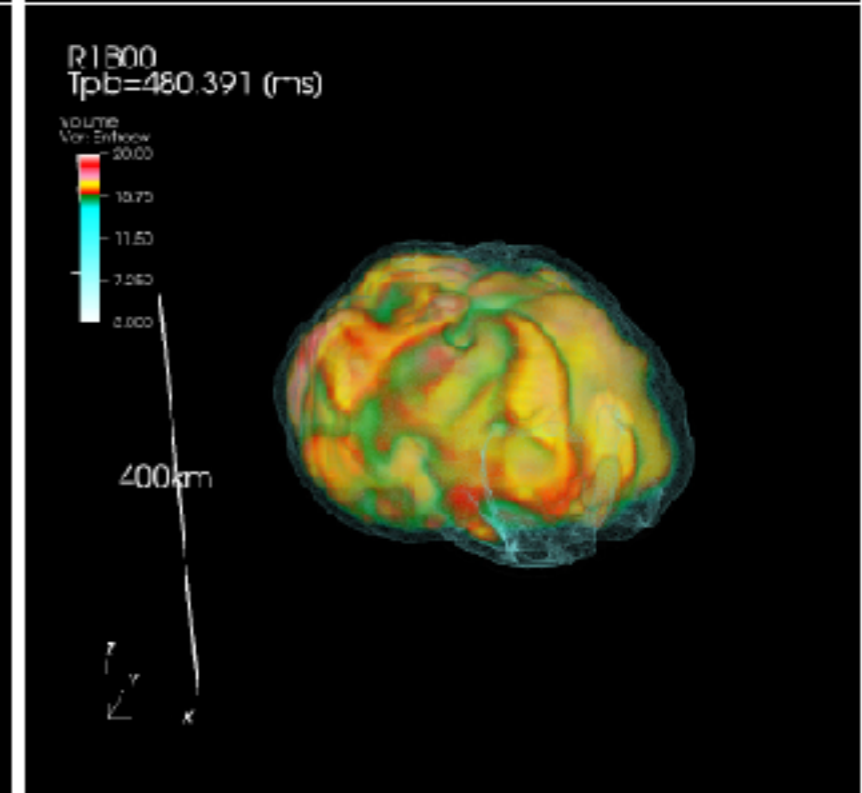
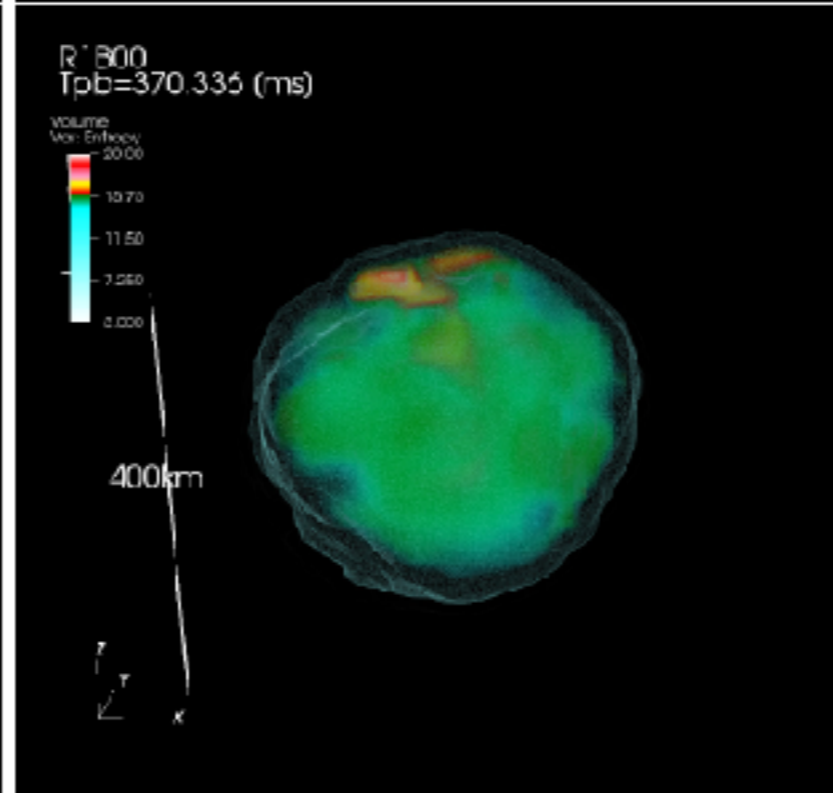
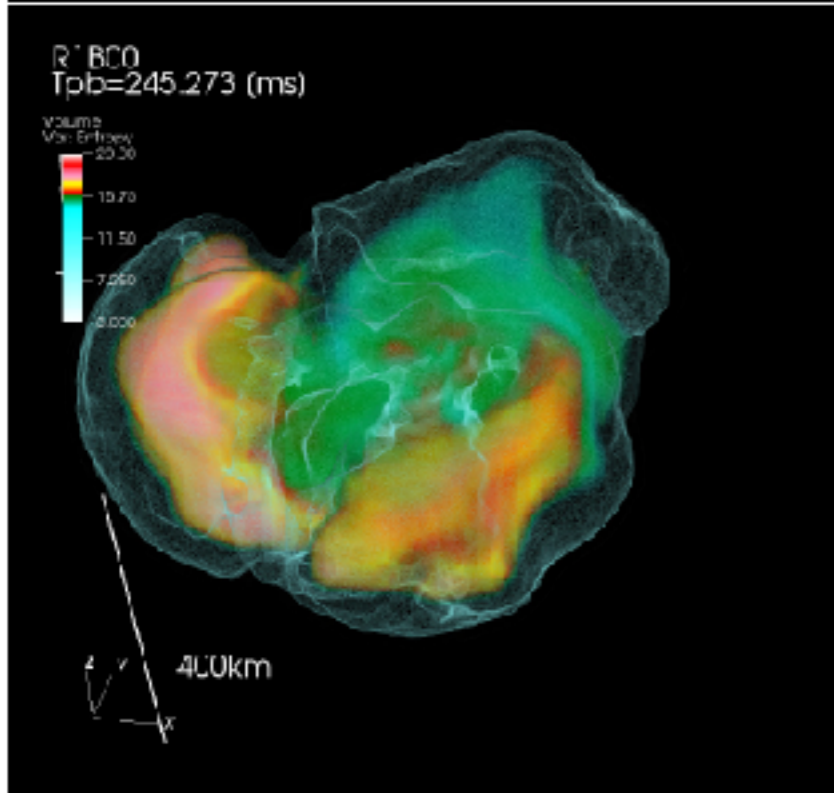


Non-magnetized models

Non-rotating model

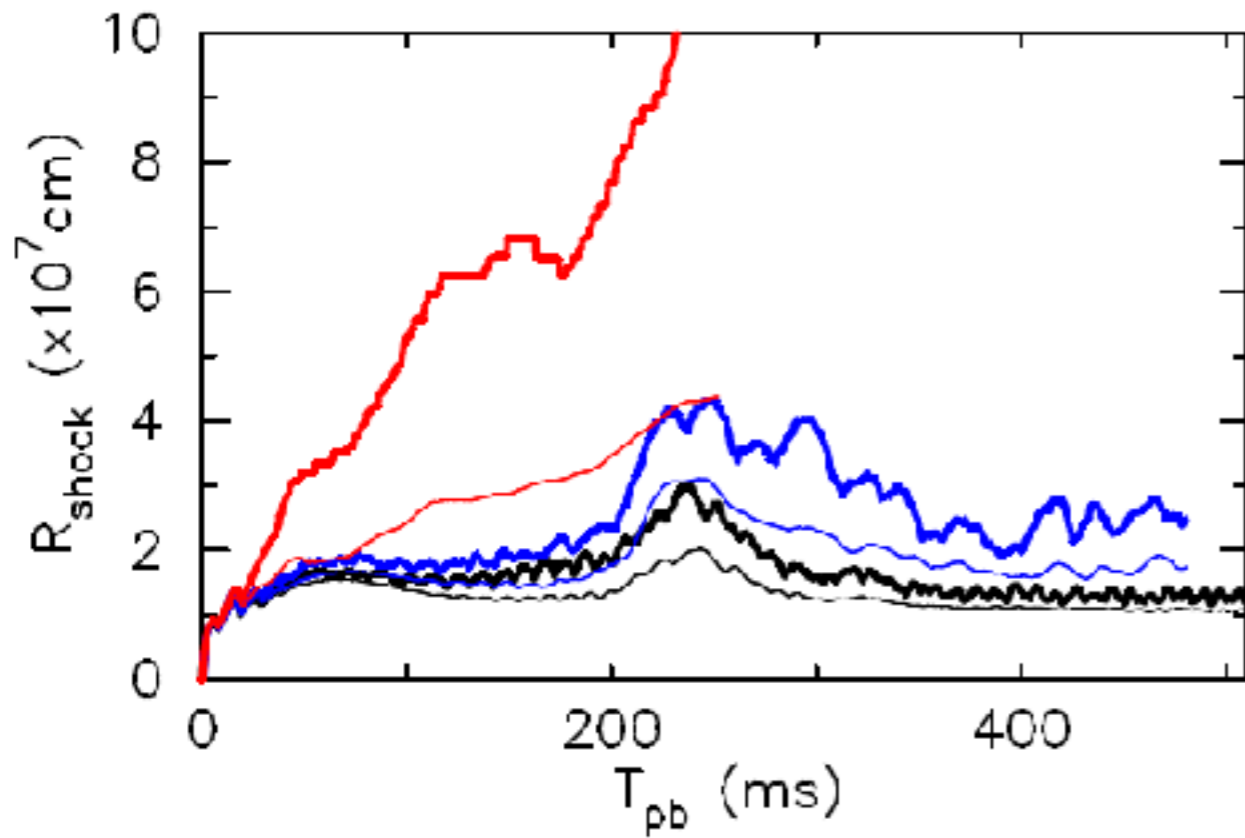


Rotating model

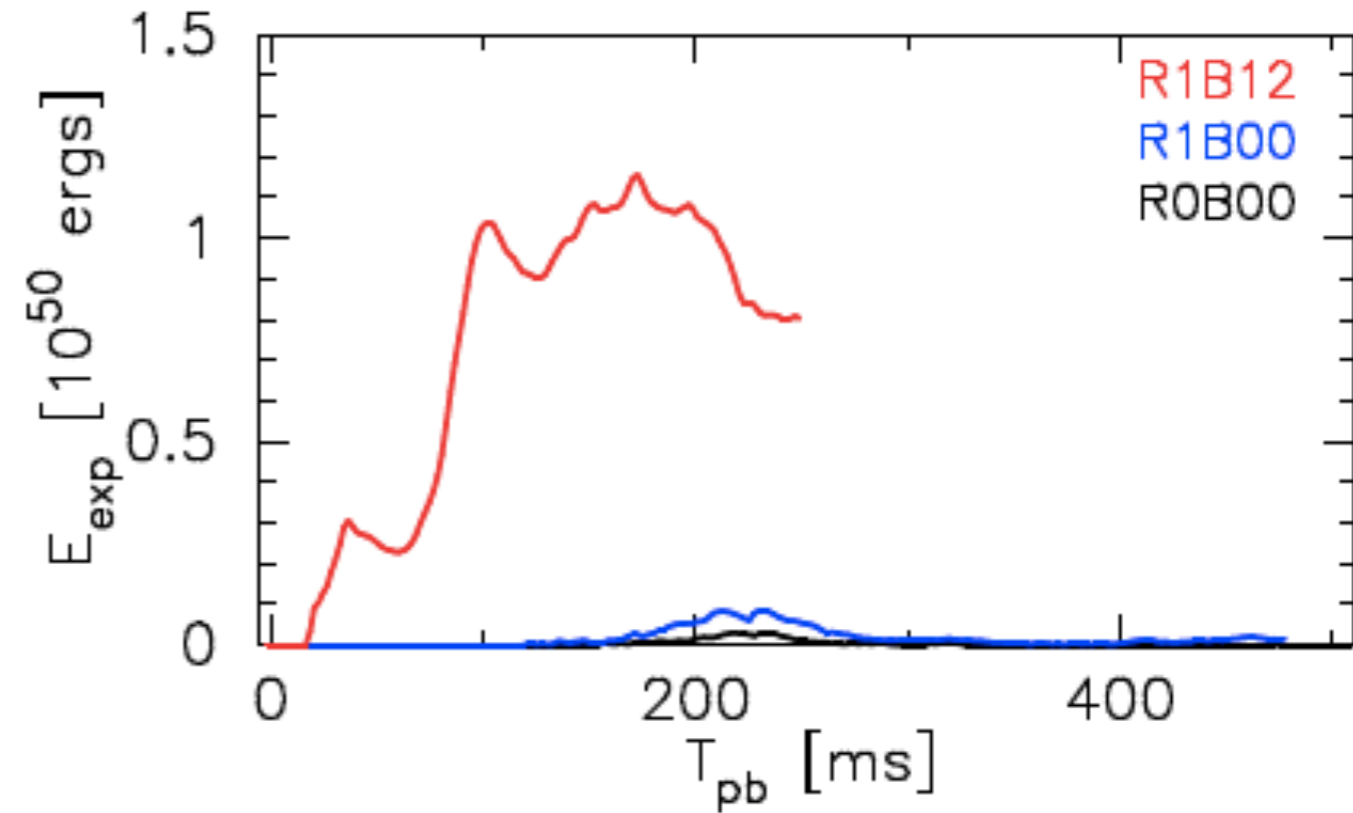


Energetics

Shock evolution

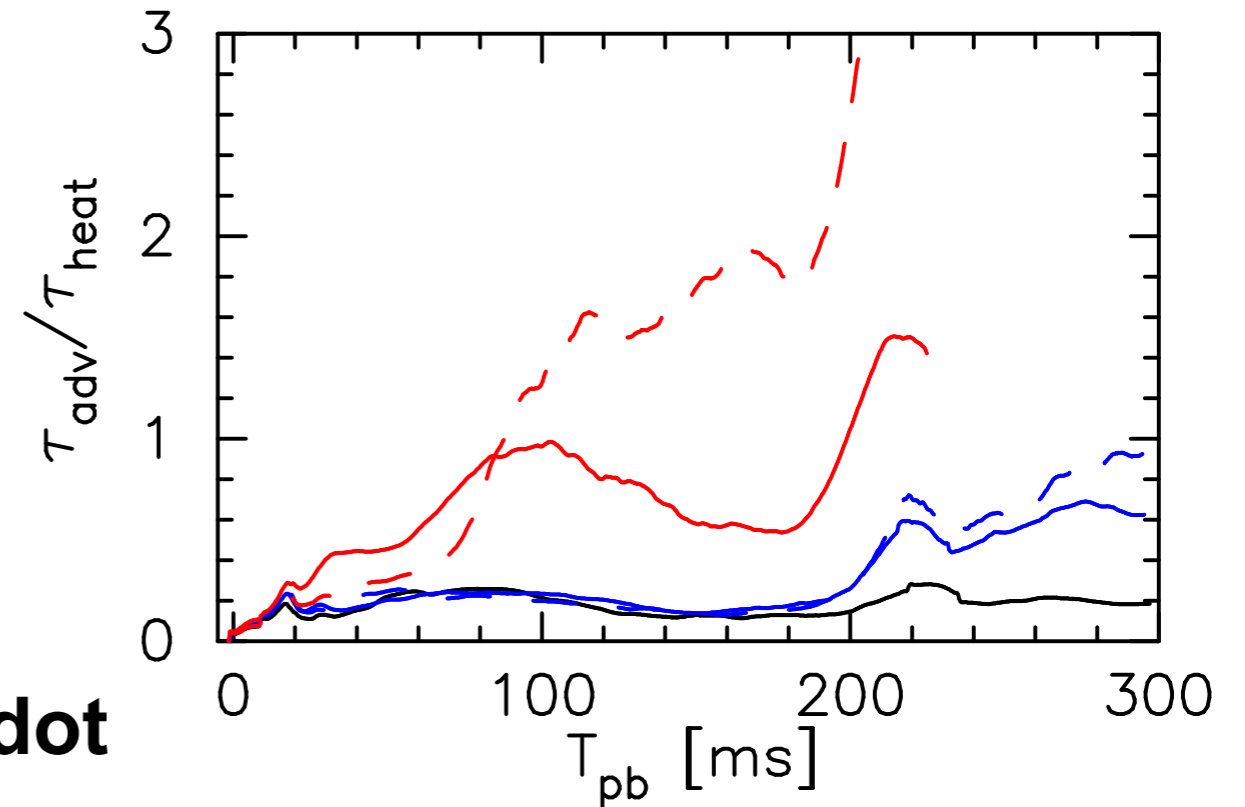
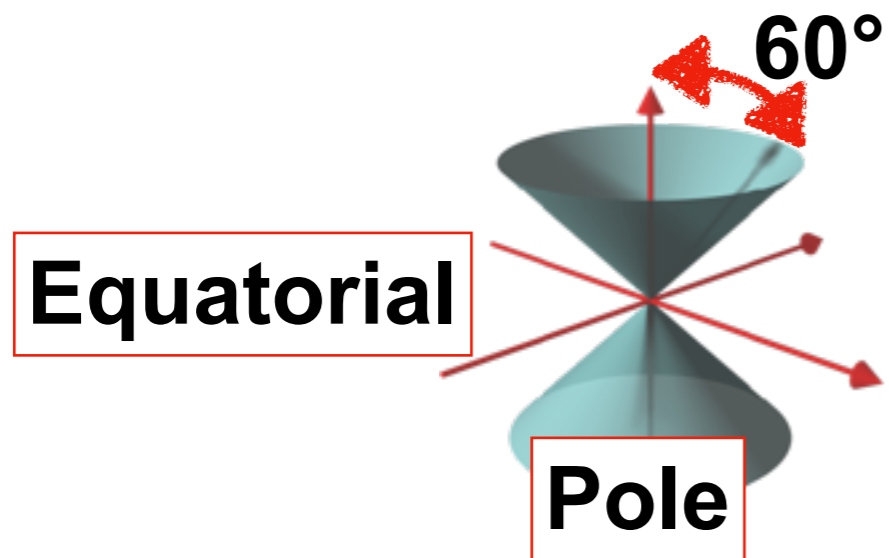
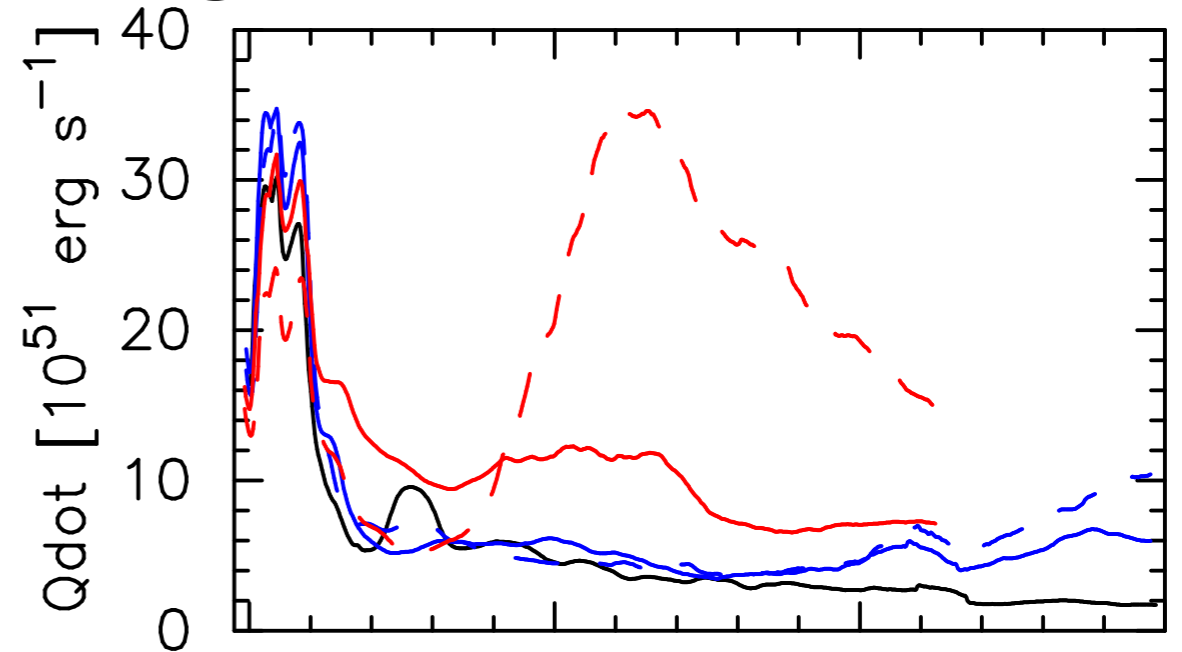
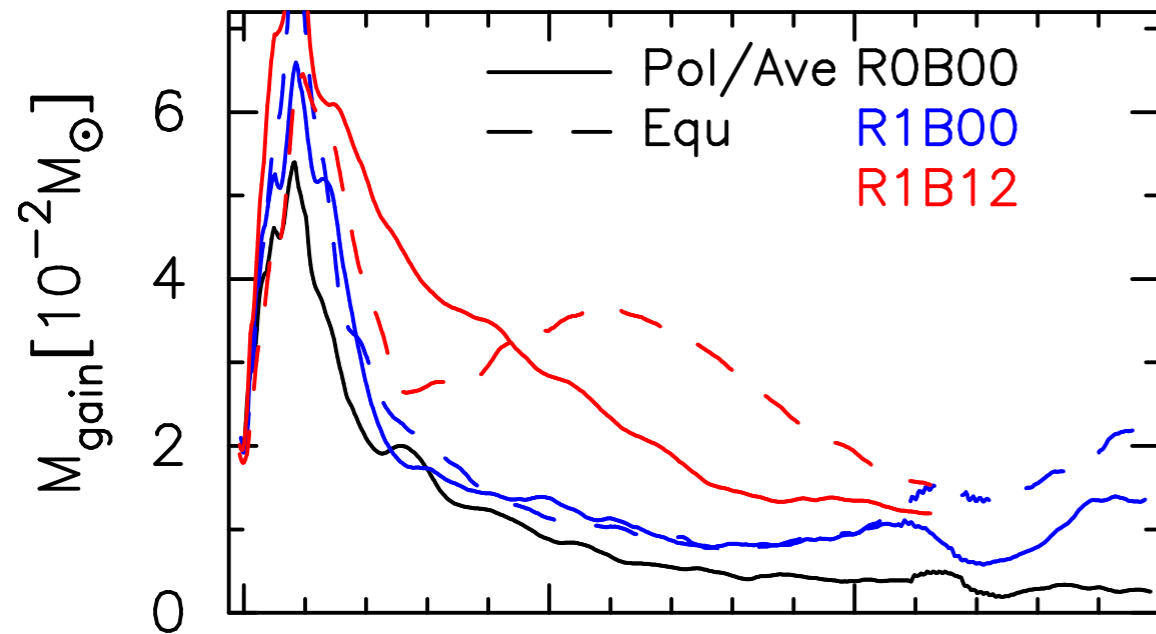


Diagnostic energy



Rotation and Magnetic fields facilitate the explosion

Neutrino heated? or magneto-driven?



- Rotation & B increase M_{gain} and \dot{Q}
- For MHD model

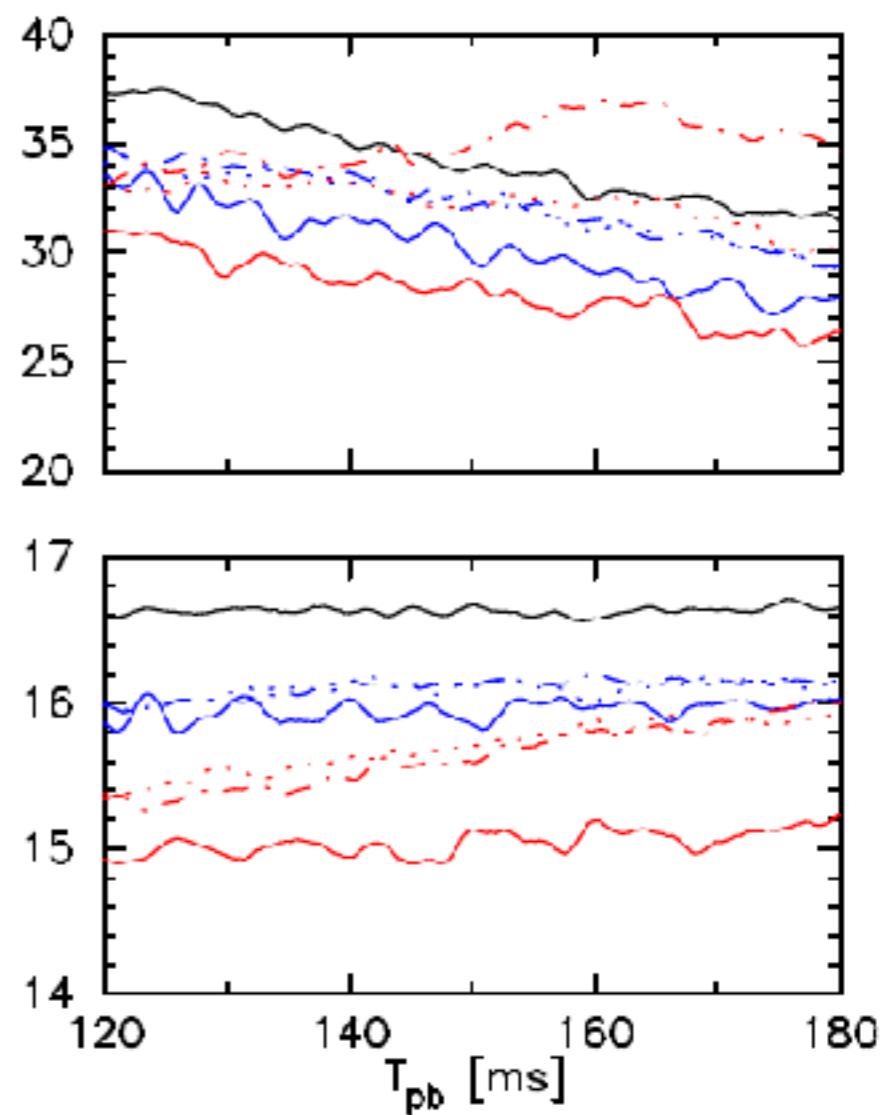
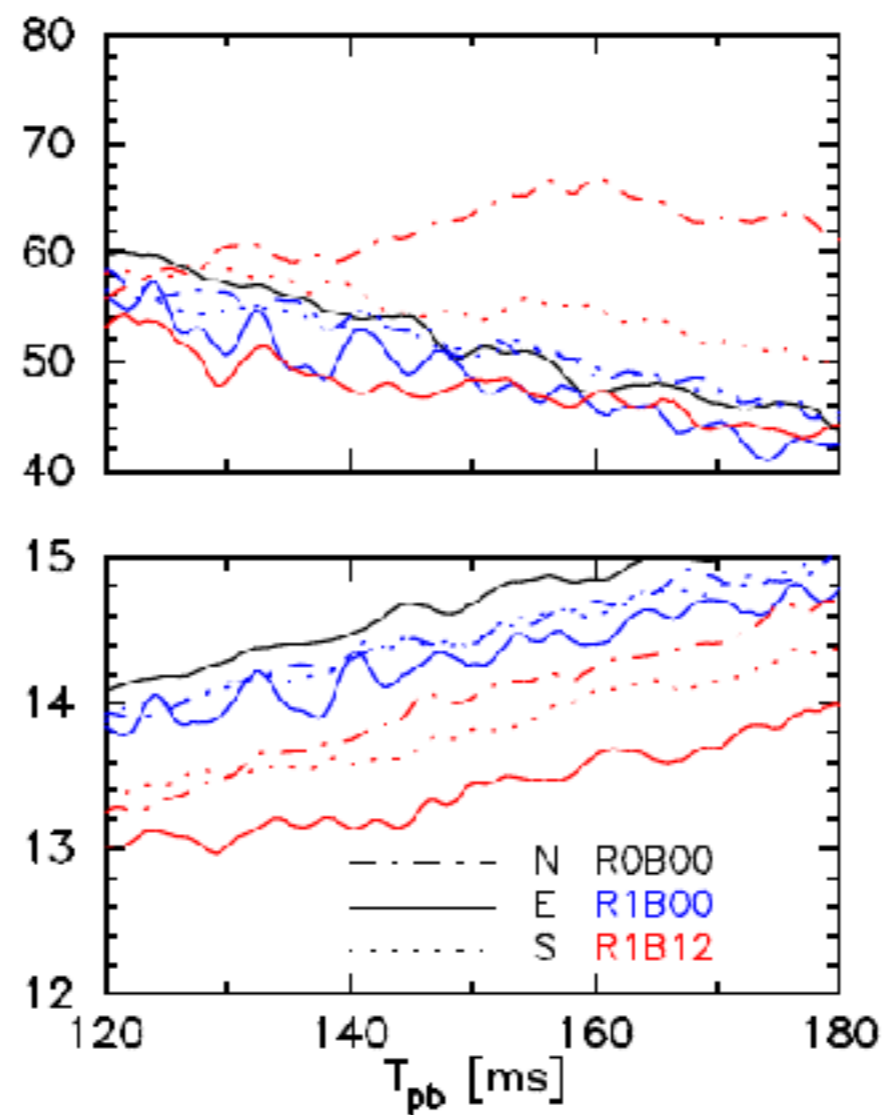
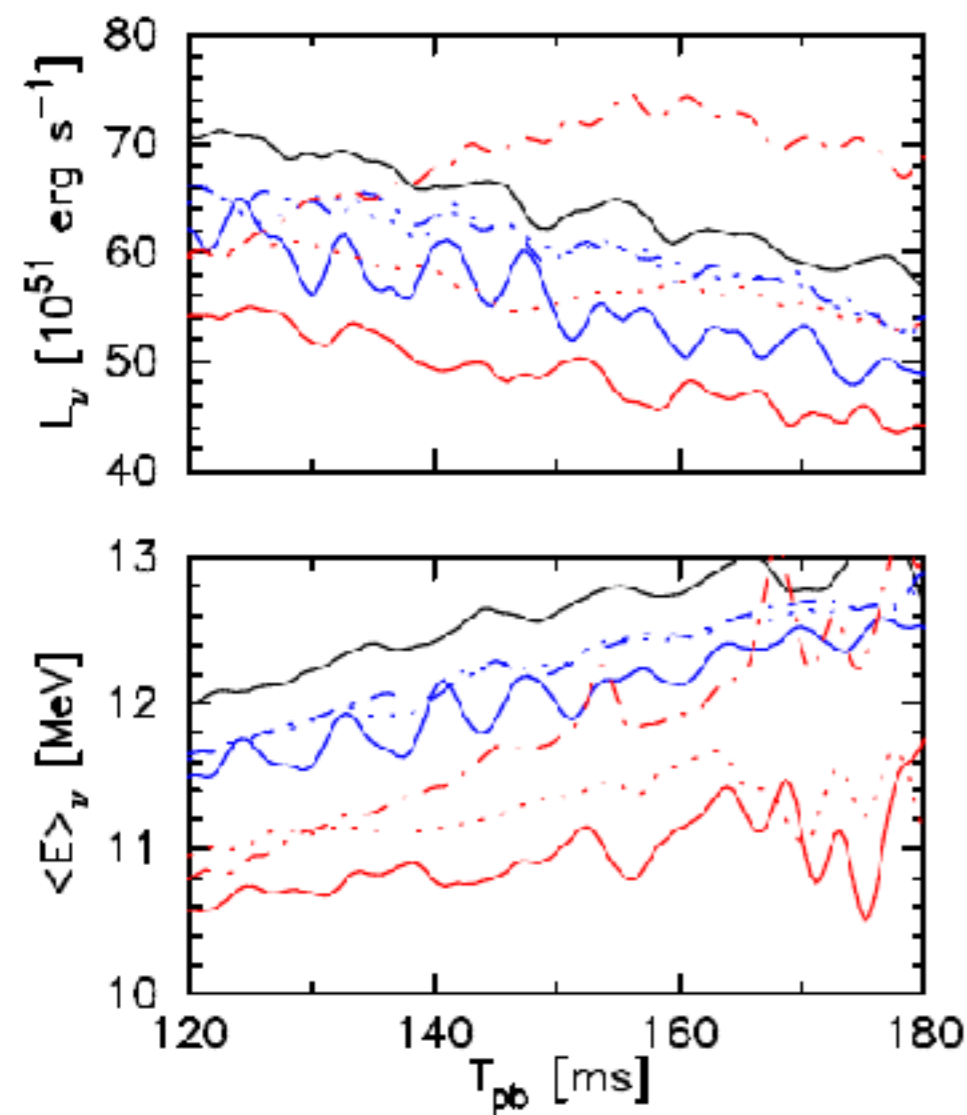
- Equatorial expansion is supported by v-heating $\tau_{\text{adv}}/\tau_{\text{heat}} > 1$
- Prompt bipolar outflow is due to magnetic field $\tau_{\text{adv}}/\tau_{\text{heat}} < 1$
- Later by v-heating $\tau_{\text{adv}}/\tau_{\text{heat}} > 1$

Neutrino emission (angle dependence)

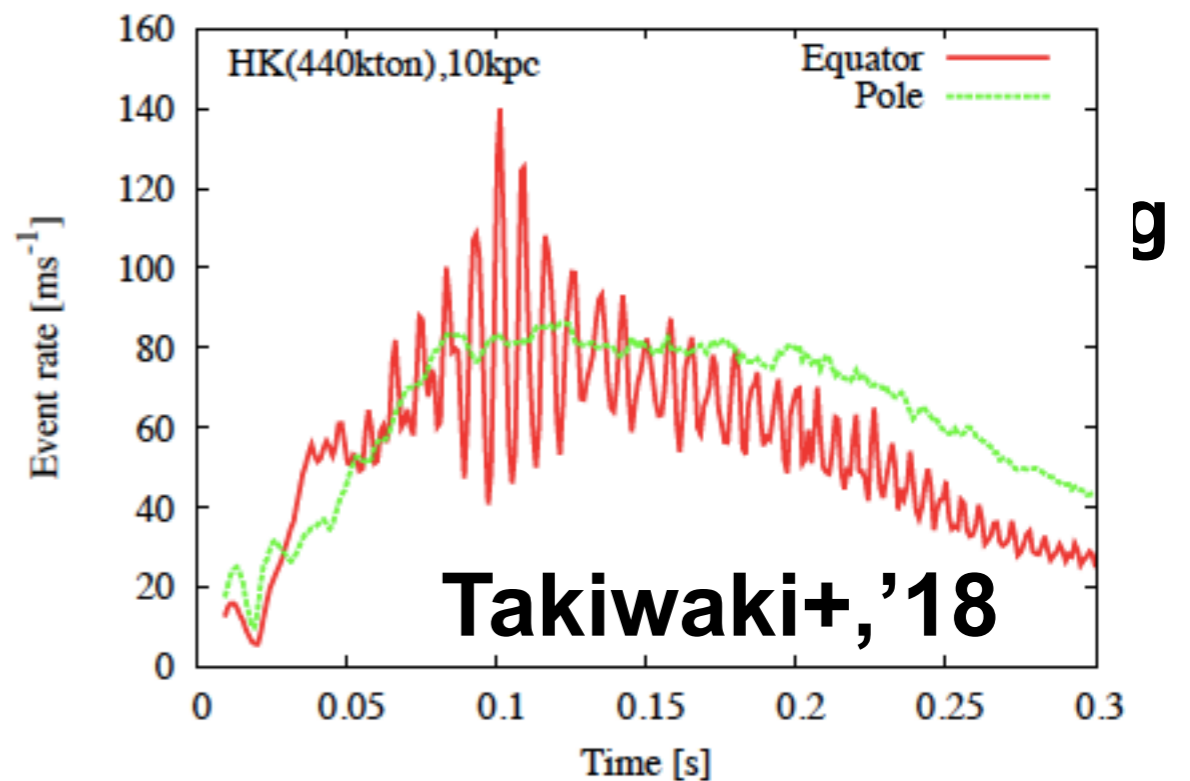
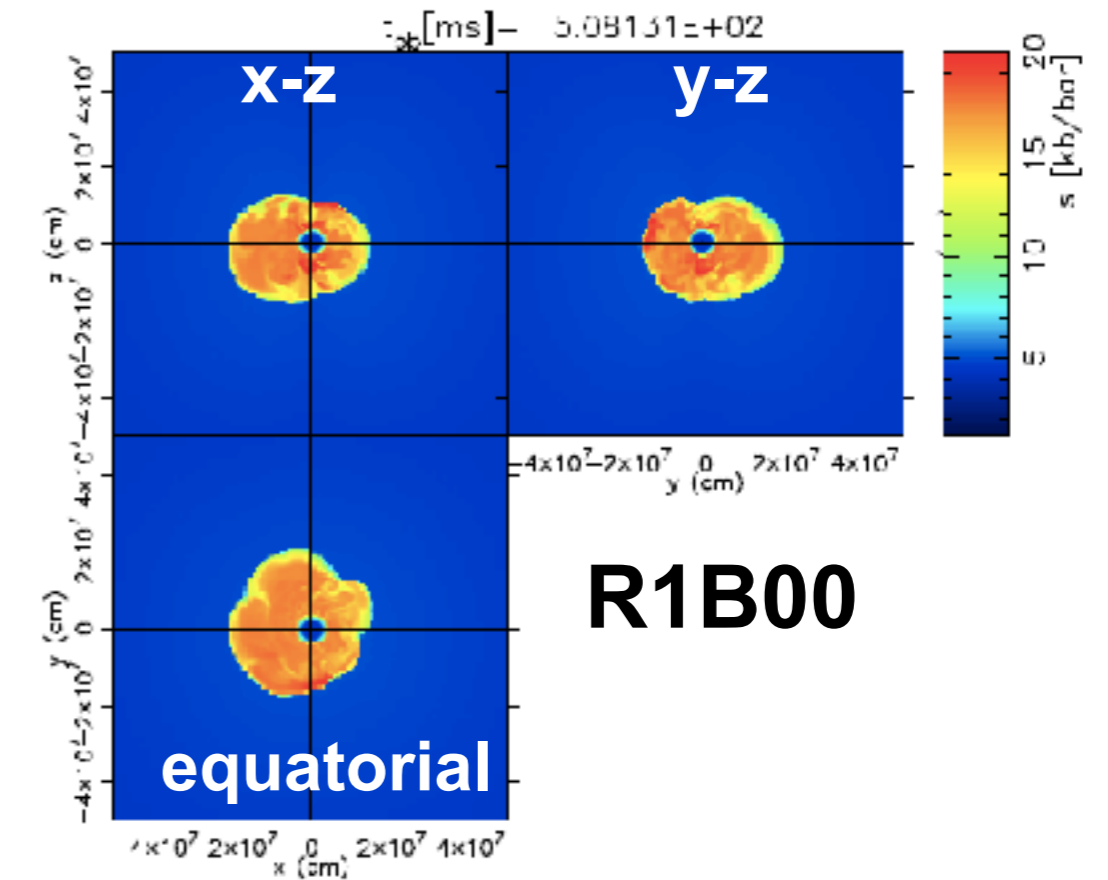
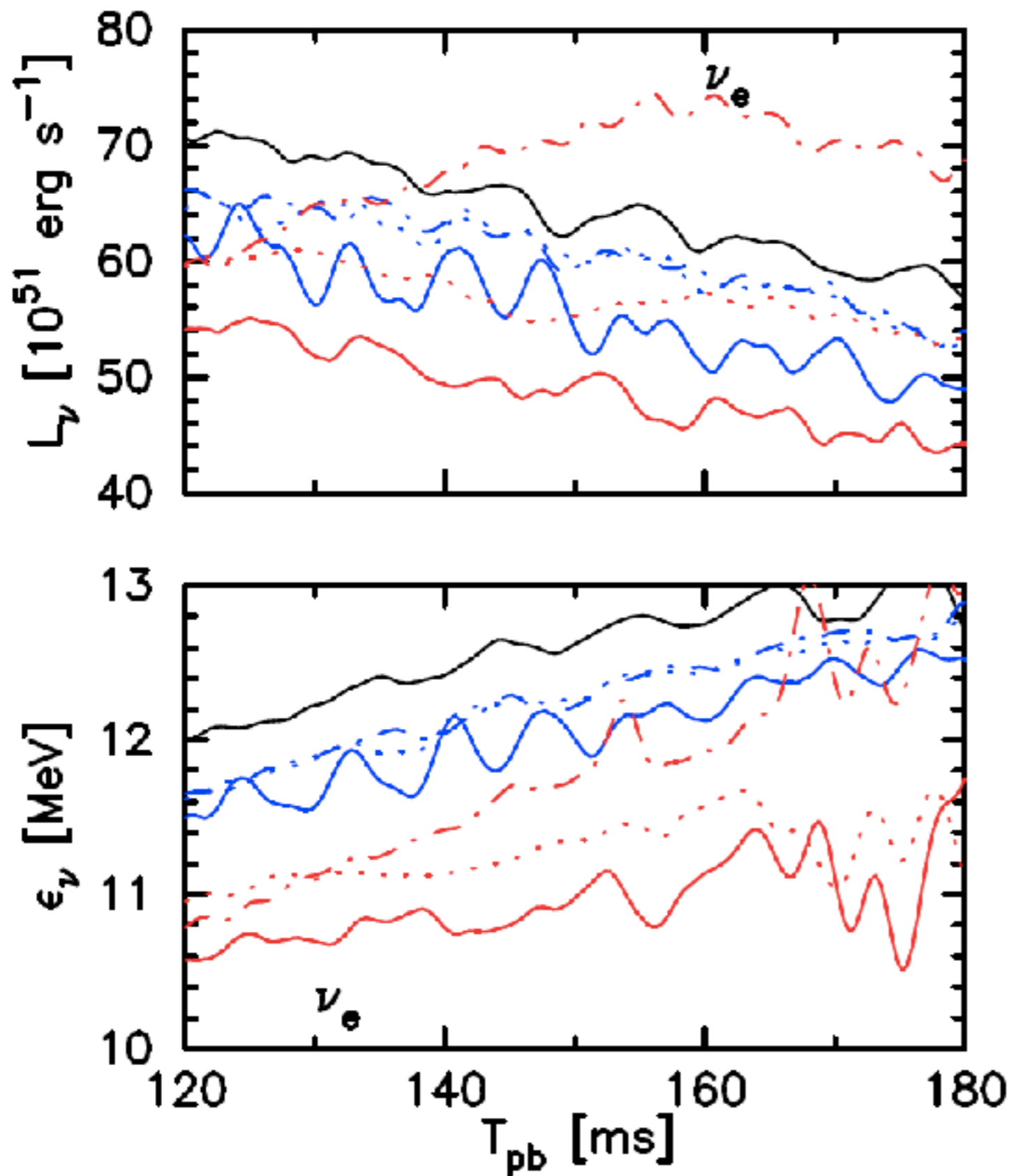
ν_e

$\bar{\nu}_e$

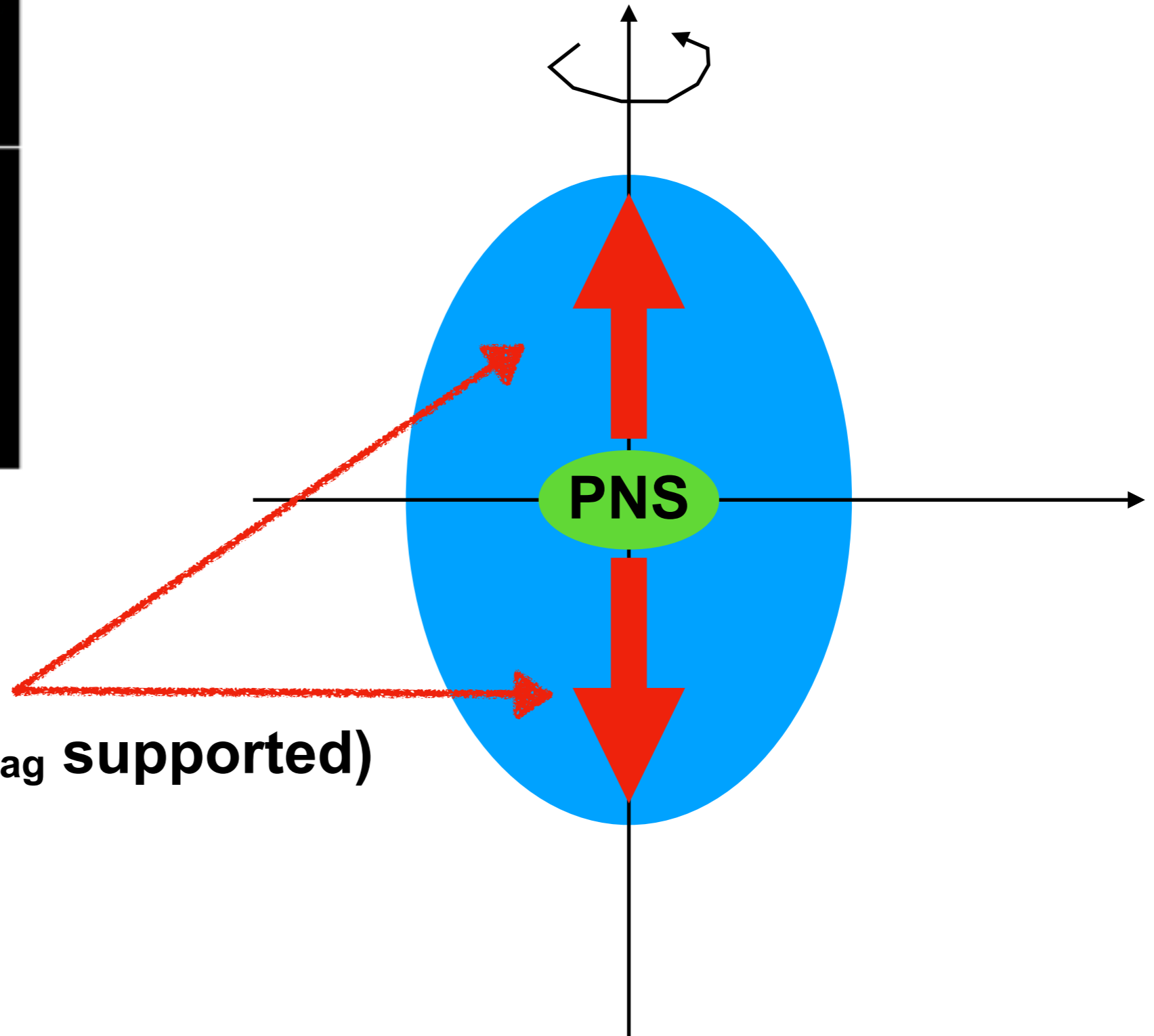
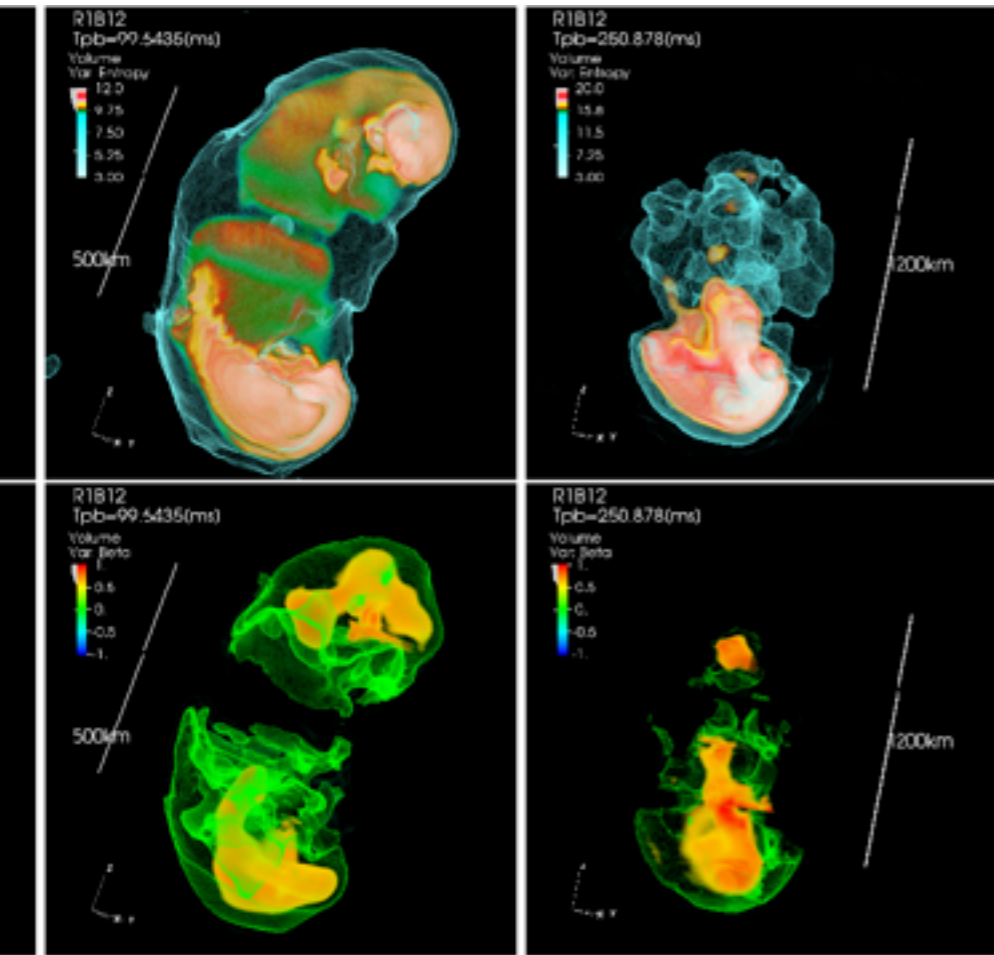
ν_x



Neutrino emission (angle dependence)

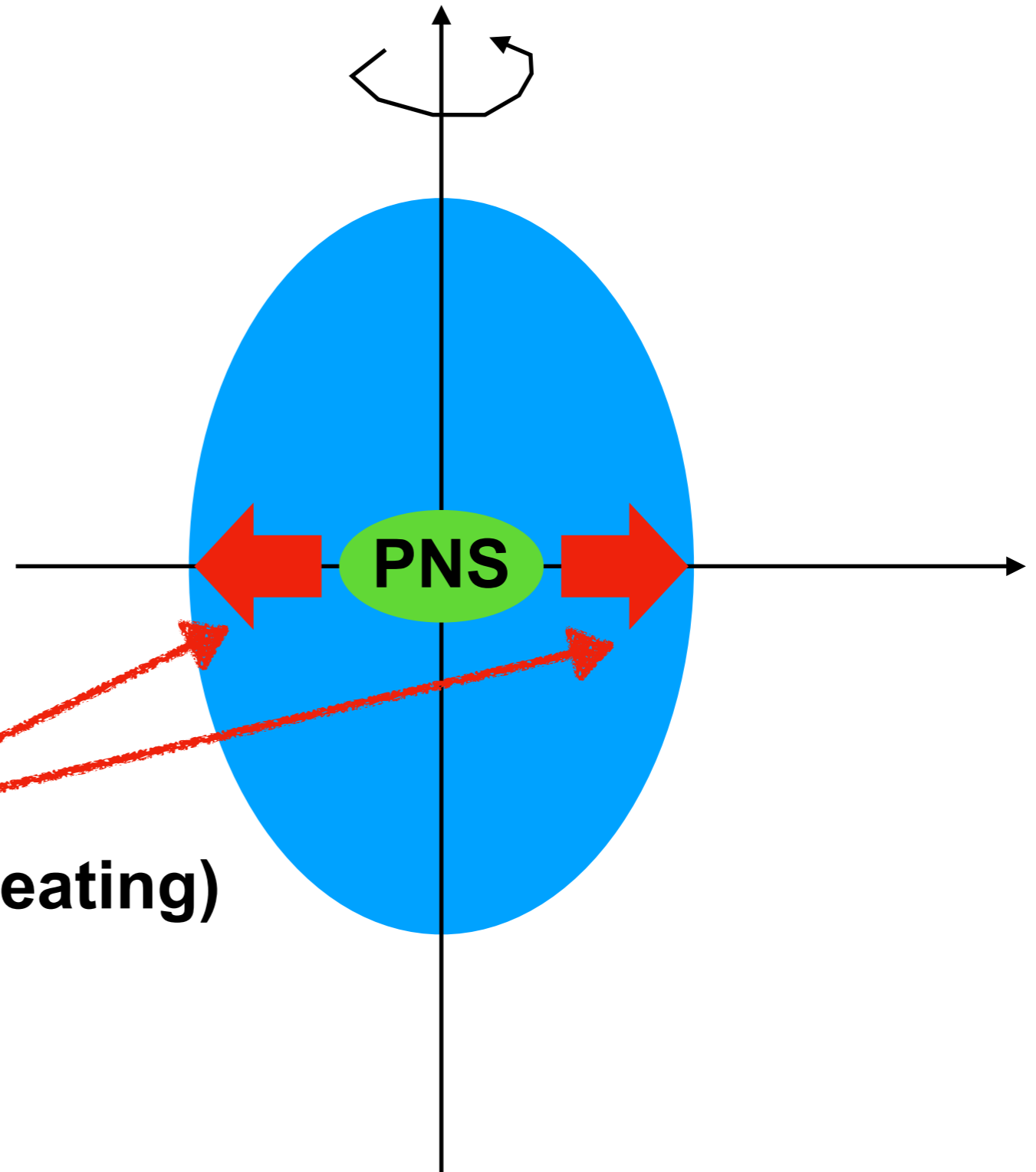
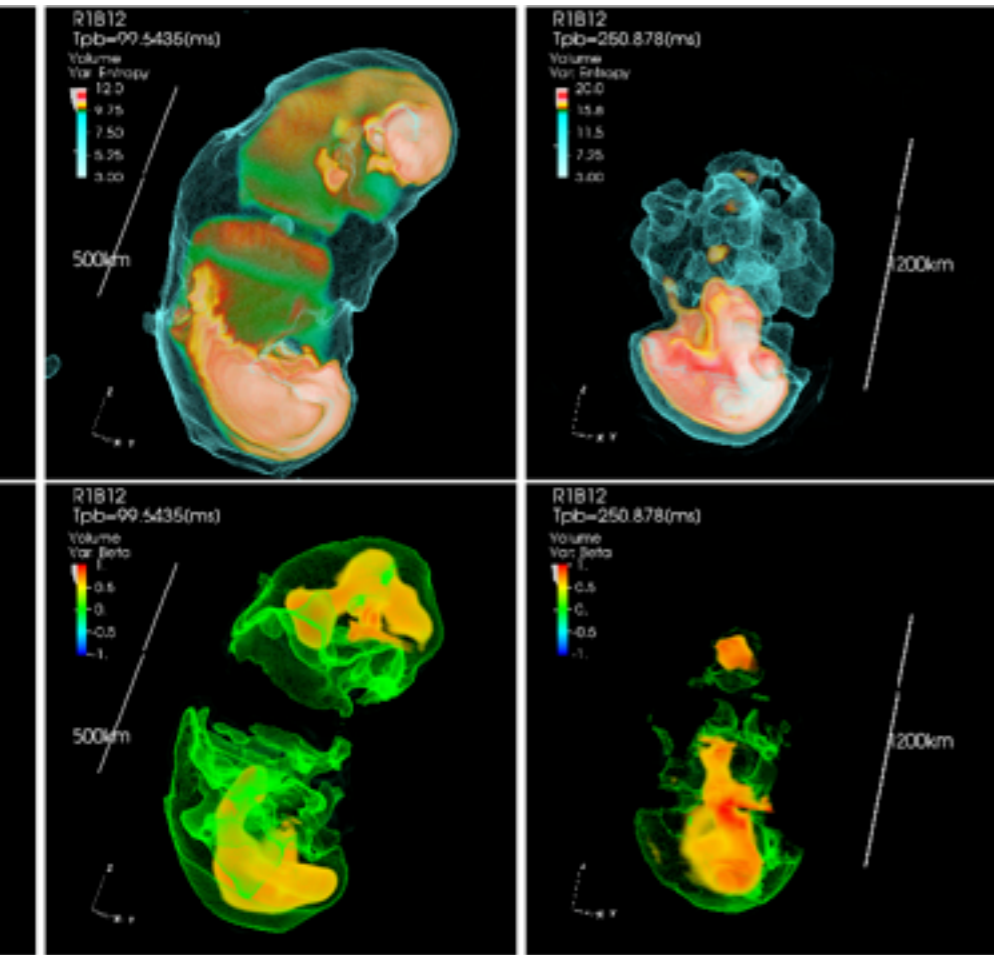


Ejecta structure



low Y_e (P_{mag} supported)

Ejecta structure



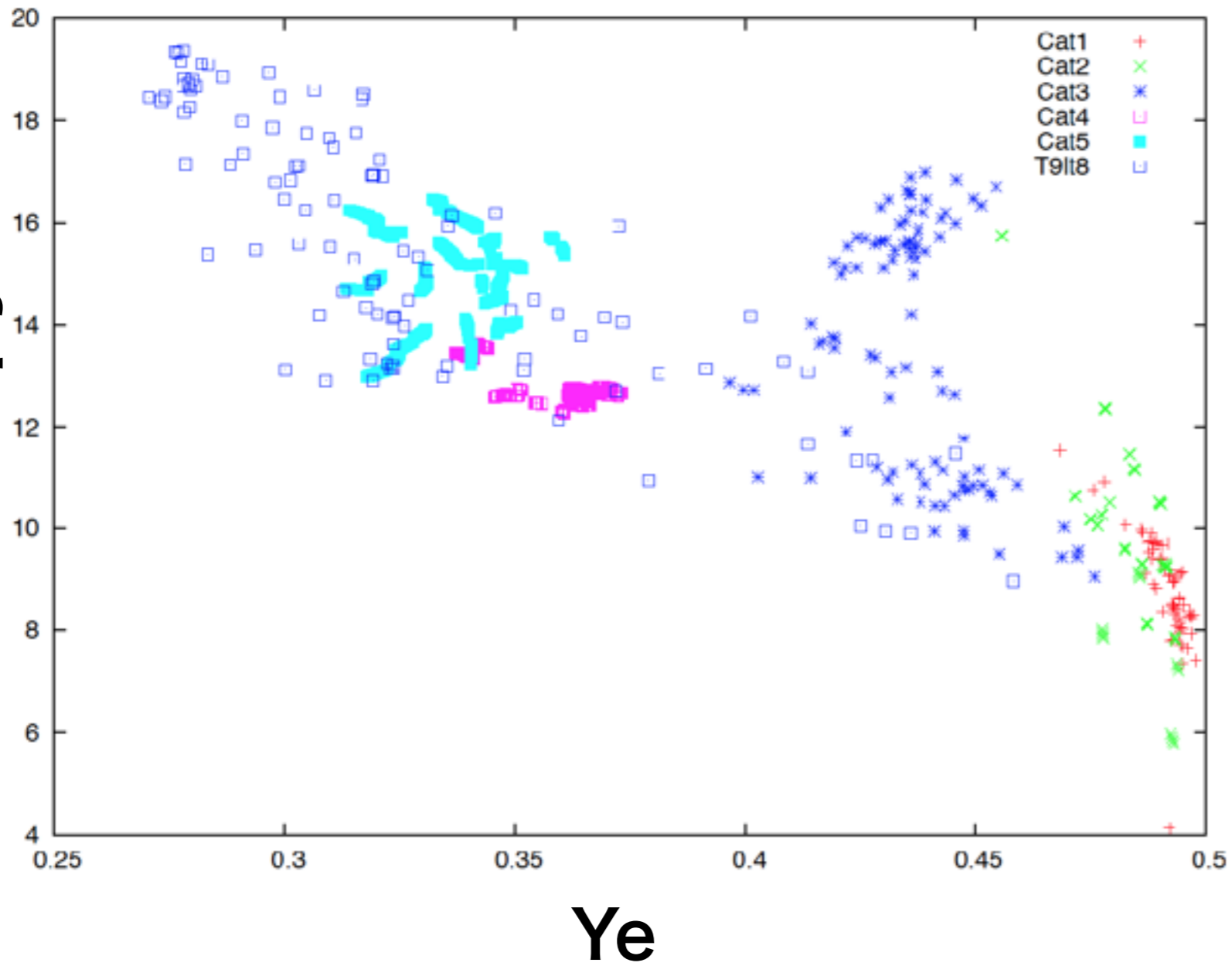
high Y_e (neutrino heating)



Selection rule

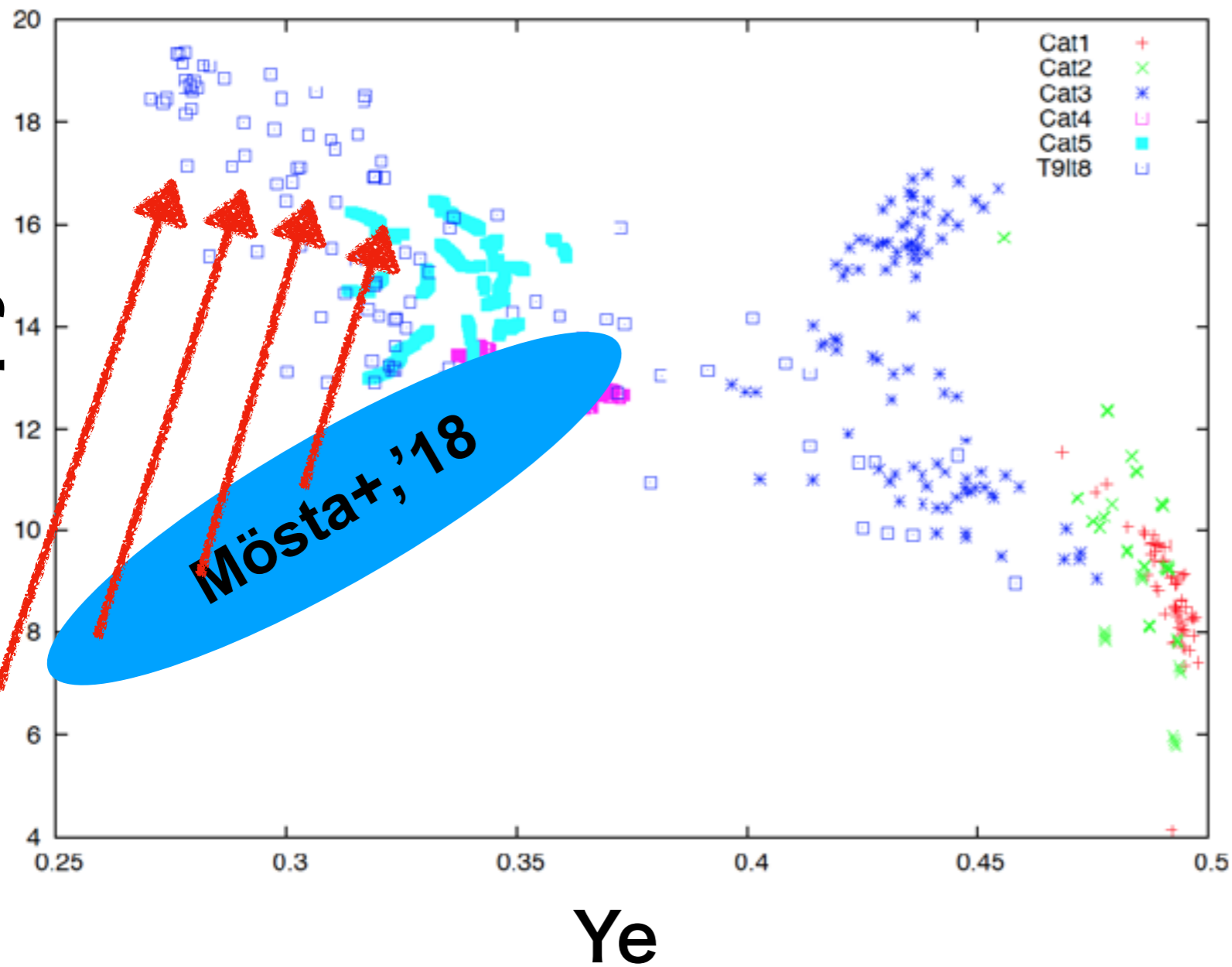
- (1) Y_e and entropy unchanged (and low peak temperature), such that the progenitor composition does not change much
- (2) Y_e unchanged, but high peak temperature, with explosive nucleosynthesis
- (3) Y_e once <0.45 and at the end >0.38
- (4) Y_e always <0.45 and entropy <15
- (5) Y_e always <0.45 and entropy >15
- (6) T_9 at t_8 : final temperature (averaged in time of ~ 10 ms) decreases below 8GK.

Entropy



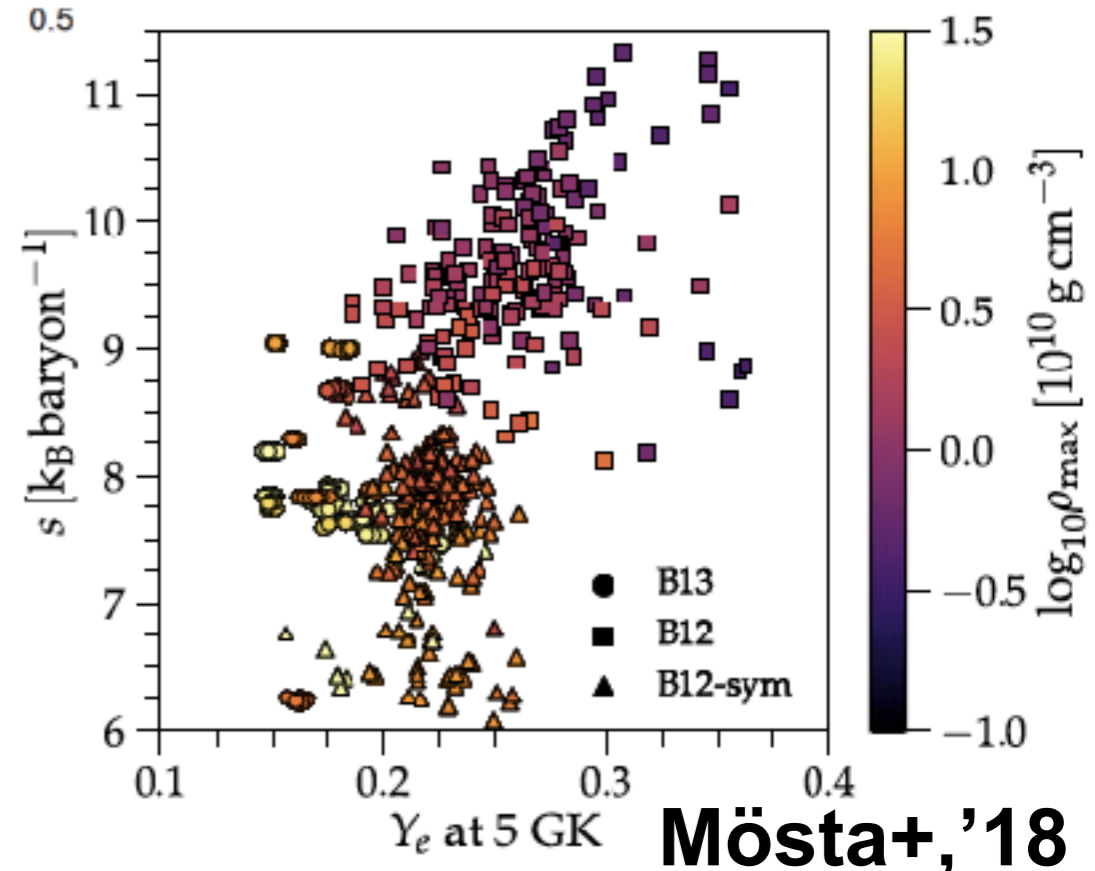
Scattering in S- Y_e plane

Entropy

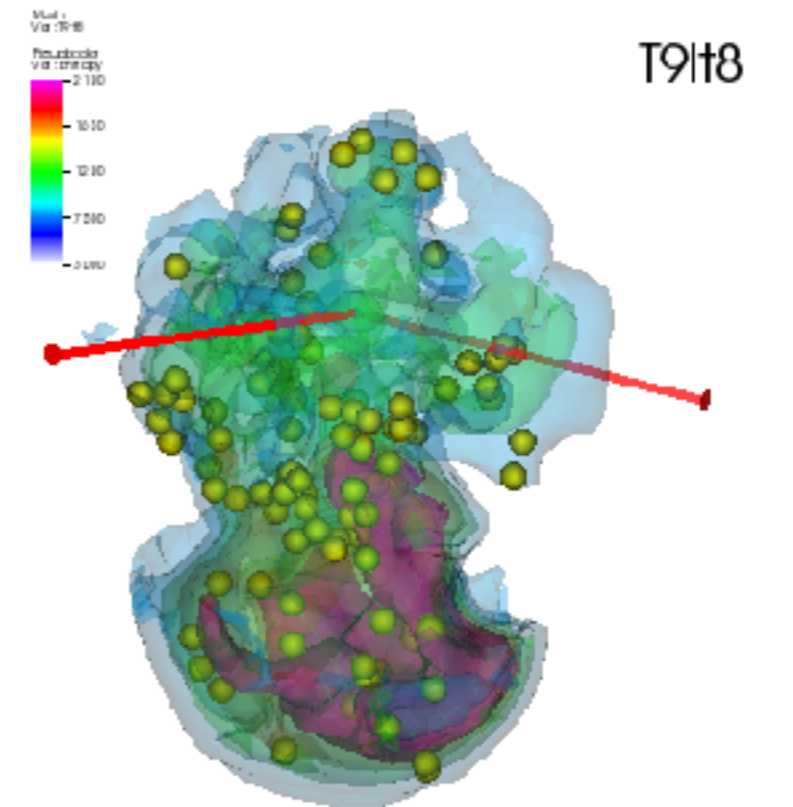
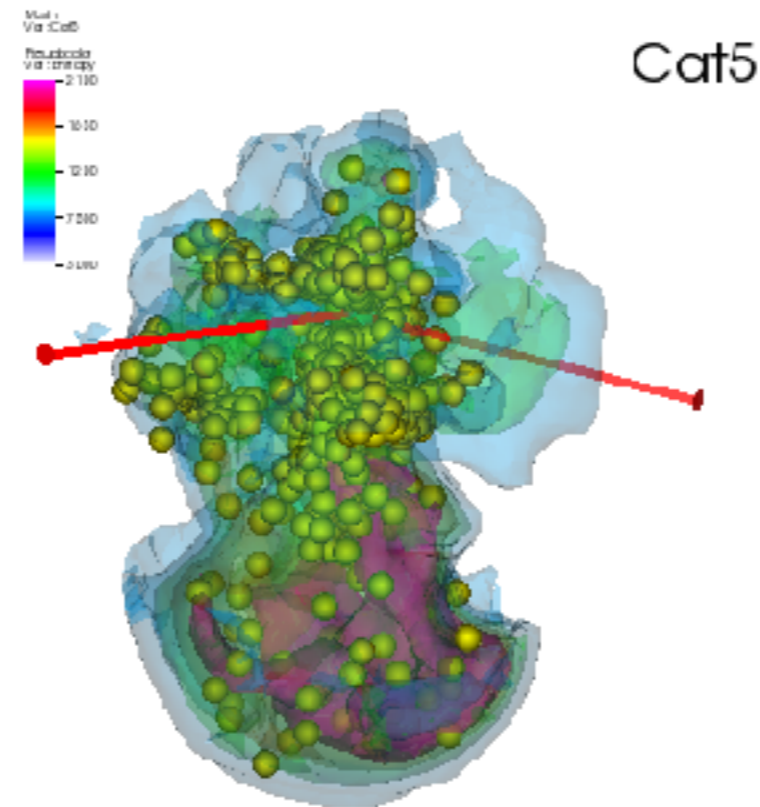
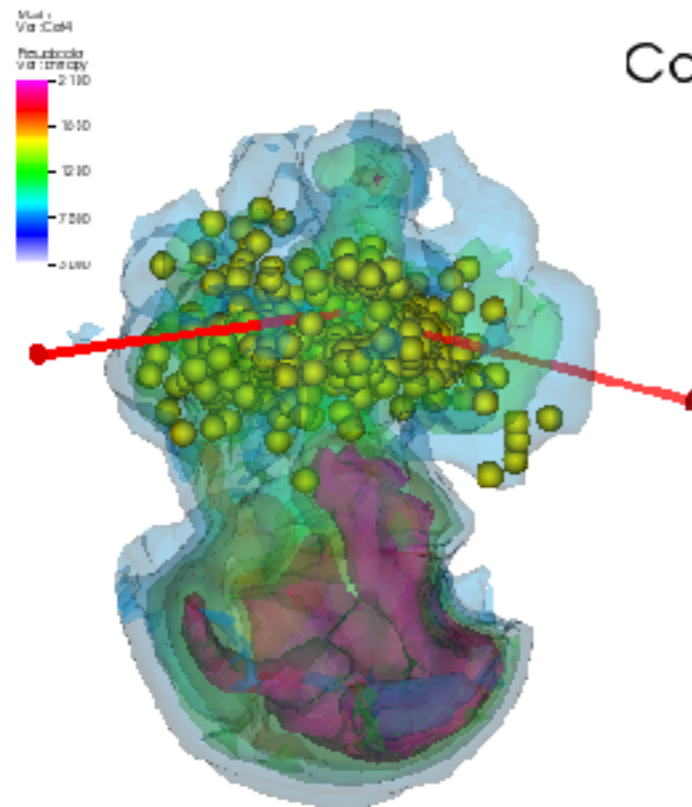
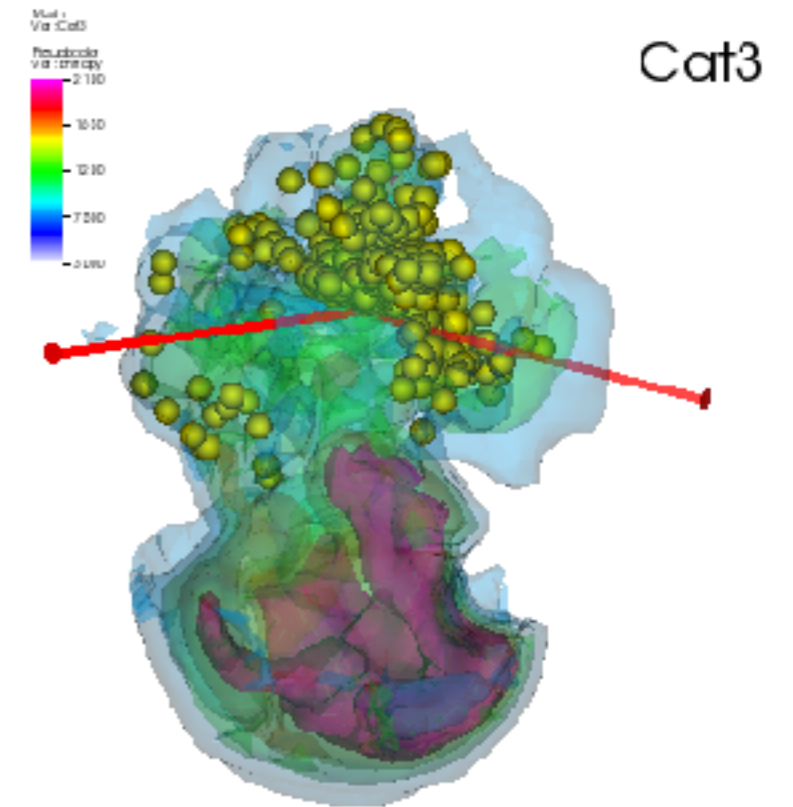
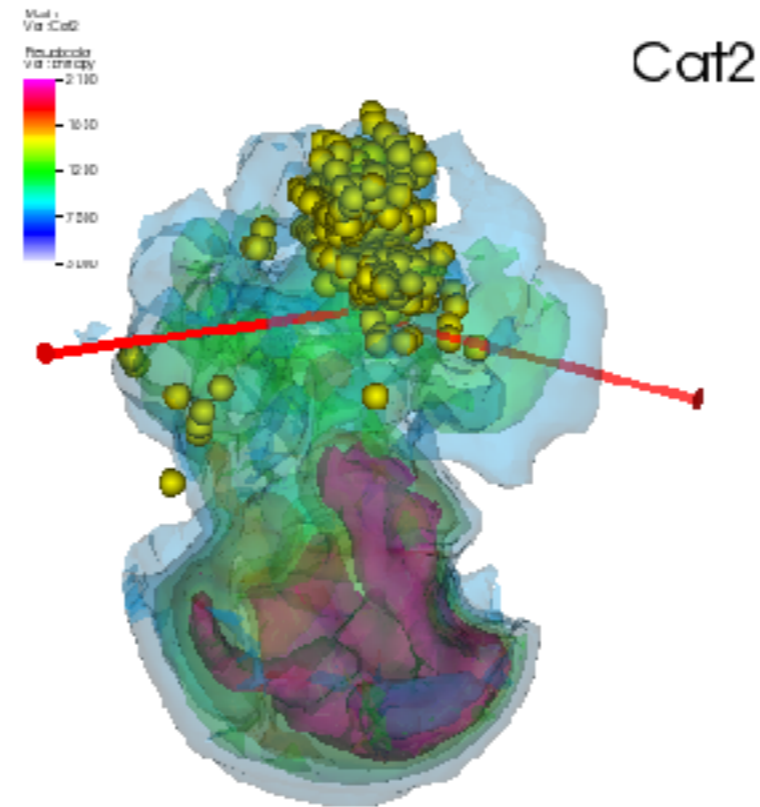
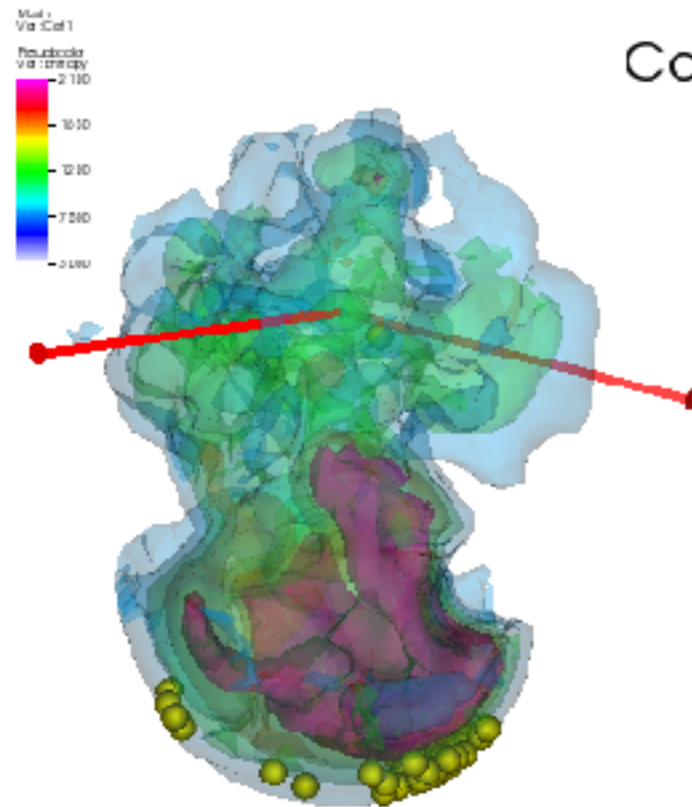


Scattering in S- Y_e plane

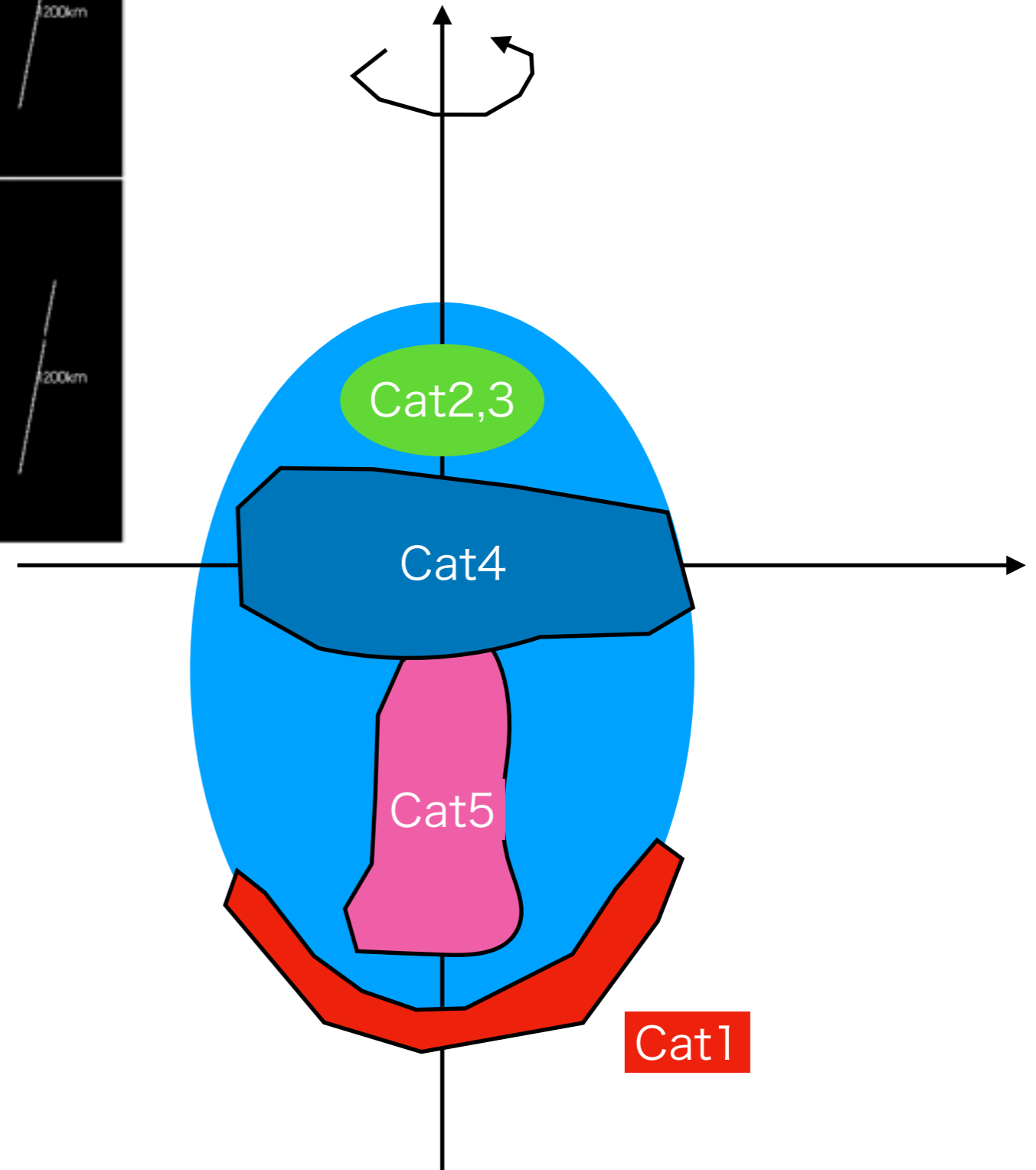
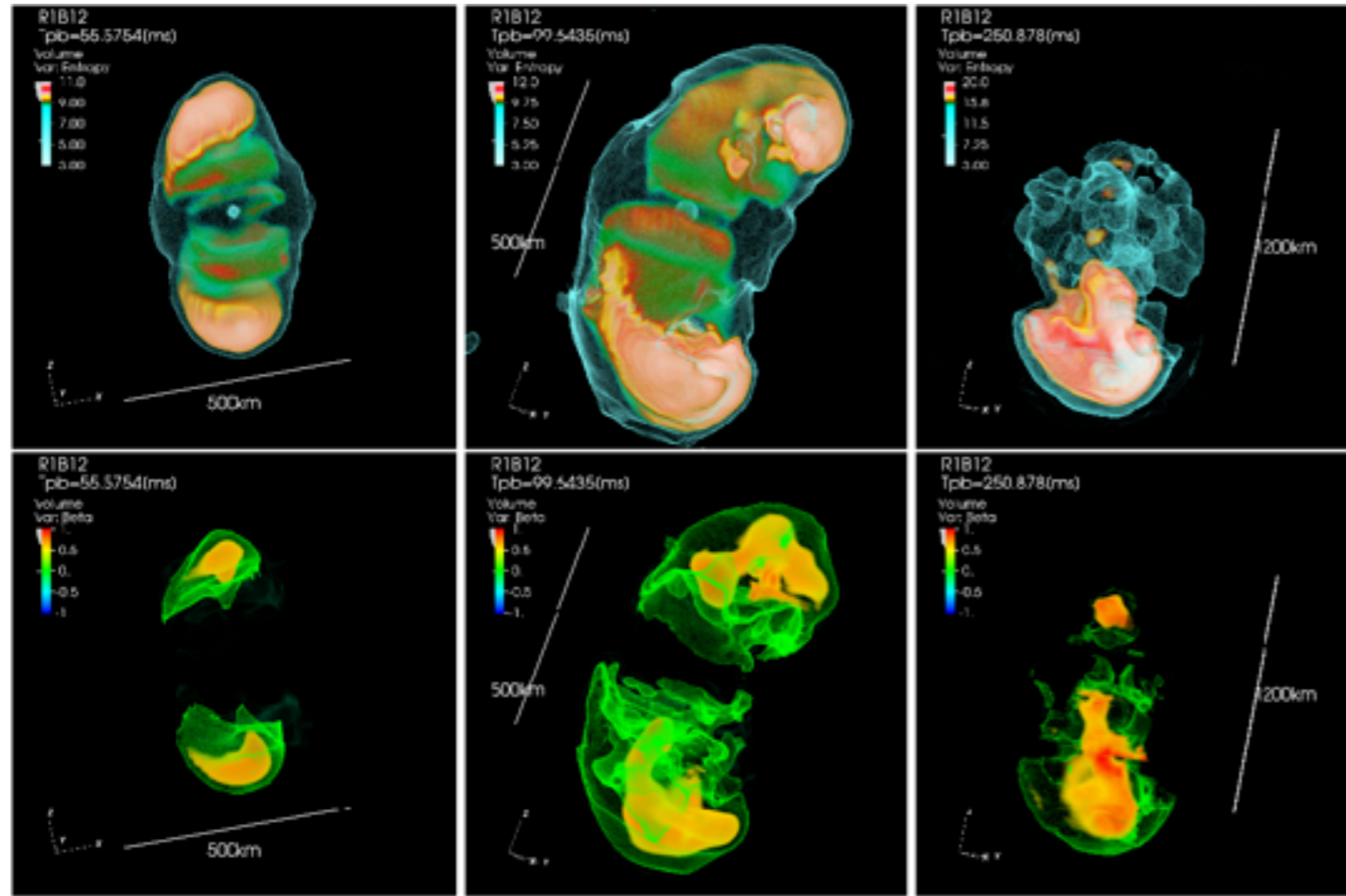
Due to the interaction $\nu_e n \rightarrow e^- p$
 Y_e & S tend to become higher



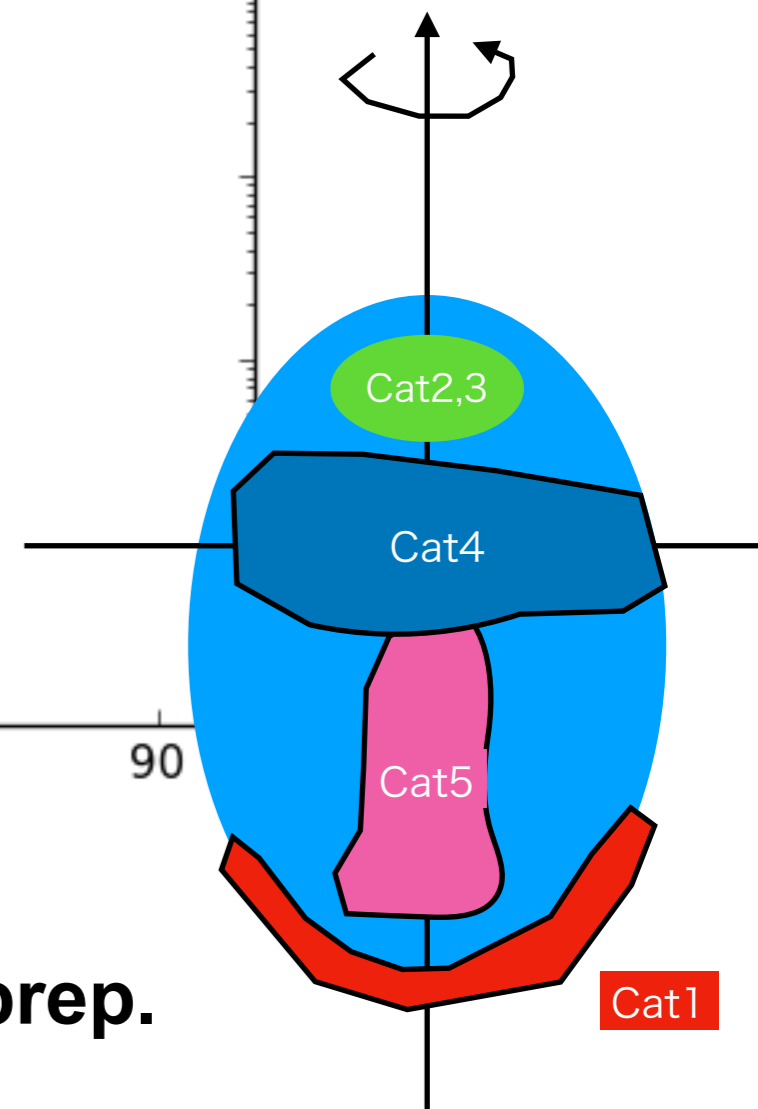
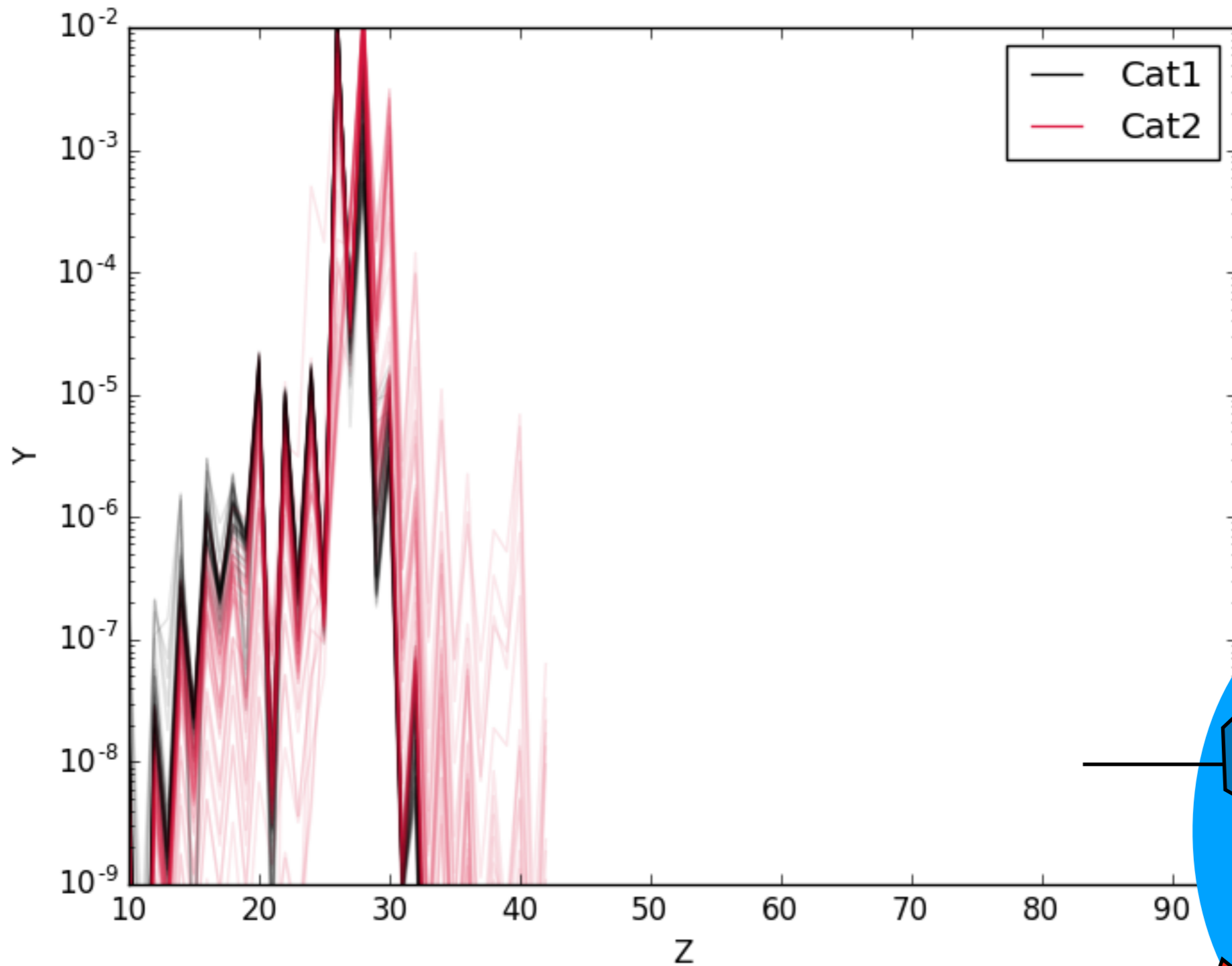
Ejecta distribution



Ejecta distribution



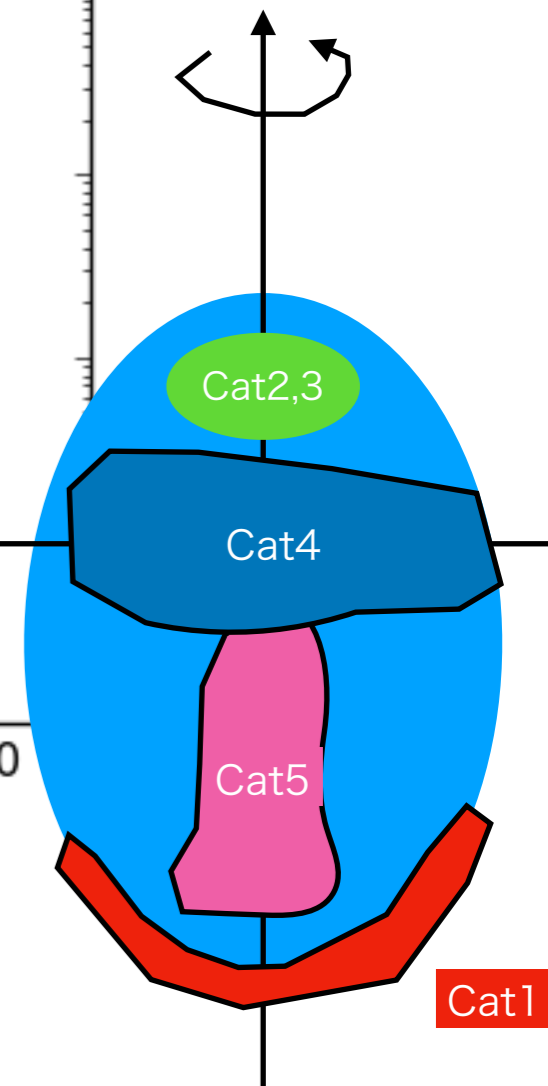
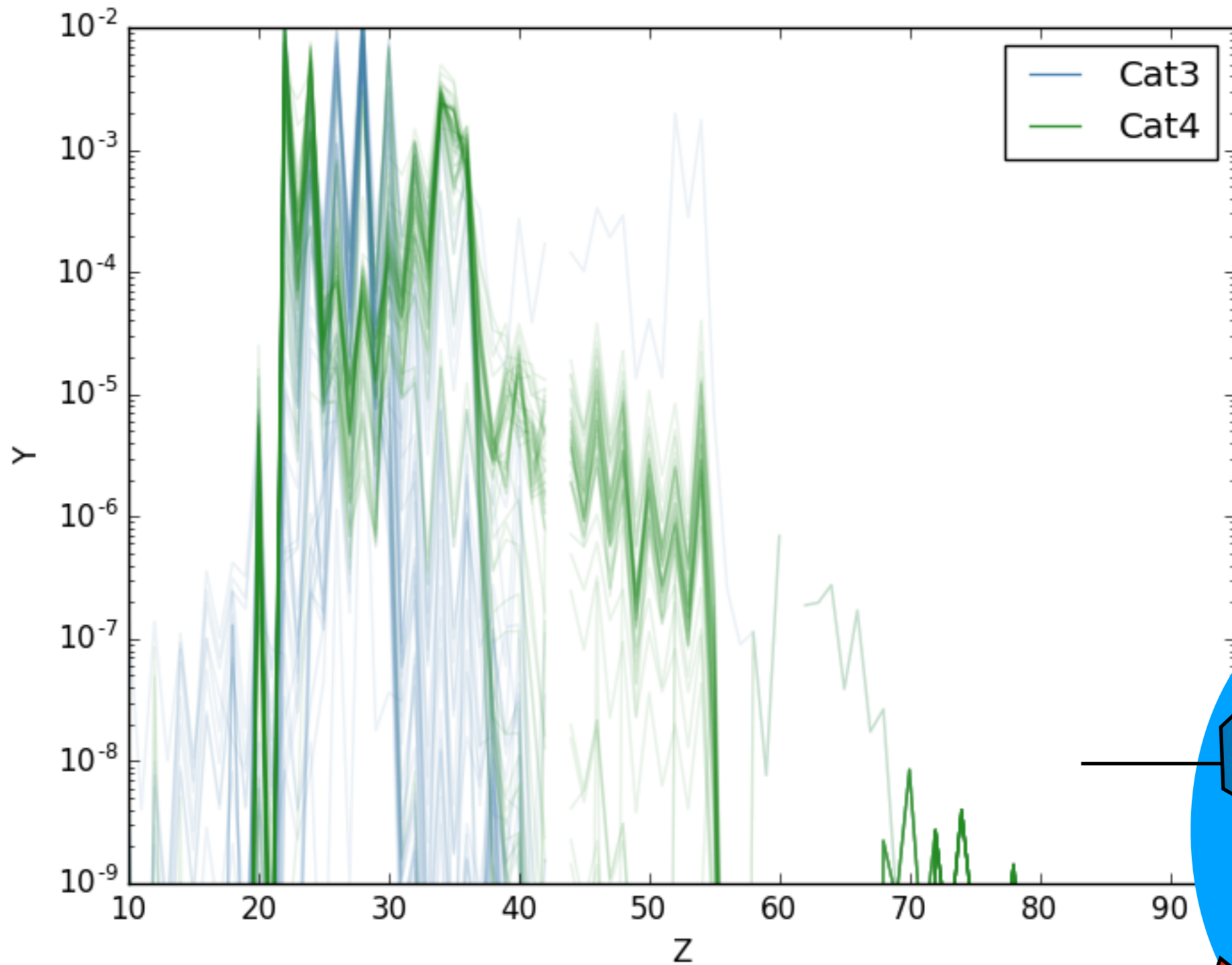
Nucleosynthesis (1st peak)



Reichert, Eichler, TK+, in prep.

Cat1

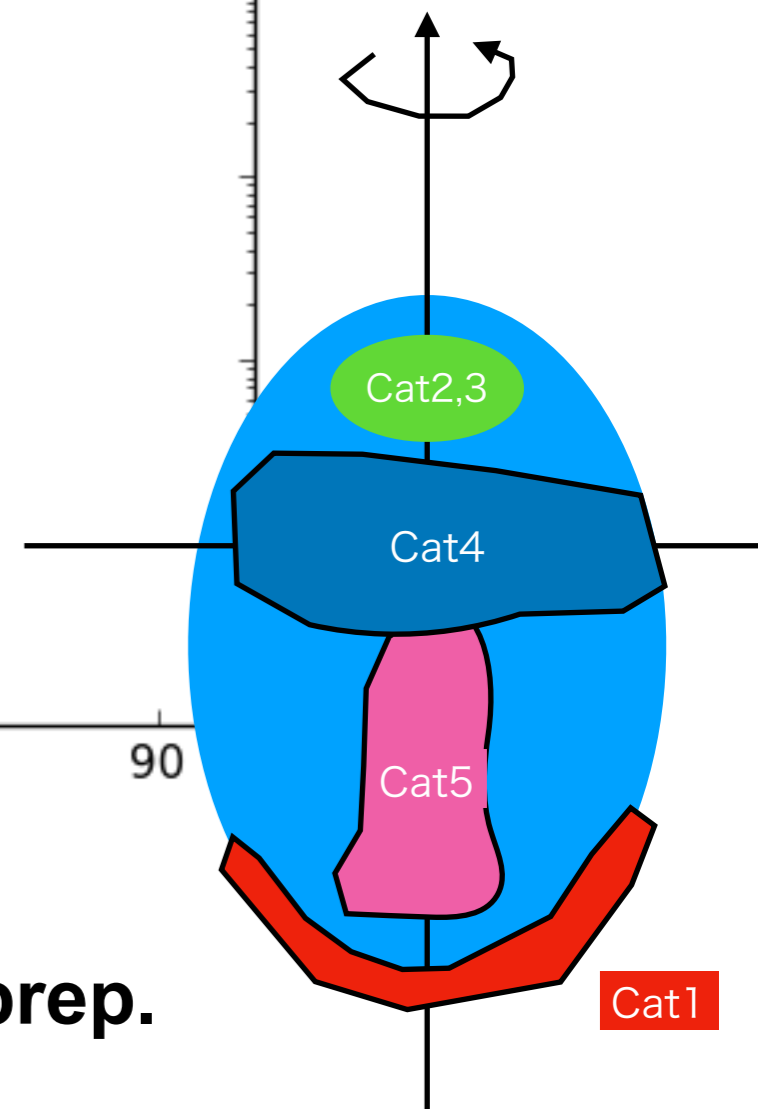
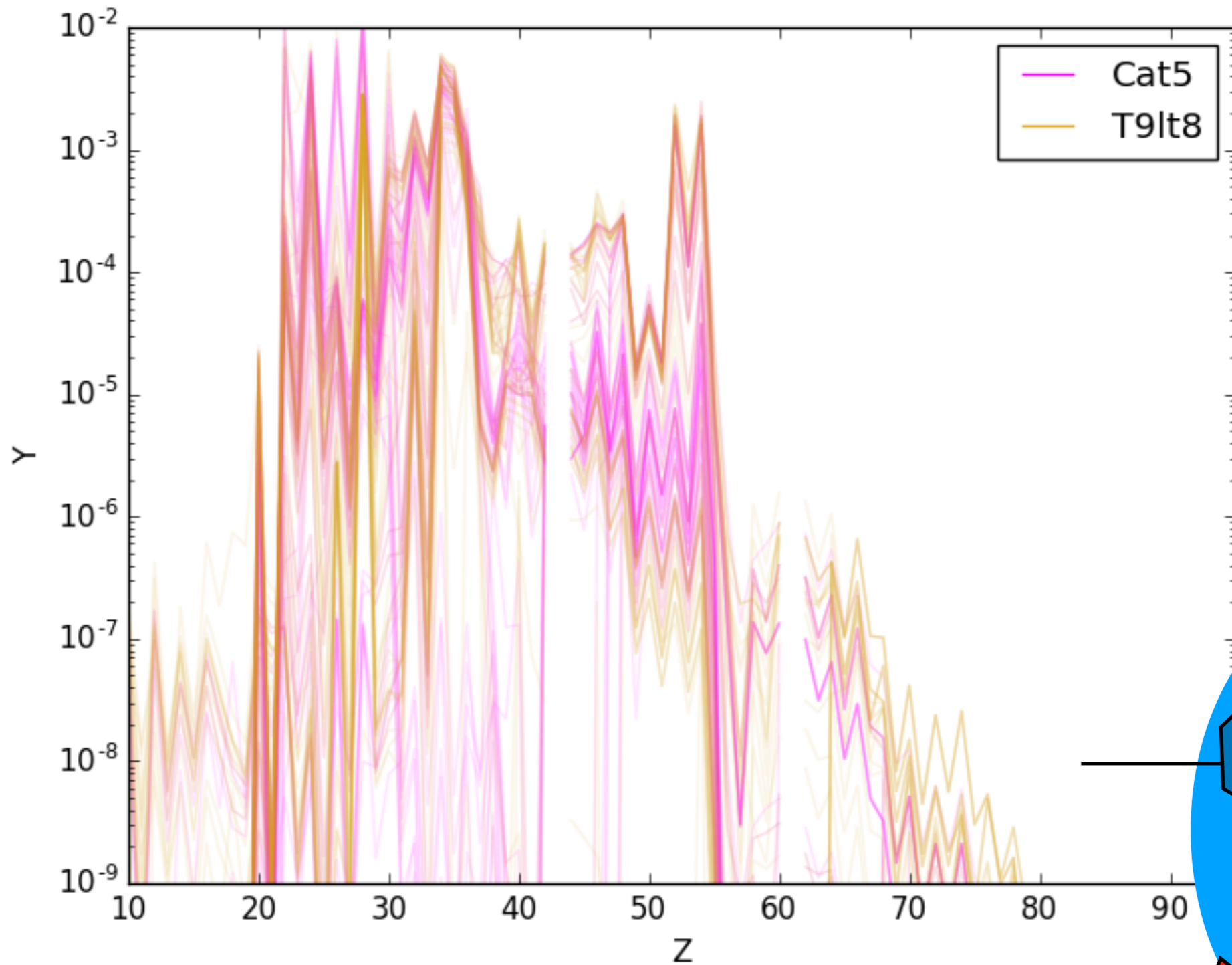
Nucleosynthesis (weak 2nd peak)



Reichert, Eichler, TK+, in prep.

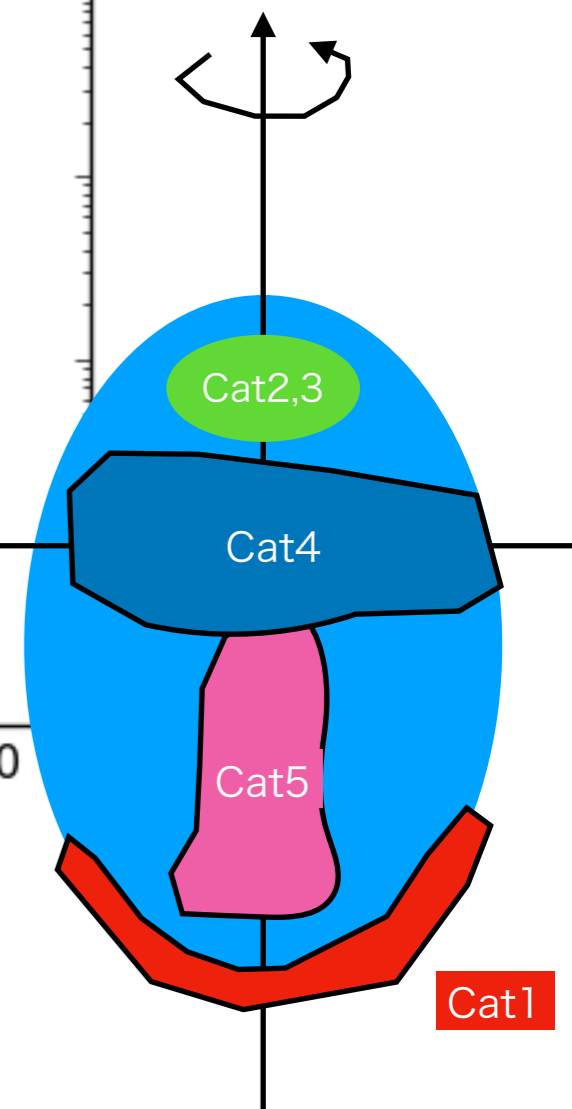
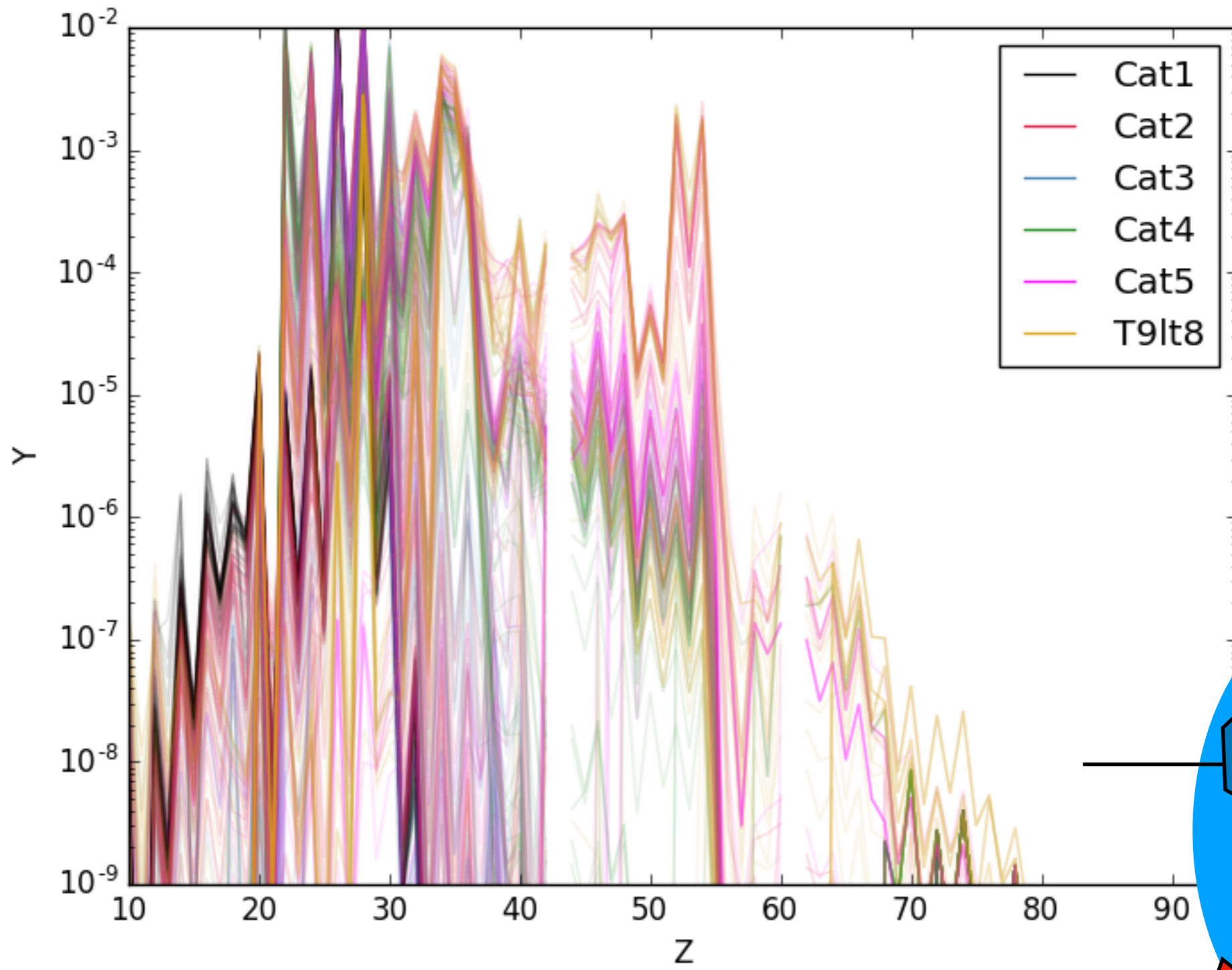
Cat1

Nucleosynthesis (2nd + weak 3rd peak)



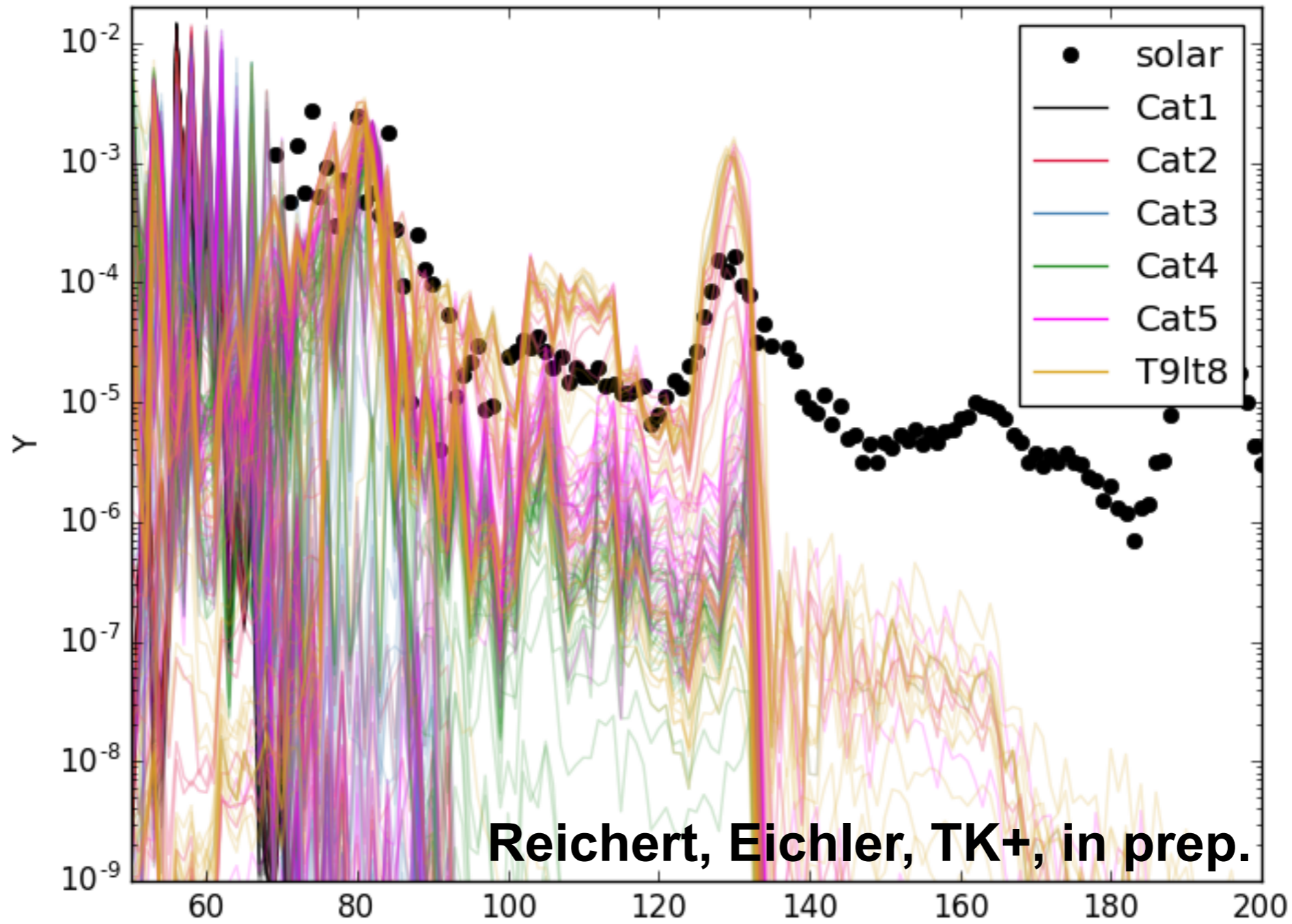
Reichert, Eichler, TK+, in prep.

Nucleosynthesis



Reichert, Eichler, TK+, in prep.

Nucleosynthesis (1st peak)



- Not enough to produce 3rd peak
- consistent with Mösta+, '18

Summary

1. SN simulations are becoming more realistic (full GR, 3D effects, sophisticated neutrino opacities) —> more reliable messages from SNe (GWs, **neutrinos**, and **heavy elements**)
2. In MHD model, the polar/**equatorial** explosion is boosted mainly by **B/neutrinos**.
3. 2nd & weak 3rd peak elements can be produced with significant **B/neutrino** effects depending on the trajectory
4. Temporal modulation in neutrinos reflecting the **SASI** motions.