Three-body calculations of the triple-alpha reaction rate at low temperatures

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1. INTRODUCTION

Triple-alpha reaction

 $^{4}\text{He} + ^{4}\text{He} + ^{4}\text{He} \rightarrow ^{12}\text{C}$

-Resonant process (T>10⁸ K) ⁸Be, ¹²C* Resonance formula

-Non-resonant process (T<10⁸ K) [A] Extension of the resonance formula with energy dependent widths. -NACRE [1]

[B] Quantum mechanical 3-body calculations
 -OKK: CDCC calculations (Ogata et al.[2]) Significant effects at low temperature





[1] C. Angulo et al., NPA**656** (1999) 3. [2] K. Ogata et al., PTP**122** (2009) 1055.

Astrophysical input: 3α reaction rate $<\alpha\alpha\alpha\alpha>$ [cm⁶/s]



n₁₂ (n₄): Number density of ¹²C (⁴He)



3-body calculations for 3α reaction

- Ccontinuum Discretized Coupled Channel (CDCC) K. Ogata et al., PTP122 (2009) 1055. [OKK]
- Hyperspherical Harmonics basis + R-matrix (HHR) N.B. Nguyen et al. arXiv:1112.2136, arXiv:1209.4999
- This workshop: K. Yabana

Ref: "Imaginary-time method for the radiative capture reaction rate" K. Yabana and Y.Funaki, PRC **85**, 055803 (2012)

Faddeev

S. Ishikawa, INPC2010, APFB2011, OMEG11 (paper in preparation)

In the present talk:

- Calculation of 3α reaction based on the Faddeev 3-body theory.
- Discussion about the difference from the OKK rate

CONTENTS

- (1. Introduction)
- 2. Formalism
- 3. Calculations and Results
- 4. Discussion -- Comparison with CDCC results
- 5. Summary



2. FORMALISM

3 α reaction $\alpha + \alpha + \alpha \rightarrow {}^{12}C(2^+) + \gamma$

1. Inverse reaction: Photo induced 3α breakup of ${}^{12}C(2^+)$ ${}^{12}C(2^+) + \gamma \rightarrow \alpha + \alpha + \alpha$

 $E = E_q + E_p$

 E_p

 $X E_q$

3. Photodisintegration cross section

2

$$\sigma_{\gamma}(E) \propto \iint d\hat{x} d\hat{y} \int_{E_q > 0} dE_q \sqrt{E_q E_p} \left| f^{(B)}(E_q; \hat{x}, \hat{y}) \right|^2$$

4. Reaction rate

$$\left\langle \alpha \alpha \alpha \right\rangle = 240 \left(3\right)^{3/2} \pi \left(\frac{\hbar}{mc}\right)^3 c \int_0^\infty \frac{dEE_{\gamma}^2}{\left(kT\right)^3} e^{-E/k_B T} \underline{\sigma}_{{}^{12}C(2_1^2)+\gamma\to 3\alpha}\left(E_{\gamma}\right) \qquad \left(E_{\gamma} = E - E_{{}^{12}C(2^+)}\right)$$

- 5. Apply the Faddeev formalism [1] to solve the equation for the 3-body disintegration process.
- Apply the Sasakawa-Sawada method [2] to accommodate the long-range Coulomb interaction.
- 7. An approximation is made to treat a long-range contribution

[1] L.D. Faddeev, Soviet Phys. JETP **12** (1961) 1041.[2] T. Sasakawa and T. Sawada, PRC 20 (1979) 1954.

Faddeev eq. (1961) Multiple scattering with rearrangements 2 2 $\Phi^{(3)}(3,12)$ Channel-1 Channel-3 $\Phi^{(1)}(1,23)$ 3 (Faddeev component) _2 $\Phi^{(2)}(2,31)$ Channel-2 Totally symmetric 3 $\Psi(123) = \Phi^{(1)}(1,23) + \Phi^{(2)}(2,31) + \Phi^{(3)}(3,12)$ Symmetric for 2<->3

An approximation

• A term $\left(\frac{1}{x_3} - \frac{1}{y_1}\right)$ appeared in the integral

kernel, which is expected to be short range because of a cancellation. But, the cancellation is not perfect for breakup channels.

→ treat this problem approximately by a (mandatory) cutoff procedure

$$\left(\frac{1}{x_3} - \frac{1}{y_1}\right) \times e^{-(x_3/R_{\rm cut})^4}$$

$$R_{\rm cut} = 20 \,\mathrm{fm}$$
 - 35 fm





3. CALCULATIONS AND RESULTS

3α model • αα-potential Ali-Bodmer type (2-range Gaussian)

$$V_{\alpha\alpha}(x) = \left(V_R^{(0)}\hat{P}_{L=0} + V_A^{(2)}\hat{P}_{L=2}\right)e^{-(x/a_R)^2} + V_A e^{-(x/a_A)^2}$$

	a _R (fm)	V _R ⁽⁰⁾ (MeV)	V _R ⁽²⁾ (MeV)	a _A (fm)	V _A (MeV)
AB(A')	1.53	125.0	20.0	2.85	-30.18
AB(D)	1.40	500.0	320.0	2.11	-130.0

• **3-body potential** [1] to reproduce the binding energy and resonance energy

$$V_{\alpha\alpha\alpha} = \left(W_0 \hat{P}_{L=0} + W_2 \hat{P}_{L=2} \right) e^{-(\rho/3.9)^2}$$

$$\rho^2 = 3.97 \sum_{i=1}^{3} r_i^2$$

[1] D.V. Fedorov and A. S. Jensen, PLB **389** (1996) 631



$\alpha\alpha\alpha$ reaction rate (AB-A', AB-D)



Comparison with CDCC results

4. DISCUSSION

CDCC calculation of photo induced 3α breakup of ${}^{12}C(2^+)$ ${}^{12}C(2^+) + \gamma \rightarrow \alpha + \alpha + \alpha$

Wave function for (photo-) disintegration process

$$\begin{split} \left|\Psi\right\rangle &\equiv \frac{1}{E + i\varepsilon - H_0 - V} H_{\gamma} \left|\Psi_b\right\rangle \quad \xrightarrow[R \to \infty]{} \quad \frac{e^{iKR}}{R^{5/2}} f^{(B)}\left(E_q, x, y\right) \\ \left(E - H_0 - V\right) \left|\Psi\right\rangle &= H_{\gamma} \left|\Psi_b\right\rangle \end{split}$$

Discretized α - α functions $u_n(x)$: $[T_x + V_{\alpha\alpha}(x)]u_n(x) = E_{q_n}u_n(x)$ $\Psi(x, y) = \sum_n u_n(x)\varphi_n(y)$

 $\sigma_{\gamma}(E) \propto \sum_{n < n_{\gamma}} \frac{|T_n|^2}{p_n}$

 $\varphi_n(y) \rightarrow [\text{Outgoing wave}] \times T_n$

$$\sum_{n'} \left[\left(E_p - T_y \right) \delta_{n,n'} - V_{n,n'}(y) \right] \varphi_{n'}(y) = \left\langle u_n \left| H_\gamma \right| \Psi_b \right\rangle$$

of base functions = 120 (~OKK)

Photodisntegration cross section (AB-A') (CDCC calculations by S.I.)





- At low temperatures (T<10⁸ K): $<\alpha\alpha\alpha>_{NACRE} \sim <\alpha\alpha\alpha>_{Faddeev} << <\alpha\alpha\alpha>_{CDCC}$
- Explanation of this enhancement by Ogata: Coulomb barrier between αα-pair and α: non-resonant pair vs. resonant pair



Reason for the enhancement (Ogata)

• Coulomb potential between $\alpha\alpha$ -pair and α -particle



Model space of CDCC calculation

 Only one set of Jacobi coordinate is used: Neglects of rearrangement channels as well as symmetrization of the wave functions



Rearrangement effect



3α decay mechanism of the Hoyle state

The enhancement of σ_γ(E) by the CDCC calculation at low energies is due to the reduction of Coulomb barrier between α and non-resonant αα-pair.

 This reduction may cause an enhancement of non-resonant (direct) process of 3α-decay of the Hoyle state.



3α decay of the Hoyle state

- Direct decay or Sequential two-step process
- Ad.R. Raduta et al., PLB **705**, 65 (2011). 40 Ca + 12 C at 25MeV/nucleon Direct-decay contribution: 7.5 ± 4.0 %
- O. S. Kirsebom et al. PRL 108, 202501 (2012).
 ¹¹B(³He,d)
 "no evidence for direct-decay branches"
- J. Manfredi et al. PRC 85, 037603 (2012).
 ¹⁰C + ¹²C "An upper limit of 0.45%"

Decomposition of the cross section

$$\sigma_{\gamma}(E) \propto \iint d\hat{x} d\hat{y} \int_{E_q > 0} dE_q \sqrt{E_q E_p} \left| f\left(E_q; \hat{x}, \hat{y}\right) \right|^2$$

$$\sigma_{\gamma}^{R}(E) \propto \iint d\hat{x} d\hat{y} \int_{E_{q}=E_{r}\pm\Delta E} dE_{q} \sqrt{E_{q}E_{p}} \left| f\left(E_{q};\hat{x},\hat{y}\right) \right|^{2}$$

$$\sigma_{\gamma}^{\mathrm{NR}}(E) = \sigma_{\gamma}(E) - \sigma_{\gamma}^{\mathrm{R}}(E)$$



Faddeev vs. CDCC (SI)

3α-decay of the Hoyle state
 Sequential decay vs. Direct decay

--Faddeev: Sequential decay-dominant

--CDCC: large contribution from Direct decay 67% at E=380keV

Non-resonant contribution (CDCC)



Non-resonant contribution (Faddeev)



5. SUMMARY

- Quantum mechanical 3-body calculations of 3α-reactionas photodisintegration of ¹²C(2⁺) Faddeev method, CDCC method
- Faddeev calculation: similar to the NACRE 3α rate
- CDCC calculations: Increase of the cross section at low energies (similar to Ogata's CDCC results)
- 3α-decay of Hoyle resonance
 Faddeev: Sequential decay (via ⁸Be) dominant
 CDCC: A large contribution from the Direct decay
 → This may be tested by experiments.
- Future problem:

• Higher energies (theoretical calculations of ¹²C-resonance other than the Hoyle state)

- . ⁹Be(α - α -n), ⁶He(α -n-n), n-n-n (3-n potential)
- •4 α problem, ¹²C(α , γ)¹⁶O