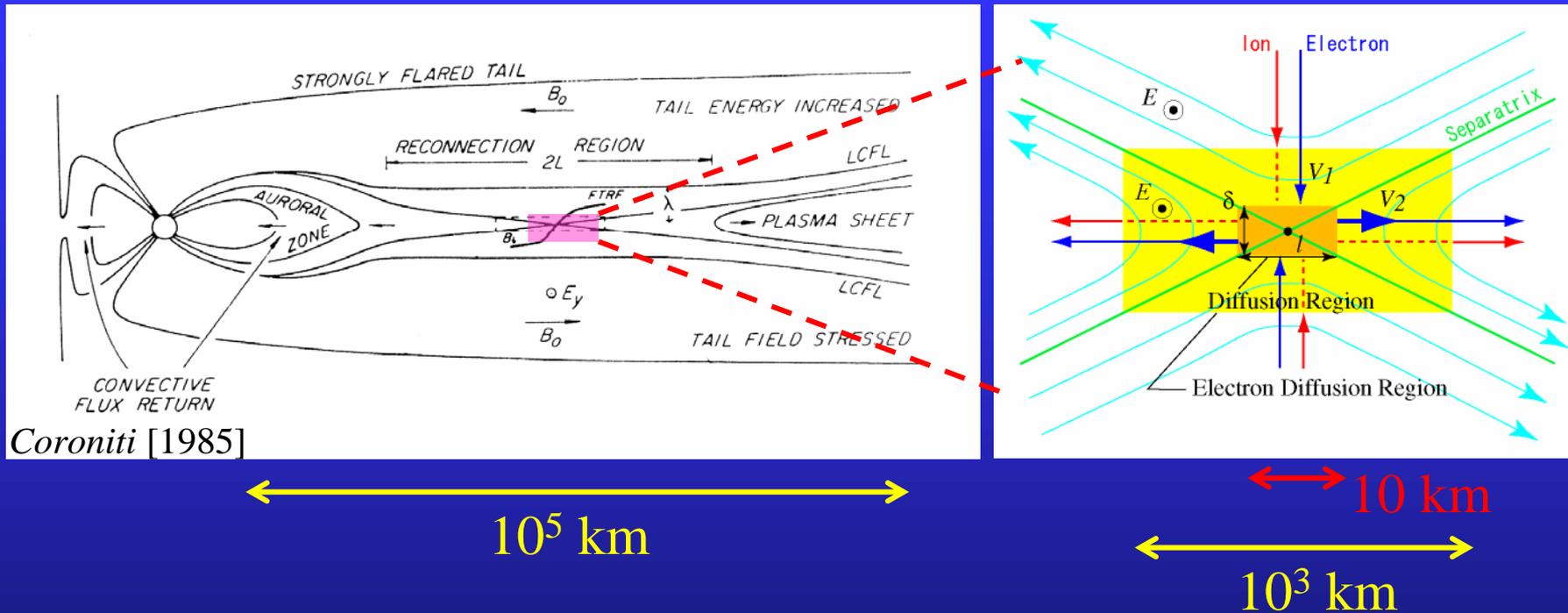


Dissipation Mechanism in 3D Magnetic Reconnection

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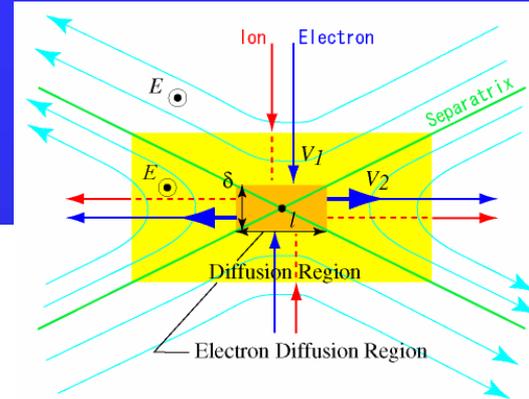
Reconnection (in the Earth Magnetosphere)



- Can induce global-scale convection causing change in the field line topology,
- Is intrinsically multi-scale process.

Why is the dissipation mechanism important?

Our goal is to understand **macroscopic dynamics** in a variety of systems.



Diffusion equation: $\frac{\partial \vec{B}}{\partial t} = \frac{\eta}{\mu_0} \nabla^2 \vec{B}$

η : Electric resistivity \rightarrow Macroscopic parameter originating from **microscopic (kinetic)** processes.

In the MHD framework...

The reconnection rate depends on the resistivity model.

(Biskamp, 1986; Ugai, 1995)

Global responses in substorm and flares are sensitive to the parameterization of the resistivity.

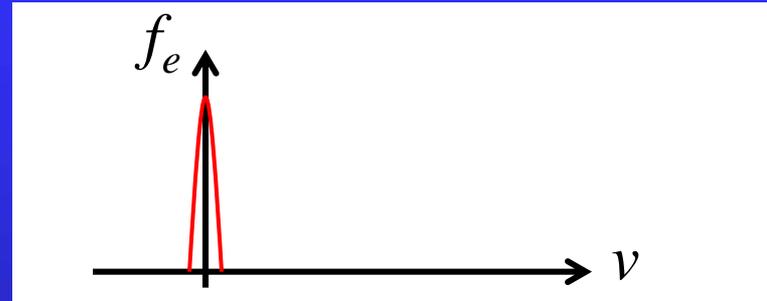
(Raeder et al., 2001; Kuznetsova et al., 2007)

How is the resistivity generated?

The motion of charged particles supporting the current density must be disturbed by “collision”.

Collision with other particles

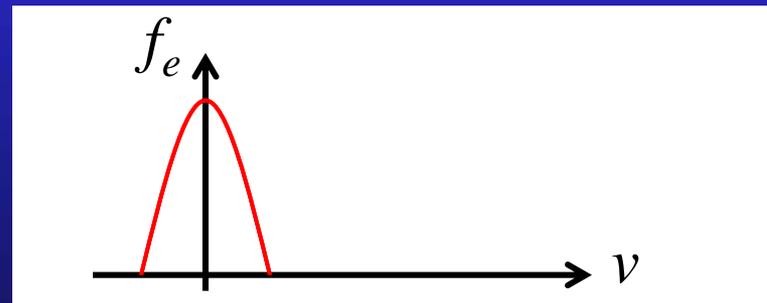
$$\eta = \frac{m_e \nu_c}{n_e e^2}, \quad \nu_c = \frac{1}{\tau_c}$$



➤ *How does the “collision” occur in collisionless plasmas?
= How does the momentum transport occur?*

Inertia resistivity

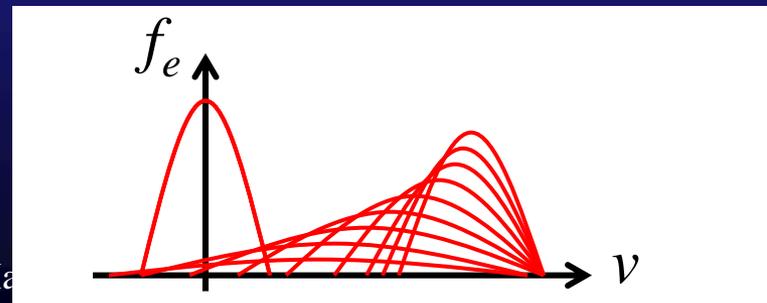
$$\eta = \frac{m_e \nu_T}{n_e e^2}, \quad \nu_T = \frac{1}{\tau_T} \sim \frac{V}{L}$$



Anomalous resistivity
(Wave-particle interaction)

$$\eta = \frac{m_e \nu_w}{n_e e^2}, \quad \nu_w \approx \frac{R_e^{an}}{n_e m_e V_e}$$

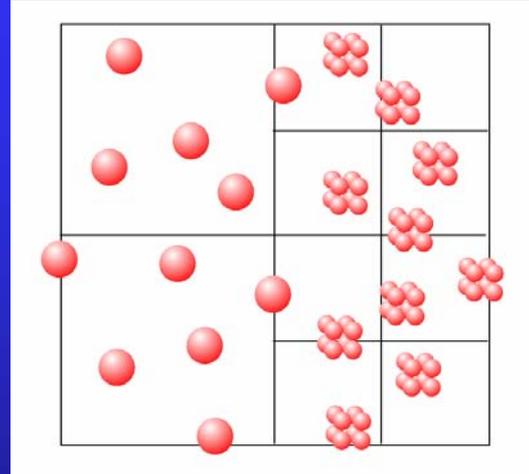
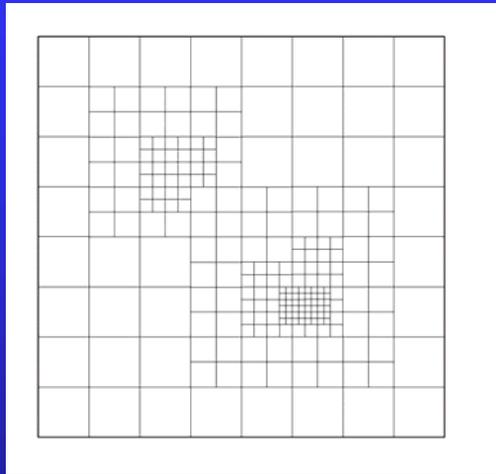
$$R_e^{an} = -e \left(\langle \delta n_e \delta \vec{E} \rangle + \langle \delta(n_e \vec{V}_e) \times \delta \vec{B} \rangle \right)$$



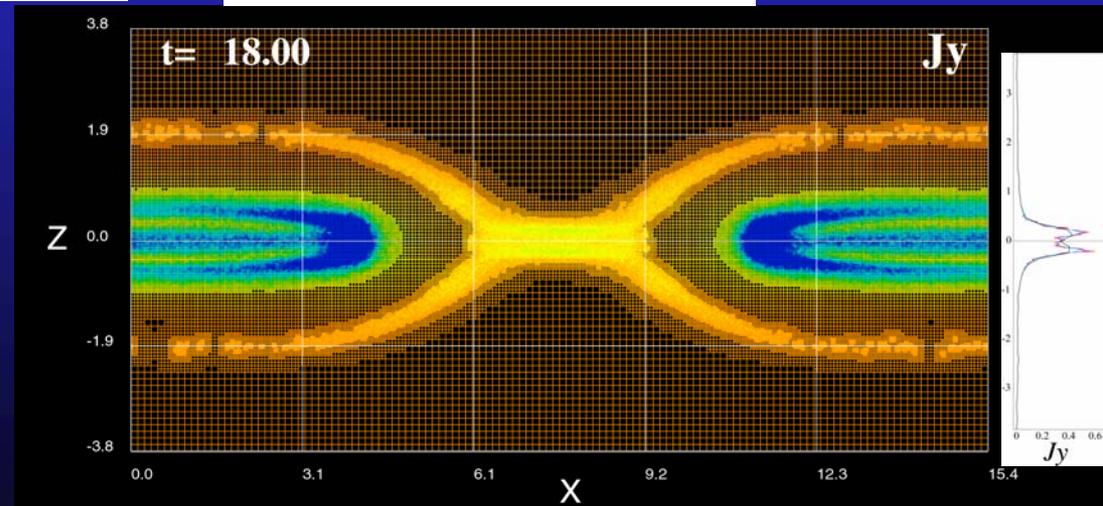
Simulation Model

PIC + Adaptive Mesh Refinement (AMR)

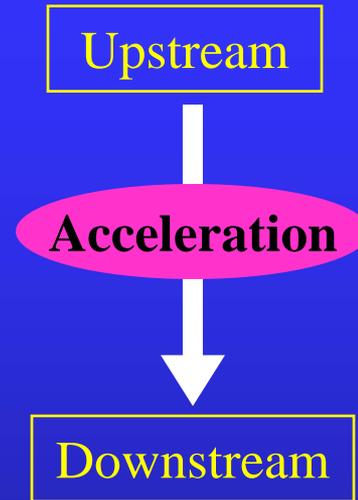
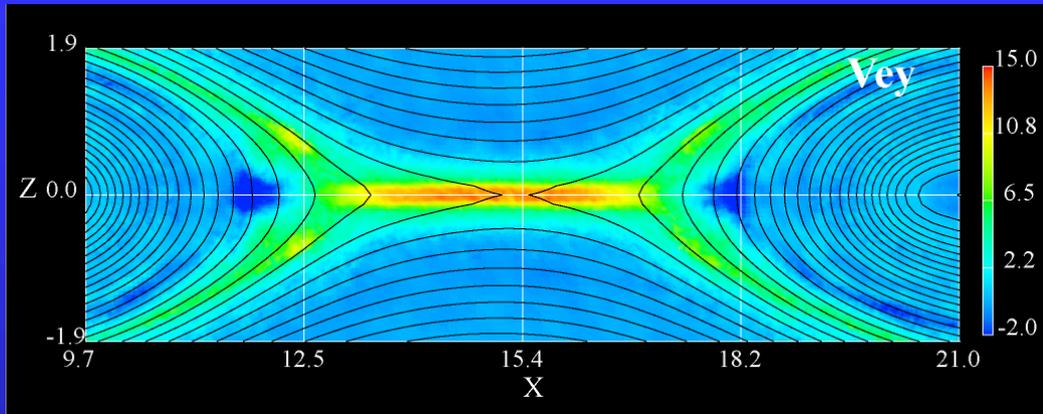
[Fujimoto & Machida, JCP, 2006; Fujimoto & Sydora, CPC, 2008]



Refinement cells are selectively allocated around the **X-line** and **separatrices**.



2D Reconnection

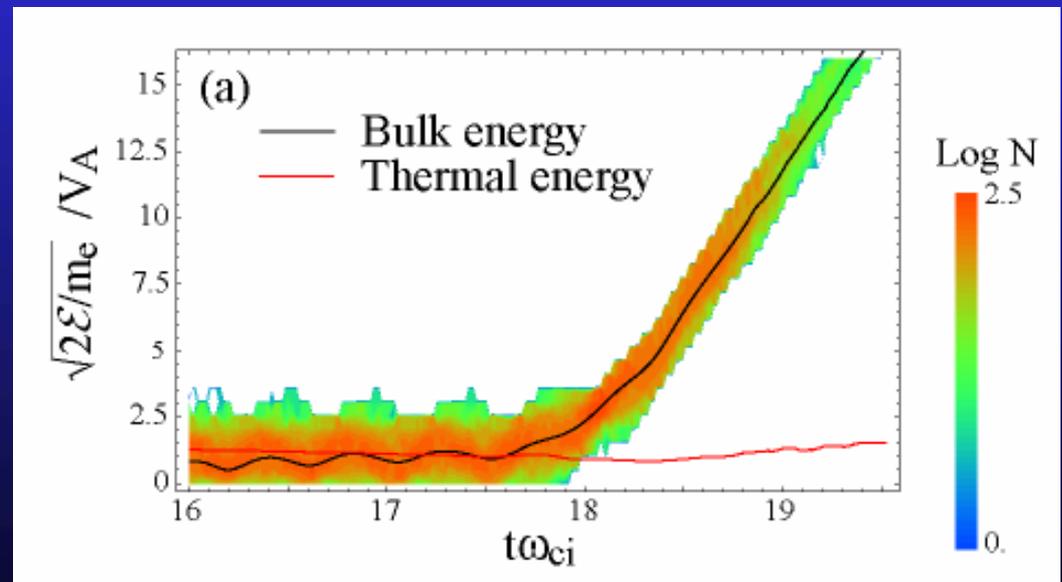


Electrons are...

- Coherently accelerated in the diffusion region,
- Not thermalized.



Inertia resistivity



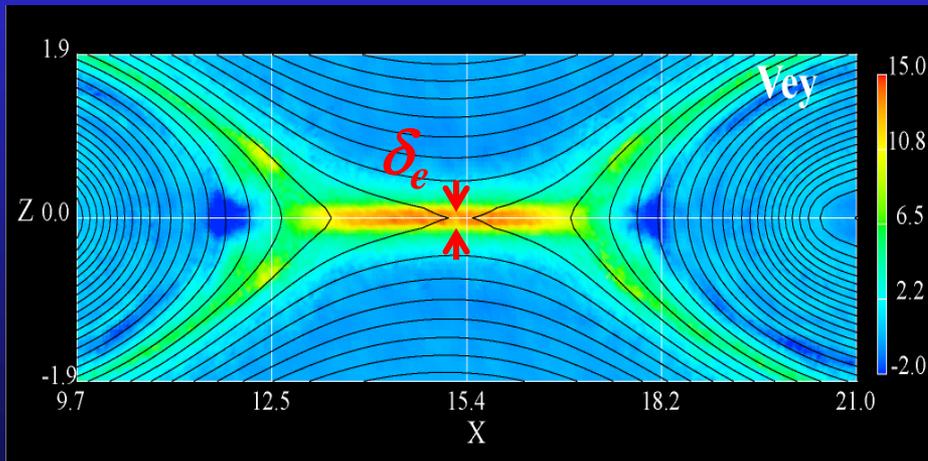
Electron Inertia Resistivity

[Speiser, 1970; Tanaka, 1995; Fujimoto & Sydora, 2009]

$$\eta_{in} = \frac{m_e \nu}{n_e e^2}, \quad \nu = \frac{1}{\tau_{tr}}$$

$$E_R = \eta_{in} j \approx \frac{m_e V_{ey}}{e \tau_{tr}}$$

The electrons must be accelerated **quickly** up to **electron Alfvén velocity**.



$$\delta_e \approx \lambda_e (= c/\omega_{pe})$$

Magnetotail: $\sim 10\text{km}$ ($\Leftrightarrow 10^5\text{km}$)

Solar Flare: $\sim 10^{-5}\text{km}$ ($\Leftrightarrow 10^4\text{km}$)

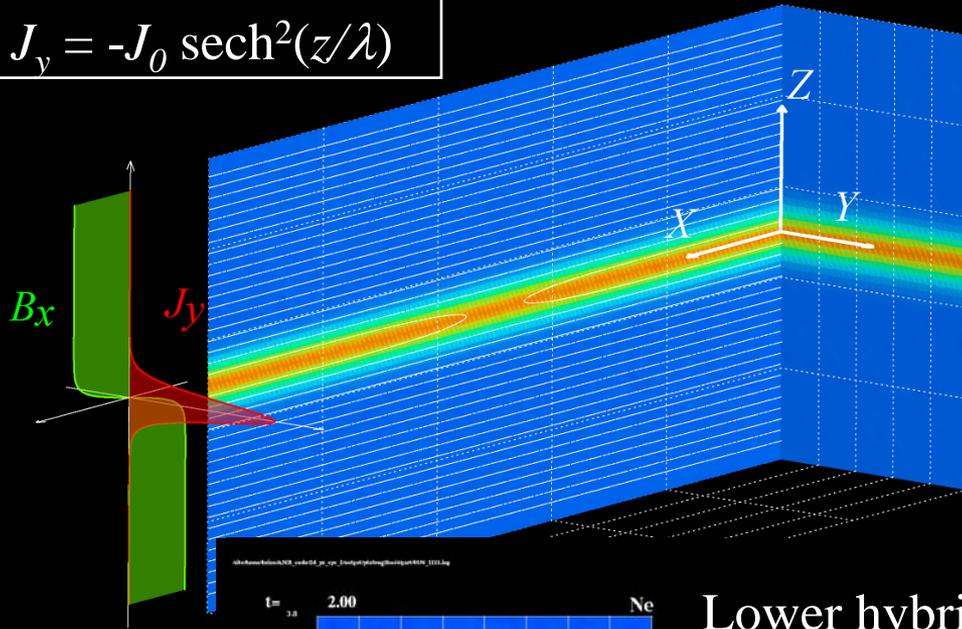


Is such a thin current sheet really stable in 3D system?

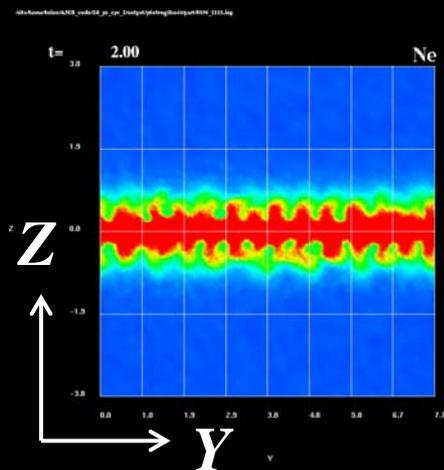
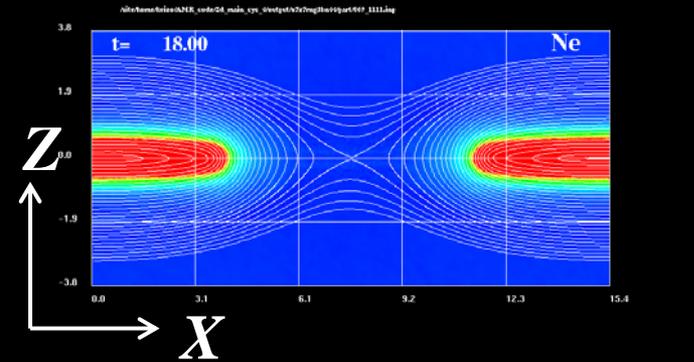
Unstable Modes Expected in the Current Sheet

$$B_x = -B_0 \tanh(z/\lambda)$$

$$J_y = -J_0 \operatorname{sech}^2(z/\lambda)$$



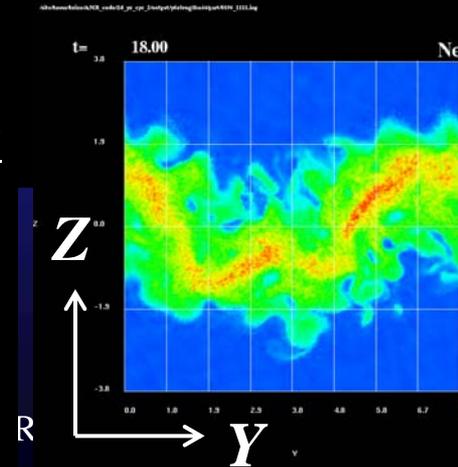
Tearing instability



Lower hybrid drift instability (LHDI)

$$k_y r_{Le} \sim 1$$

$$\gamma \sim \omega_{lh}$$

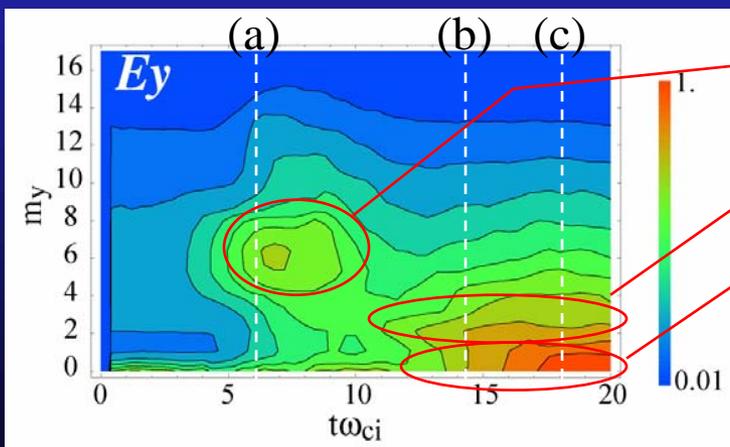
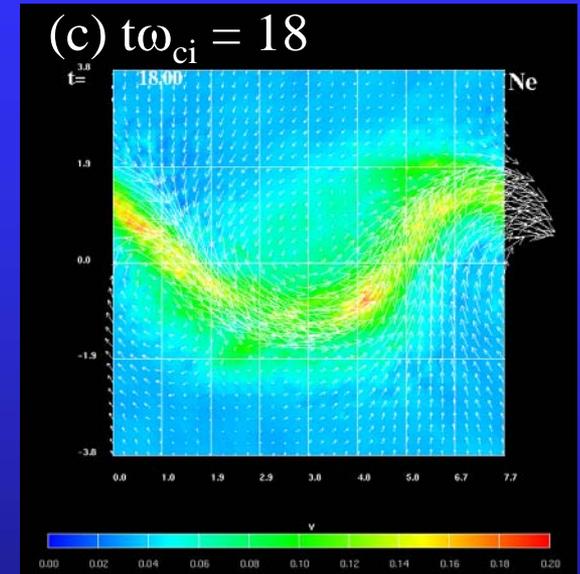
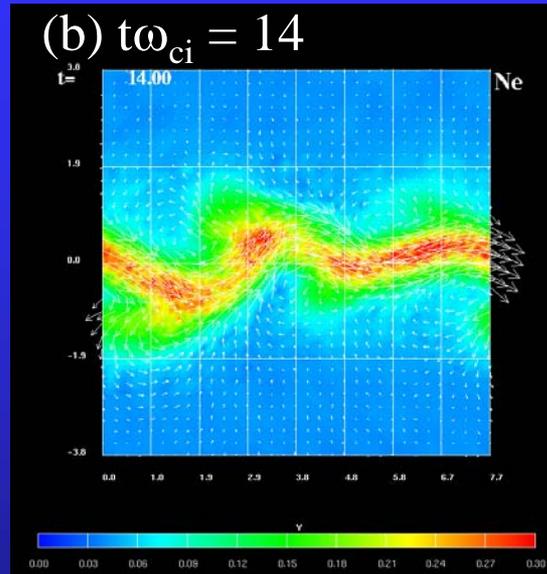
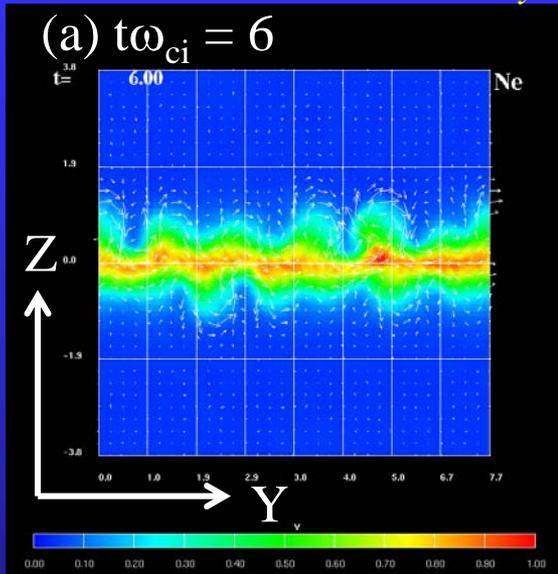


Kink-type instability

$$k_y L \sim 1$$

Time Evolution of the Current Sheet in the YZ Plane

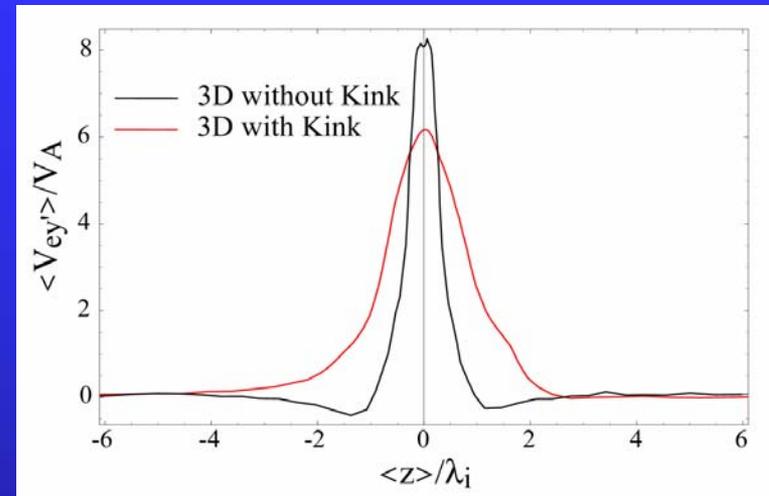
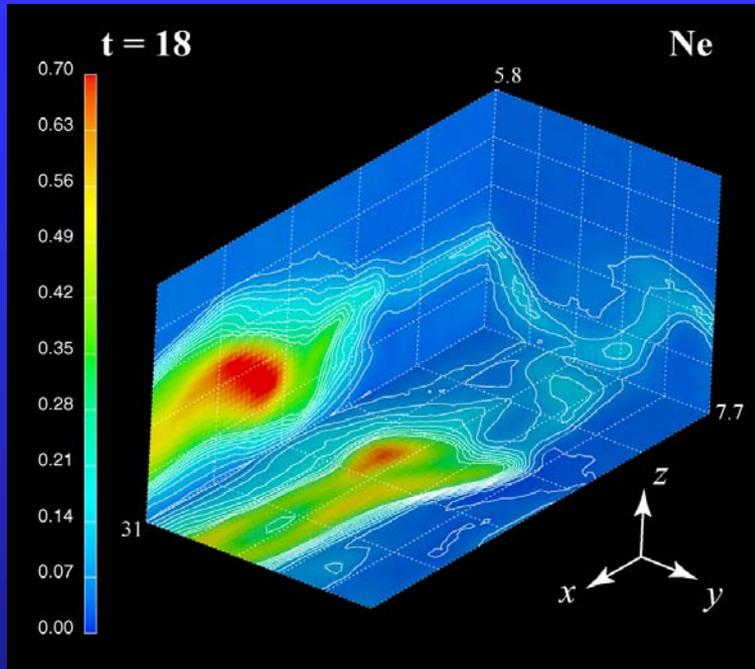
$$m_i/m_e = 25, L_x \times L_y \times L_z = 31 \times 7.7 \times 31$$



- Lower hybrid drift instability (LHDI)
- Kink-type instability
- Induction field due to tearing instability

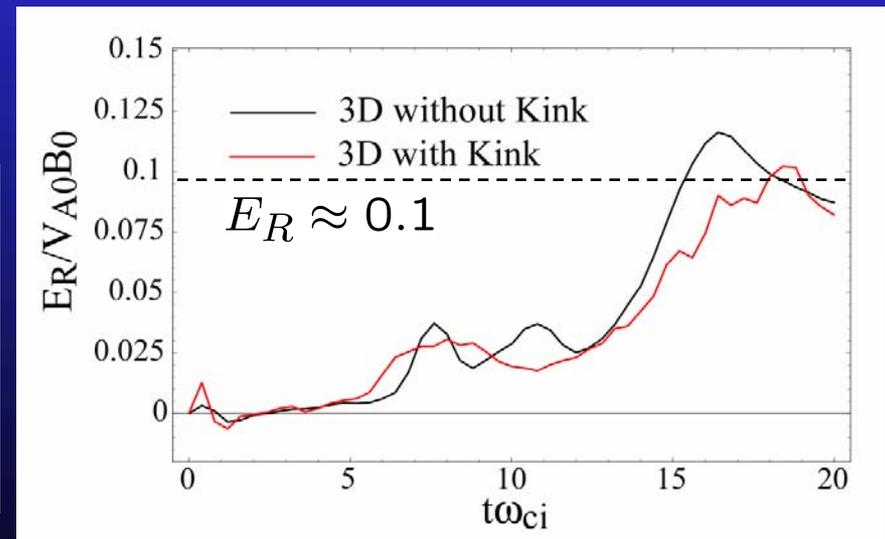
Kink instability coexists with the tearing mode.

Reconnection Rate & Current Sheet Width



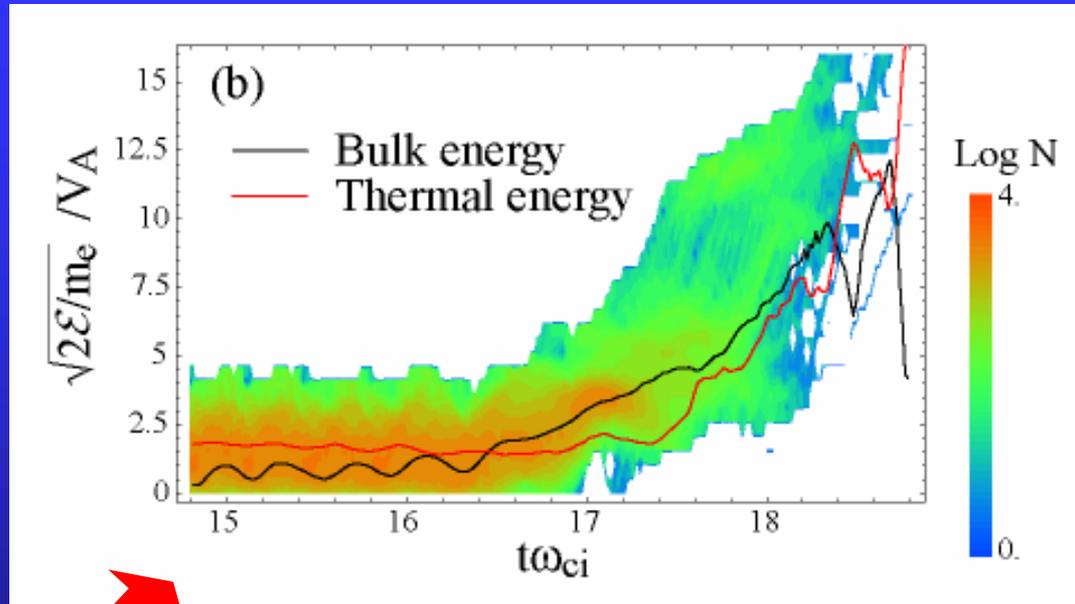
If the inertia resistivity alone supported the dissipation, the E_R would be significantly reduced.

$$\frac{E_{R,kink}}{E_{R,nokink}} \propto \left(\frac{V_{e,kink}}{V_{e,nokink}} \right)^2 \approx 0.56$$

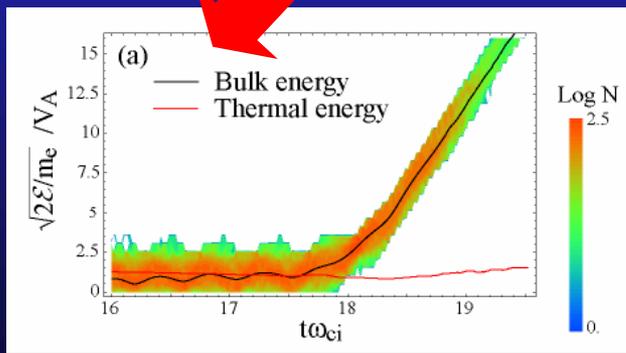
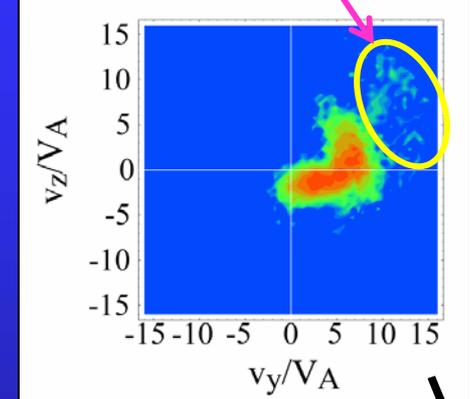


Dissipation Mechanism in 3D Reconnection

[Fujimoto, POP, 2009]

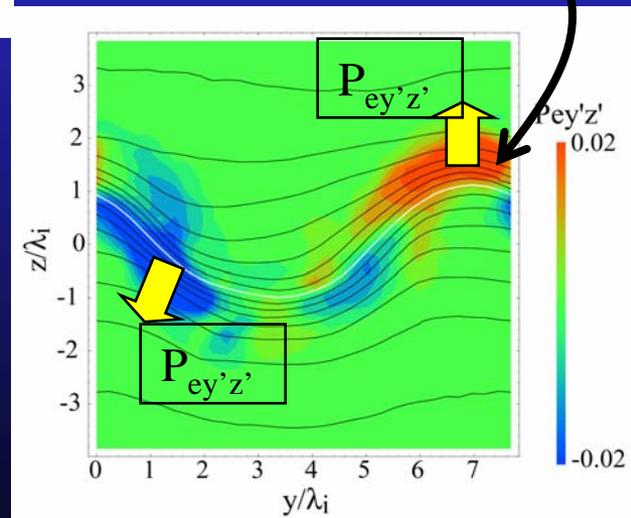


High-energy electrons



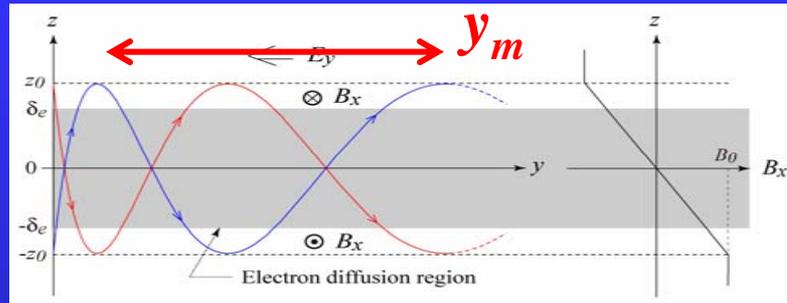
(2D reconnection case)

The electrons are intensely thermalized as well as accelerated in bulk.



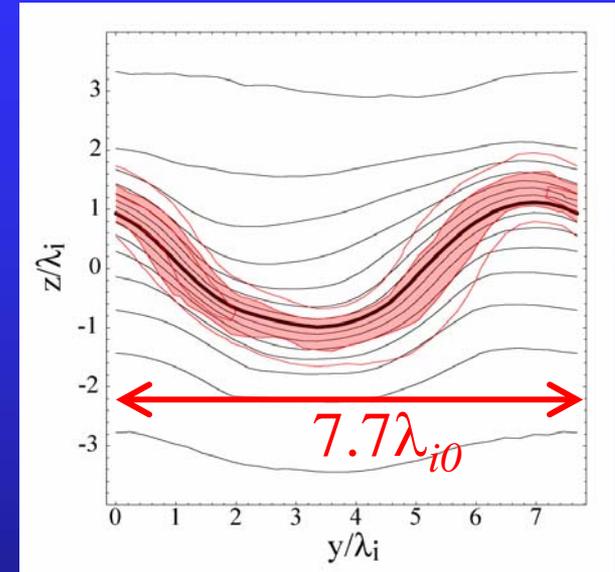
Meandering Scale & Wavelength

➤ Electron meandering scale



$$\omega \approx \frac{2}{3} \frac{V_{ey'}}{c} \omega_{pe} \quad [\text{Speiser, 1965}]$$

$$y_m \approx V_{ey'} \frac{2\pi}{\omega} = 3\pi\lambda_e \approx 9\lambda_{i0}$$



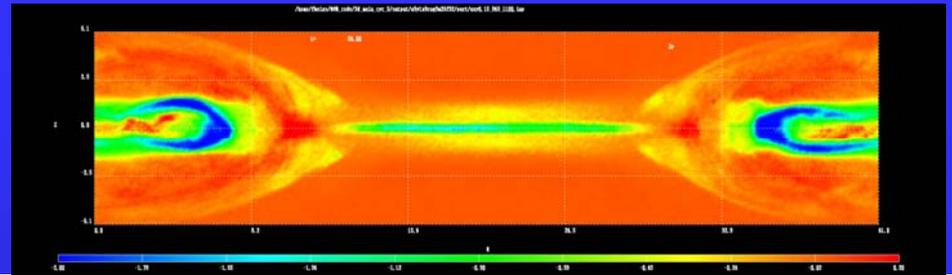
Meandering scale \sim Wavelength of the kink mode



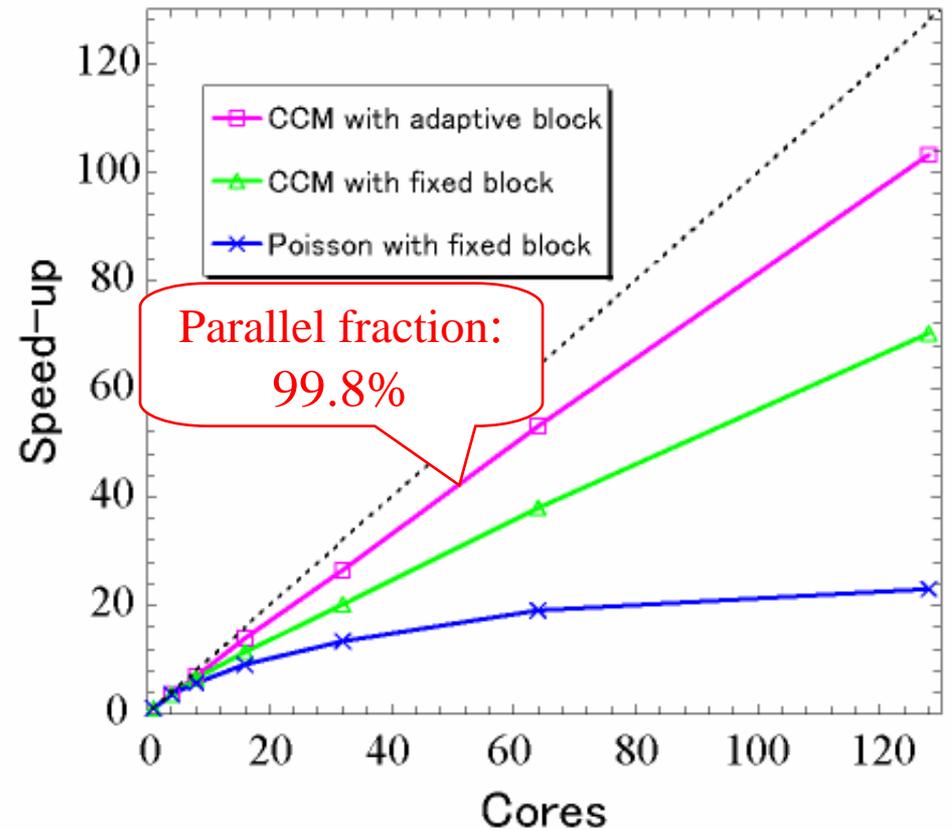
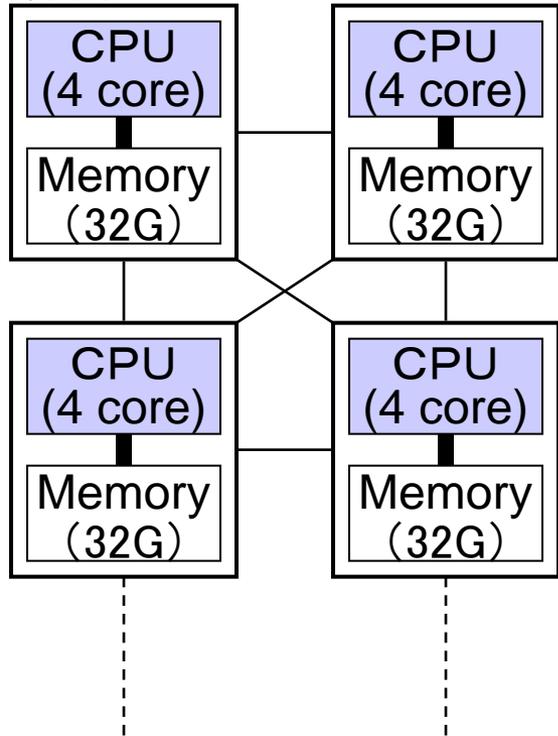
Does the wave-particle interaction still occur in higher mass ratio cases?

Massively parallel AMR-PIC code

Fujitsu FX1



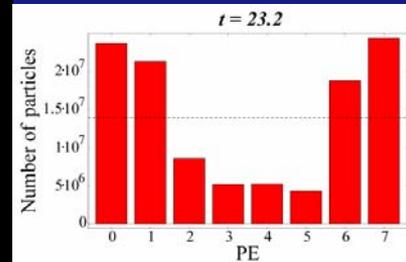
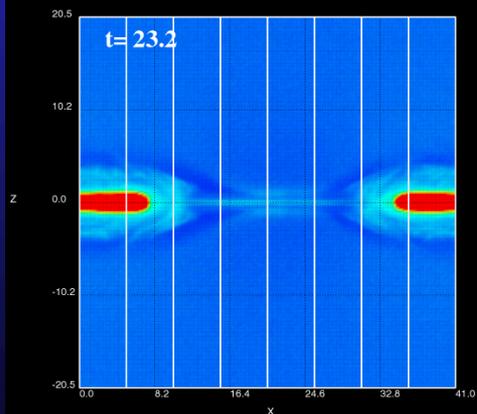
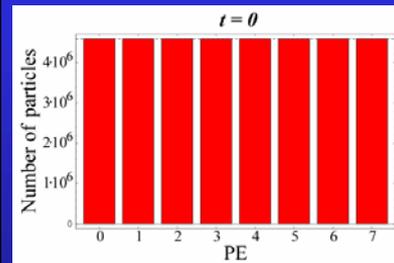
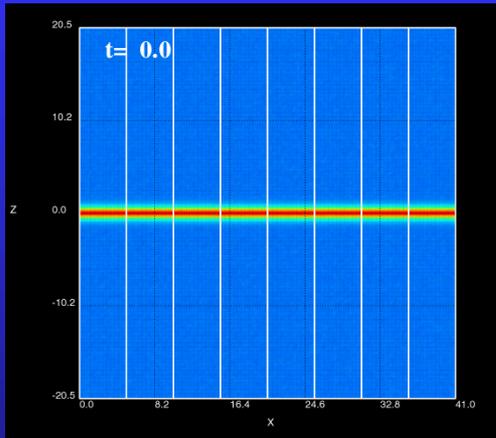
Node



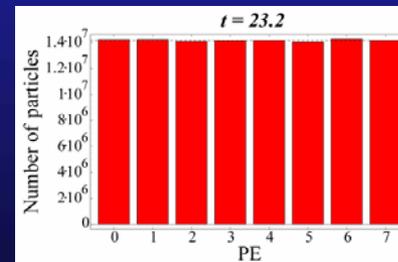
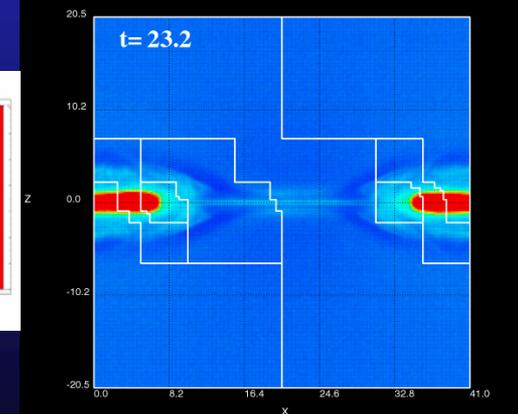
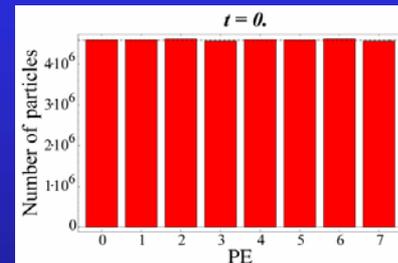
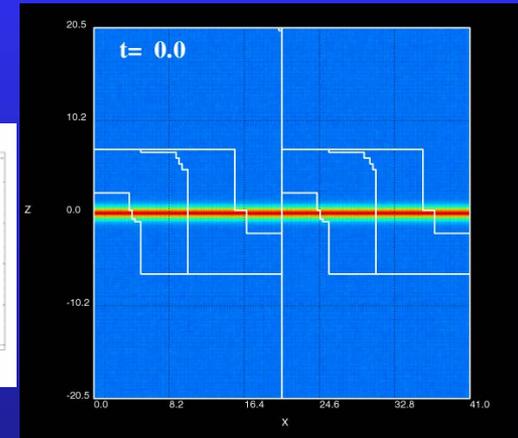
Load Balancing Technique

Example of 8-node parallelization

Fixed block

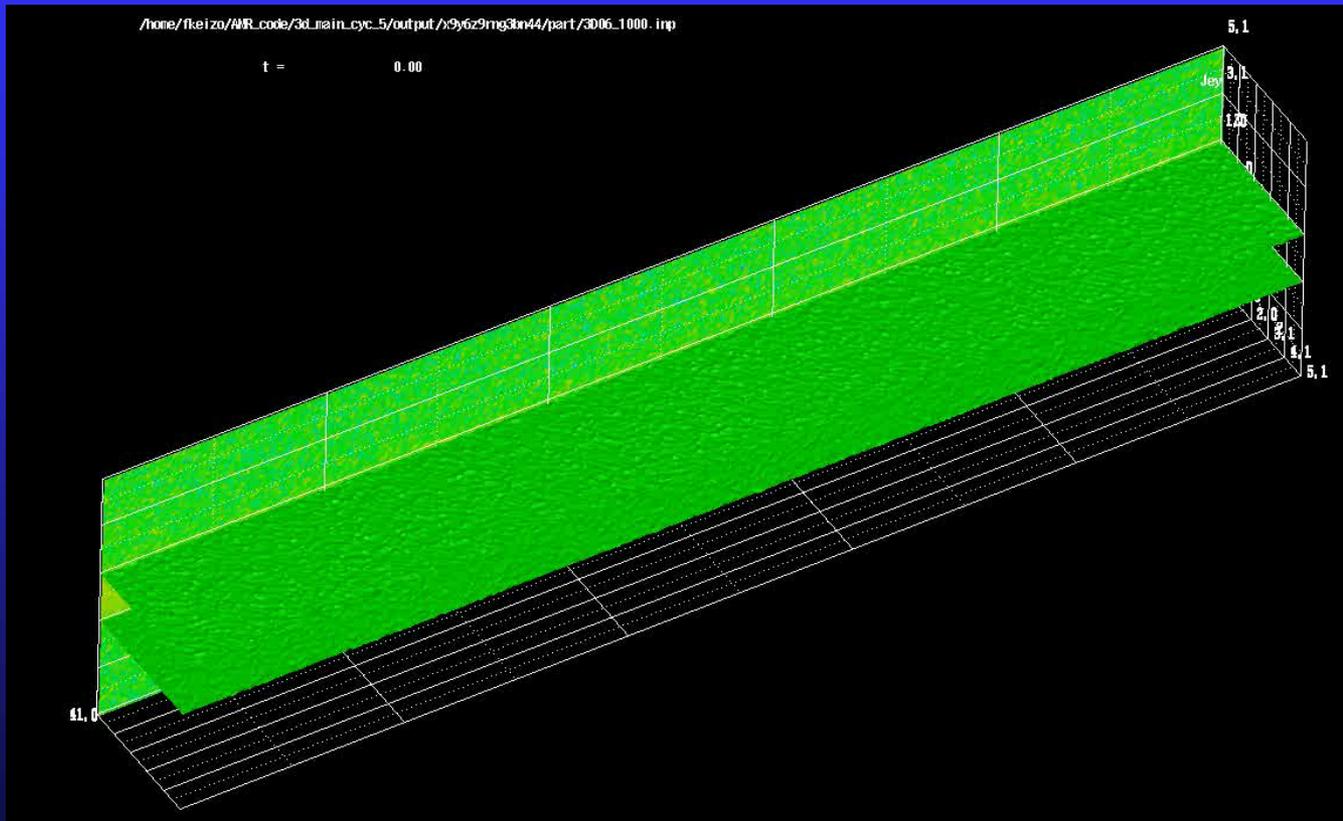


Adaptive block

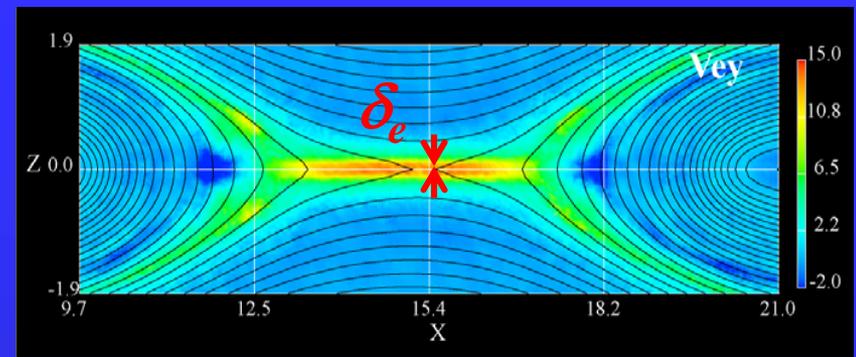


Large-scale 3D simulation: First result

$$m_i/m_e = 100, \quad L_x \times L_y \times L_z = 41\lambda_i \times 5.1\lambda_i \times 41\lambda_i$$

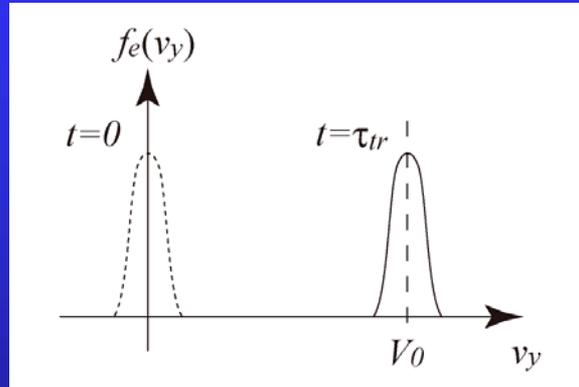


Current Sheet Width (δ_e)



2D case

Inertia
resistivity

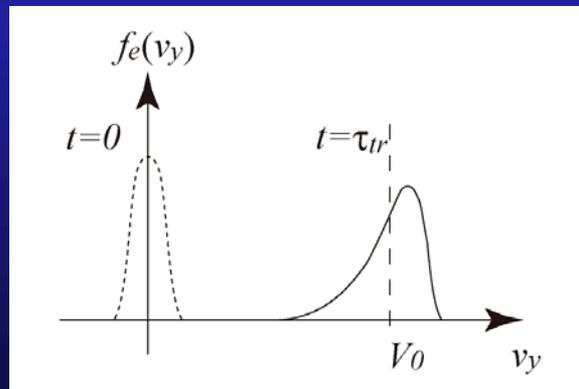


$$\frac{e|E_y|}{m_e} \tau_{tr} = V_0 \quad \tau_{tr} \sim \frac{\delta_e}{V_{in}}$$

→ $\delta_e \sim c/\omega_{pe}$

3D case

Inertia +
anomalous
resistivity

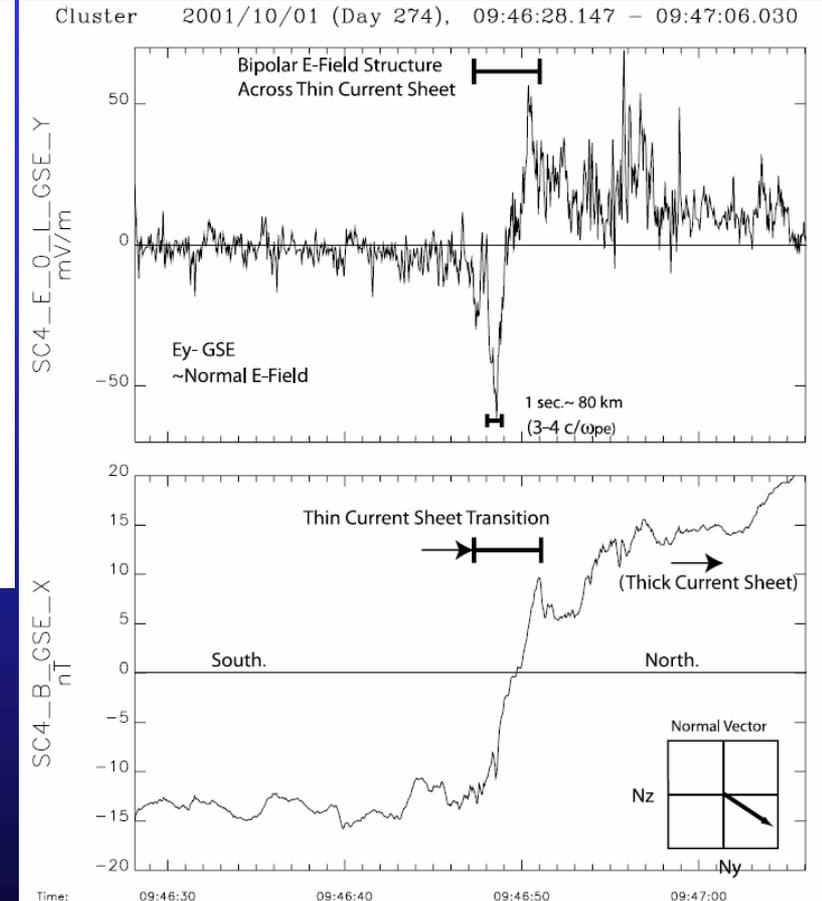
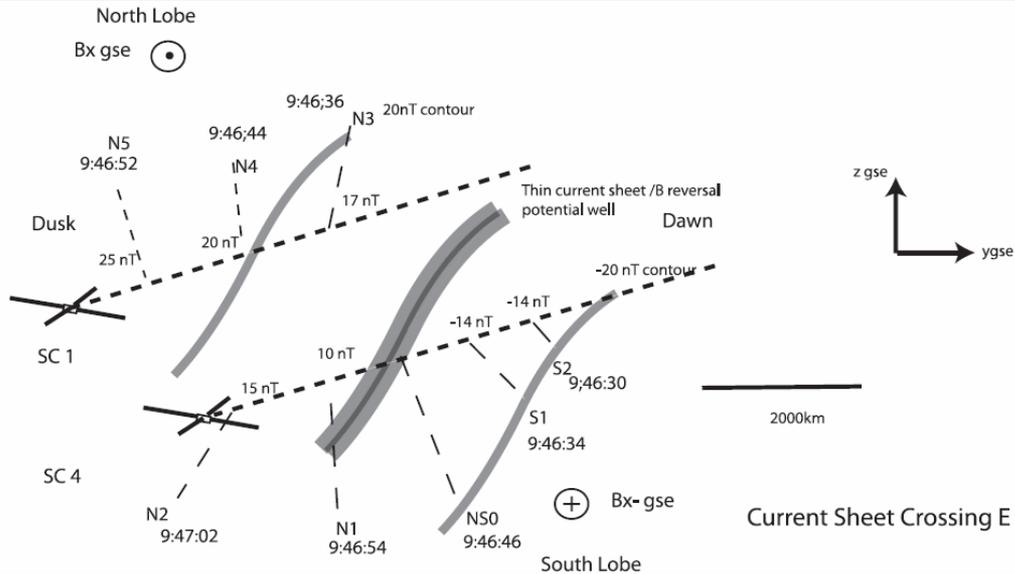


$$\frac{e|E_y|}{m_e} \tau_{tr} > V_0 \quad \tau_{tr} \sim \frac{\delta_e}{V_{in}}$$

→ $\delta_e > c/\omega_{pe}$

In the Earth Magnetotail

Cluster observation of normal electric field in the kinked current sheet [Wygant et al., 2005]



Spiky normal electric field

→ Duration = 1 sec ~ 80 km

→ $3 - 4 c/\omega_{pe} > c/\omega_{pe}$

Summary and Conclusions

The present study has investigated the dissipation mechanism in 3D magnetic reconnection in comparison with 2D reconnection, using a large-scale PIC simulations.

- Reconnection rate

$E_R \sim 0.1$ both the cases of 2D and 3D reconnections

- Dissipation mechanism

2D reconnection → Inertia resistivity

3D reconnection → Inertia resistivity +
Anomalous resistivity (Electron heating
due to wave-particle interaction)

→ Current sheet width larger than c/ω_{pe}

Both the 3D simulation and observation studies indicate the existence of some wave-particle interaction at the X-line.